

# THE CONTINUITY EQUATION AND PARTICLE MOTION

TRISTRAM DE PIRO

ABSTRACT. We define a diffusion using nonstandard analysis and the formula, used in [3];

$$\rho \bar{v} = \bar{J} (*)$$

where  $\rho$  is charge density and  $\bar{J}$  is current, to define the individual particle velocities, with the initial charge distribution given by  $\rho_0$ . We prove that given  $\bar{J}$ , with  $\rho'$  defined by the diffusion, that the continuity equation;

$$\frac{\partial \rho'}{\partial t} + \text{div}(\bar{J}) = 0$$

is satisfied. It follows inductively, as  $\rho'_0 = \rho_0$ , that  $\rho = \rho'$ . The result is useful for answering questions arising from [6].

## 1. THE CONTINUITY EQUATION

**Definition 1.1.** We let  $S(\mathcal{R}^3 \times \mathcal{R}_{>0}) = \{f \in C^\infty(\mathcal{R}^3 \times \mathcal{R}_{>0}) : \text{there exist constants } C_{\bar{I}}, D_{\bar{I}} \geq 0 \text{ with } \bar{I} \in \mathcal{Z}_{\geq 0}^3 \text{ and } |\frac{\partial f}{\partial x_{\bar{I}}}| \leq \frac{C_{\bar{I}}}{|\bar{x}|^2}, \text{ for } |\bar{x}| \geq D_{\bar{I}}\}$

**Lemma 1.2.** Let  $\{\rho, j_1, j_2, j_3\} \subset S(\mathcal{R}^3 \times \mathcal{R}_{>0})$ , with  $\bar{J} = (j_1, j_2, j_3)$ , such that the pair  $(\rho, \bar{J})$  satisfies the continuity equation;

$$\frac{\partial \rho}{\partial t} = -\text{div}(\bar{J})$$

Then, for  $t_0 \in \mathcal{R}_{>0}$ , there exists a constant  $c_0 \geq 0$ , such that  $\rho + c_0 > 1$  and  $(\rho + c_0, \bar{J})$  satisfies the continuity equation on  $\mathcal{R}^3 \times (0, t_0)$ .

*Proof.* As  $\rho \in S(\mathcal{R}^3 \times \mathcal{R}_{>0})$ , taking  $\bar{I} = \bar{0}$ , there exists a ball  $B(0, r_0)$ ,  $r_0 > 0$ , and  $C_{\bar{0}}$  with  $|\rho| \leq \frac{C_{\bar{0}}}{|\bar{x}|^2}$ . As  $\rho \in C^\infty(\mathcal{R}^3 \times \mathcal{R}_{>0})$ ,  $\rho$  is bounded on  $B(0, r_0) \times (0, t_0)$  and therefore on  $\mathcal{R}^3 \times (0, t_0)$ . In particular, there exists a constant  $d > 0$  with  $|\rho| < d$  on  $\mathcal{R}^3 \times (0, t_0)$ . Then  $\rho + d > 0$ ,

$\rho + d + 1 > 1$ , and clearly  $\frac{\partial(\rho+c)}{\partial t} = \frac{\partial\rho}{\partial t} = -\text{div}(\bar{J})$  on  $\mathcal{R}^3 \times (0, t_0)$ , where  $c = d + 1$ , as required.  $\square$

**Definition 1.3.** Let  $(\rho, \bar{J})$  satisfy the continuity equation on  $\mathcal{R}^3 \times (0, t_0)$ , with  $\rho > 1$ . Fix  $\{\eta, \nu\}$  positive infinite, with  $\eta$  odd,  $\nu$  an integer,  $\epsilon$  infinitesimal. We let;

$$\mathcal{T}_\nu = \bigcup_{0 \leq j \leq \nu^2 - 1} \left[ \frac{j}{\nu}, \frac{j+1}{\nu} \right)$$

$$\mathcal{T}_{\nu,0} = \bigcup_{0 \leq j \leq \lfloor t_0 \nu \rfloor} \left[ \frac{j}{\nu}, \frac{j+1}{\nu} \right)$$

$$\mathcal{R}_\eta = \bigcup_{-\frac{(\eta-1)}{2} \leq i \leq \frac{(\eta-1)}{2}} \left[ \frac{i}{\sqrt{\eta}}, \frac{i+1}{\sqrt{\eta}} \right) \cup \left( -\frac{(\eta+1)}{2}, -\frac{(\eta-1)}{2} \right)$$

We define the integer part  $[, ]$  for  $x \in \mathcal{R}$ , by;

$$[x] = \max(y : y \in \mathcal{Z}, y \leq x) \text{ if } x \geq 0$$

$$[x] = \min(y : y \in \mathcal{Z}, y \geq x) \text{ if } x < 0$$

and extend the definition to  ${}^*\mathcal{R}$  by transfer. If  $\bar{z} \in {}^*\mathcal{R}^3$ , we define,  $(\bar{z})_\eta = \bar{y}$ , where  $y_i = \frac{\lfloor z_i \sqrt{\eta} \rfloor}{\sqrt{\eta}}$ . Note that  $|(\bar{z})_\eta| \leq |\bar{z}|$ , as  $|\frac{\lfloor z_i \sqrt{\eta} \rfloor}{\sqrt{\eta}}| \leq |z_i|$ , for  $1 \leq i \leq 3$ .

We define

$$\bar{J}_\eta(x_1, x_2, x_3, t) = \bar{J}^* \left( \frac{\lfloor x_1 \sqrt{\eta} \rfloor}{\sqrt{\eta}}, \frac{\lfloor x_2 \sqrt{\eta} \rfloor}{\sqrt{\eta}}, \frac{\lfloor x_3 \sqrt{\eta} \rfloor}{\sqrt{\eta}}, \frac{\lfloor t \nu \rfloor}{\nu} \right)$$

$$\rho_\eta(x_1, x_2, x_3, t) = \rho^* \left( \frac{\lfloor x_1 \sqrt{\eta} \rfloor}{\sqrt{\eta}}, \frac{\lfloor x_2 \sqrt{\eta} \rfloor}{\sqrt{\eta}}, \frac{\lfloor x_3 \sqrt{\eta} \rfloor}{\sqrt{\eta}}, \frac{\lfloor t \nu \rfloor}{\nu} \right)$$

where  $\{\rho^*, \bar{J}^*\}$  are the transfers of  $\{\rho, \bar{J}\}$  to  $(\mathcal{R}^*)^3 \times \mathcal{T}_{\geq 0}^*$

For  $\bar{x} \in \mathcal{R}_\eta^3$ , with  $\frac{\lfloor x_i \sqrt{\eta} \rfloor}{\sqrt{\eta}} = x_i$ , for  $1 \leq i \leq 3$ , we let;

$$N_{\epsilon, \bar{x}} = \left[ \frac{\rho_\eta(\bar{x}, 0)}{\epsilon} \right]$$

$$M_\epsilon = \sum_{\bar{x} \in \mathcal{R}_\eta^3, \frac{\lfloor x_i \sqrt{\eta} \rfloor}{\sqrt{\eta}} = x_i, 1 \leq i \leq 3} N_{\epsilon, \bar{x}}$$

For  $\bar{x} \in \mathcal{R}_\eta^3$  finite, with  $\frac{\lfloor x_i \sqrt{\eta} \rfloor}{\sqrt{\eta}} = x_i$ , for  $1 \leq i \leq 3$ , define a trajectory  $\bar{s}_{\bar{x}}$  inductively on  $\mathcal{T}_{\nu,0}$  by;

$$\bar{s}_x(\frac{j+1}{\nu}) - \bar{s}_x(\frac{j}{\nu}) = \frac{1}{\nu} \frac{\bar{J}_\eta(\bar{s}_x(\frac{j}{\nu}), \frac{j}{\nu})}{\rho_\eta(\bar{s}_x(\frac{j}{\nu}), \frac{j}{\nu})}, \text{ for } 0 \leq j \leq [t\nu] - 1$$

$$\bar{s}_x(0) = \bar{x}$$

**Lemma 1.4.** *For  $t_0$  finite, there exists a constant  $G \geq 0$  with  $|\frac{\bar{J}_\eta}{\rho_\eta}| \leq G$  on  $\mathcal{R}_\eta^3 \times [0, t_0]$ .*

*Proof.* As  $\rho > 1$ , and  $\bar{J} \in S(\mathcal{R}^3 \times \mathcal{R}_{>0})$ , for  $B(\bar{0}, 1)$ , using smoothness of  $\bar{J}$  on  $B(\bar{0}, 1) \times [0, t_0]$ , and the fact that  $|\frac{1}{\rho}| \leq 1$ , that there exists a constant  $G_1$  with  $|\frac{\bar{J}}{\rho}| \leq G_1$  on  $B(\bar{0}, 1) \times [0, t_0]$ . Moreover, as  $\bar{J} \in S(\mathcal{R}^3 \times \mathcal{R}_{>0})$  and  $|\frac{1}{\rho}| \leq 1$ , we can find a constant  $G_2$ , with  $|\frac{\bar{J}}{\rho}| \leq \frac{G_1}{|\bar{x}|^2} \leq G_1$  on  $B(\bar{0}, 1)^c \times [0, t_0]$ . Then the result follows by transfer.  $\square$

**Definition 1.5.** *Let  $R = \{\bar{x} \in \mathcal{R}_\eta^3 : \frac{[x_i \sqrt{\eta}]}{\sqrt{\eta}} = x_i, \text{ for } 1 \leq i \leq 3, \text{ with } \bar{x} \text{ finite}\}$ .*

*Let  $T_\nu = \{t \in \mathcal{T}_\nu : [t\nu] = t\nu\}$*

**Lemma 1.6.** *Trajectories are bounded*

*Let  $\bar{x} \in R$ , and let  $s_{\bar{x}}$  be a trajectory, then there exists a constant  $D \in \mathcal{R}$ , with  $D \geq 0$ , such that;*

$$|\bar{s}_x(\frac{j}{\nu})| \leq D, \text{ for } 0 \leq j \leq [t\nu].$$

*Proof.* We have, using Lemma 1.4, that;

$$\begin{aligned} |\bar{s}_x(\frac{j}{\nu}) - \bar{s}_x(0)| &= |\sum_{0 \leq k \leq j-1}^* (\bar{s}_x(\frac{k+1}{\nu}) - \bar{s}_x(\frac{k}{\nu}))| \\ &\leq \sum_{0 \leq k \leq j-1}^* |\bar{s}_x(\frac{k+1}{\nu}) - \bar{s}_x(\frac{k}{\nu})| \\ &= \frac{1}{\nu} \sum_{0 \leq k \leq j-1}^* \left| \frac{\bar{J}_\eta(\bar{s}_x(\frac{k}{\nu}), \frac{k}{\nu})}{\rho_\eta(\bar{s}_x(\frac{k}{\nu}), \frac{k}{\nu})} \right| \\ &\leq \frac{Gj}{\nu} \leq Gt_0 \end{aligned}$$

so that;

$$\begin{aligned} |\bar{s}_x(\frac{j}{\nu})| &\leq |\bar{s}_x(0)| + Gt_0 \\ &= \bar{x} + Gt_0 = D \end{aligned}$$

where  $|\frac{\bar{J}_\eta}{\rho_\eta}| \leq G$  on  $\mathcal{R}_\eta^3 \times \mathcal{T}_{\nu,0}$ .

□

**Lemma 1.7.** *We have, for  $t_0$  finite and fixed  $\{\bar{x}, \bar{y}\} \subset \mathcal{R}^3$ ,  $1 \leq i \leq 3$ , that  $\frac{j_{i,0}}{\rho_0}(\bar{y}) = \frac{j_{i,0}}{\rho_0}(\bar{x}) + \epsilon_i + \delta_i$ , where  $|\epsilon_i| \leq D_i|\bar{z}|$ ,  $|\delta_i| \leq F_i|\bar{z}|^2$ ,  $\bar{z} = \bar{y} - \bar{x}$ , and  $D_i, F_i \geq 0$ , and  $(\frac{j_{1,0}}{\rho_0}, \frac{j_{2,0}}{\rho_0}, \frac{j_{3,0}}{\rho_0}) = \frac{\bar{J}}{\rho}|_{t_0}$ . Moreover, varying  $\{\bar{x}, \bar{y}\}$  in a ball  $B(\bar{0}, r)$ , we have that  $\frac{j_{i,0}}{\rho_0}(\bar{y}) = \frac{j_{i,0}}{\rho_0}(\bar{x}) + \epsilon_{i,\bar{x},\bar{y}} + \delta_{i,\bar{x},\bar{y}}$ , where  $|\epsilon_{i,\bar{x},\bar{y}}| \leq H_r|\bar{z}|$ ,  $|\delta_{i,\bar{x},\bar{y}}| \leq K_r|\bar{z}|^2$ , and  $H_r, K_r \geq 0$ . For  $\{\bar{x}, \bar{y}\} \subset B^*(\bar{0}, r)$ ,  $(\frac{j_{i,0}}{\rho_0})_\eta(\bar{y}) = (\frac{j_{i,0}}{\rho_0})_\eta(\bar{x}) + \epsilon_{\eta,i,\bar{x},\bar{y}} + \delta_{\eta,i,\bar{x},\bar{y}}$ , where  $|\epsilon_{\eta,i,\bar{x},\bar{y}}| \leq H_r|\bar{z}_\eta|$ ,  $|\delta_{\eta,i,\bar{x},\bar{y}}| \leq K_r|\bar{z}_\eta|^2$ . We have, for  $t_0$  finite and fixed  $\{\bar{x}, \bar{y}\} \subset \mathcal{R}^3$ ,  $1 \leq i \leq 3$ , that  $|\frac{j_{i,0}}{\rho_0}(\bar{y}) - \frac{j_{i,0}}{\rho_0}(\bar{x})| \leq M_i|\bar{z}|$ ,  $M_i \geq 0$ . Moreover, varying  $\{\bar{x}, \bar{y}\}$  in a ball  $B(\bar{0}, r)$ , we have that  $M_i \leq L_r$ ,  $L_r \geq 0$ . For  $\{\bar{x}, \bar{y}\} \subset B^*(\bar{0}, r)$ ,  $|(\frac{j_{i,0}}{\rho_0})_\eta(\bar{y}) - (\frac{j_{i,0}}{\rho_0})_\eta(\bar{x})| \leq L_r|\bar{z}_\eta|$ . Finally, there exists a constant  $A_r \geq 0$ , with  $|(\frac{j_i}{\rho})(\bar{y}, t) - (\frac{j_i}{\rho})(\bar{x}, t)| \leq A_r|\bar{z}|$ , for  $\{\bar{x}, \bar{y}\} \subset B(\bar{0}, r)$  and  $0 \leq t \leq t_0$ , and with  $|(\frac{j_i}{\rho})_\eta(\bar{y}, t) - (\frac{j_i}{\rho})_\eta(\bar{x}, t)| \leq A_r|\bar{z}_\eta|$ , for  $\{\bar{x}, \bar{y}\} \subset B^*(\bar{0}, r)$  and  $0 \leq t \leq t_0$ .*

*Proof.* For the first part, let  $f_i(t) = \frac{j_{i,0}}{\rho_0}(\bar{x} + t(\bar{y} - \bar{x}))$  and  $g_i(t) = f_i(t) - f_i(0) - f'_i(t)t$ . Then  $g_i(0) = 0$ ,  $g_i(1) - g_i(0) = g_i(1) = f_i(1) - f_i(0) - f'_i(1) = \frac{j_i}{\rho_0}(\bar{y}) - \frac{j_i}{\rho_0}(\bar{x}) - \epsilon_i$ . By the mean value theorem, we have that  $\delta_i = g_i(1) - g_i(0) = g'_i(c)$ , for some  $c \in (0, 1)$ . We have that  $g'_i(t) = f'_i(t) - f''_i(t)t - f'_i(t) = -f''_i(t)t$ , so that  $|g'_i(c)| \leq \max_{c \in (0,1)} |f''_i(c)c| \leq \max_{c \in (0,1)} |f''_i(c)|$ . By the chain rule;

$$|\epsilon_i| = |f'_i(1)| = |\nabla(\frac{j_{i,0}}{\rho_0})|_{\bar{y}} \cdot (\bar{y} - \bar{x})| \leq |\nabla(\frac{j_{i,0}}{\rho_0})|_{\bar{y}}||\bar{z}| = D_i|\bar{z}|$$

$$\text{where } D_i = |\nabla(\frac{j_i}{\rho_0})|_{\bar{y}}|$$

$$\begin{aligned} |\delta_i| &\leq \max_{c \in (0,1)} |f''_i(c)| \\ &= \max_{c \in (0,1)} |(\nabla(\frac{j_{i,0}}{\rho_0})|_{\bar{x}+t(\bar{y}-\bar{x})} \cdot (\bar{y} - \bar{x}))'(c)| \\ &= \max_{c \in (0,1)} |\sum_{j=1}^3 (\nabla(\frac{\partial(\frac{j_{i,0}}{\rho_0})}{\partial x_j})|_{\bar{x}+c(\bar{y}-\bar{x})} \cdot \bar{z})z_j| \\ &\leq |\bar{z}| \max_{c \in (0,1)} \sum_{j=1}^3 |\nabla(\frac{\partial(\frac{j_{i,0}}{\rho_0})}{\partial x_j})|_{\bar{x}+c(\bar{y}-\bar{x})} \cdot \bar{z}| \\ &\leq |\bar{z}|^2 \max_{c \in (0,1)} \sum_{j=1}^3 G_{ij}(c) \end{aligned}$$

$$= F_i |\bar{z}|^2$$

where  $G_{ij}(c) = |\nabla (\frac{j_{i,0}}{\rho_0})|_{\bar{x}+c(\bar{y}-\bar{x})}|$  and  $F_i = \max_{c \in (0,1)} \sum_{j=1}^3 G_{ij}(c)$

For the second part, we can use the proof of the first part, the fact that  $B(\bar{0}, r)$  is convex, and take;

$$H_r = \max_{\bar{w} \in B(\bar{0}, r)} |\nabla (\frac{j_i}{\rho_0})|_{\bar{w}}|, K_r = \max_{\bar{w} \in B(\bar{0}, r)} \sum_{j=1}^3 |\nabla (\frac{j_{i,0}}{\rho_0})|_{\bar{w}}|$$

For the third part, we obtain, by transfer, on  $B^*(0, r)$ , that;

$$(\frac{j_{i,0}}{\rho_0})^*(\bar{y}) = (\frac{j_{i,0}}{\rho_0}(\bar{x}))^* + \epsilon_{i,\bar{x},\bar{y}}^* + \delta_{i,\bar{x},\bar{y}}^*$$

where  $|\epsilon_{i,\bar{x},\bar{y}}^*| \leq H_r |\bar{z}^*|$ ,  $|\delta_{i,\bar{x},\bar{y}}^*| \leq K_r |\bar{z}^*|^2$

In particular, we have that;

$$(\frac{j_{i,0}}{\rho_0})_{\eta}(\bar{y}) = (\frac{j_{i,0}}{\rho_0})_{\eta}(\bar{x}) + \epsilon_{\eta,i,\bar{x},\bar{y}} + \delta_{\eta,i,\bar{x},\bar{y}}$$

where  $|\epsilon_{\eta,i,\bar{x},\bar{y}}| \leq H_r |\bar{z}_{\eta}|$ ,  $|\delta_{\eta,i,\bar{x},\bar{y}}| \leq K_r |\bar{z}_{\eta}|^2$

For the fourth part, with notation as above, we have by the intermediate value theorem, that;

$$\begin{aligned} |f_i(1) - f_i(0)| &\leq \max_{c \in (0,1)} |f'_i(c)| \\ &= \max_{c \in (0,1)} |\nabla (\frac{j_{i,0}}{\rho_0})|_{\bar{x}+c(\bar{y}-\bar{x})} \cdot (\bar{y} - \bar{x})| \\ &\leq \max_{c \in (0,1)} |\nabla (\frac{j_{i,0}}{\rho_0})|_{\bar{x}+c(\bar{y}-\bar{x})}| |\bar{z}| \\ &= M_i |\bar{z}| \end{aligned}$$

where  $M_i = \max_{c \in (0,1)} |\nabla (\frac{j_{i,0}}{\rho_0})|_{\bar{x}+c(\bar{y}-\bar{x})}|$ . For the fifth part, using again the fact that  $B(0, r)$  is convex, and using the proof of the fourth part, we can take  $L_r = \max_{\bar{w} \in B(0, r)} |\nabla (\frac{j_{i,0}}{\rho_0})|_{\bar{w}}|$ . For the last part, we have, by transfer, on  $B^*(0, r)$ , that;

$$|(\frac{j_{i,0}}{\rho_0})^*(\bar{y}) - (\frac{j_{i,0}}{\rho_0}(\bar{x}))^*| \leq L_r |\bar{z}^*|$$

In particular, we have that;

$$|(\frac{j_{i,0}}{\rho_0})_\eta(\bar{y}) - (\frac{j_{i,0}}{\rho_0})_\eta(\bar{x})| \leq L_r |\bar{z}_\eta|$$

as required. For the final part, let  $A_r = \max_{\bar{w} \in B(0,r), 0 \leq t \leq t_0} |\nabla(\frac{j_i}{\rho})|_{(\bar{w},t)}$ , to deduce, as before, that;

$$|(\frac{j_i}{\rho})_\eta(\bar{y}, t) - (\frac{j_i}{\rho})_\eta(\bar{x}, t)| \leq A_r |\bar{z}_\eta|, \text{ for } 0 \leq t \leq t_0$$

□

**Lemma 1.8.** *Trajectories are injective on Initial Conditions*

For  $\{\bar{x}_1, \bar{x}_2\} \subset R$ , with  $\bar{x}_1 \neq \bar{x}_2$ , then we have that;

$$\bar{s}_{\bar{x}_1}(\frac{j}{\nu}) \neq \bar{s}_{\bar{x}_2}(\frac{j}{\nu}), \text{ for } 0 \leq j \leq [t\nu]$$

*Proof.* Suppose not, and choose  $t_1 = \frac{j_1}{\nu}$  least, with  $1 \leq j_1 \leq [t\nu]$ , such that  $\bar{s}_{\bar{x}_1}(t_1) = \bar{s}_{\bar{x}_2}(t_1)$ , so that  $\bar{s}_{\bar{x}_1}(t_1 - \frac{1}{\nu}) \neq \bar{s}_{\bar{x}_2}(t_1 - \frac{1}{\nu})$ . By the definition of trajectories, we have that;

$$\begin{aligned} \bar{s}_{\bar{x}_1}(t_1) - \bar{s}_{\bar{x}_2}(t_1) &= \bar{s}_{\bar{x}_1}(t_1 - \frac{1}{\nu}) - \bar{s}_{\bar{x}_2}(t_1 - \frac{1}{\nu}) \\ &+ \frac{1}{\nu} \left( \frac{\bar{J}_\eta(\bar{s}_{\bar{x}_1}(t_1 - \frac{1}{\nu}), t_1 - \frac{1}{\nu})}{\rho_\eta(\bar{s}_{\bar{x}_1}(t_1 - \frac{1}{\nu}), t_1 - \frac{1}{\nu})} - \frac{\bar{J}_\eta(\bar{s}_{\bar{x}_2}(t_1 - \frac{1}{\nu}), t_1 - \frac{1}{\nu})}{\rho_\eta(\bar{s}_{\bar{x}_2}(t_1 - \frac{1}{\nu}), t_1 - \frac{1}{\nu})} \right) \\ &= \bar{z} + \frac{1}{\nu} (\bar{f} + \bar{\delta}) = \bar{0} \end{aligned}$$

where, with the constants  $\{H_r, K_r, D_{t_1 - \frac{1}{\nu}}\}$  determined by Lemmas 1.7 and 1.6;

$\bar{z} = \bar{s}_{\bar{x}_1}(t_1 - \frac{1}{\nu}) - \bar{s}_{\bar{x}_2}(t_1 - \frac{1}{\nu})$ ,  $|\bar{f}| \leq D|\bar{z}_\eta|$ ,  $|\bar{\delta}| \leq F|\bar{z}_\eta|^2$ , for some  $\{D, F\} \subset \mathcal{R}$  and  $D, F \geq 0$ ,  $D = 3H_{2D_{t_1 - \frac{1}{\nu}}}$ ,  $F = 3K_{2D_{t_1 - \frac{1}{\nu}}}$ . It follows that;

$$\begin{aligned} \bar{z} + \frac{\bar{f}}{\nu} &= -\frac{\bar{\delta}}{\nu} \\ \|\bar{z}\| - \|\frac{\bar{f}}{\nu}\| &\leq \frac{F|\bar{z}_\eta|^2}{\nu} \\ \|\bar{z}\| &\leq \frac{D|\bar{z}_\eta|}{\nu} + \frac{F|\bar{z}_\eta|^2}{\nu} \end{aligned}$$

$$1 \leq \frac{D|\bar{z}_\eta|}{\nu|\bar{z}|} + \frac{F|\bar{z}_\eta|^2}{\nu|\bar{z}|} \simeq 0$$

which is a contradiction.  $\square$

**Lemma 1.9.** *Trajectories are  $S$ -continuous on initial conditions*

Let  $\{\bar{x}_1, \bar{x}_2\} \subset R$ , with  $\bar{x}_1 \simeq \bar{x}_2$ , then for finite  $t \in \mathcal{T}_\nu$ , we have that  $s_{\bar{x}_1}(t) \simeq s_{\bar{x}_2}(t)$ . Moreover,  $|s_{\bar{x}_1}(t) - s_{\bar{x}_2}(t)| \leq (e^{A_{D_t} \circ t} + 1)|\bar{x}_1 - \bar{x}_2|$ , where  $A_{D_t}$  is the constant from Lemma 1.6.

*Proof.* We have, using the definition of trajectories, with constants  $\{D_t, A_{D_t}\}$  from Lemmas 1.6, 1.7 and the remark in Definition 1.3 that, for  $0 \leq i \leq [t\nu] - 1$ ;

$$\begin{aligned} |s_{\bar{x}_1}(\frac{i+1}{\nu}) - s_{\bar{x}_2}(\frac{i+1}{\nu})| &= |s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu}) + \frac{1}{\nu}(\frac{\bar{J}_\eta}{p_\eta}(\bar{x}_1(\frac{i}{\nu}), \frac{i}{\nu}) - \frac{\bar{J}_\eta}{p_\eta}(\bar{x}_2(\frac{i}{\nu}), \frac{i}{\nu}))| \\ &\leq |s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu})| + \frac{1}{\nu}|(\frac{\bar{J}_\eta}{p_\eta}(\bar{x}_1(\frac{i}{\nu}), \frac{i}{\nu}) - \frac{\bar{J}_\eta}{p_\eta}(\bar{x}_2(\frac{i}{\nu}), \frac{i}{\nu}))| \\ &\leq |s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu})| + \frac{A_{D_t}}{\nu}|(s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu}))_\eta| \\ &= (1 + \frac{A_{D_t}}{\nu})|(s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu}))_\eta| \\ &\leq (1 + \frac{A_{D_t}}{\nu})|s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu})| \end{aligned}$$

Suppose that  $|s_{\bar{x}_1}(0) - s_{\bar{x}_2}(0)| = |\bar{x}_1 - \bar{x}_2| = \epsilon \simeq 0$ , then, by a simple induction, we obtain that, for  $0 \leq i \leq [t\nu]$ ;

$$|s_{\bar{x}_1}(\frac{i}{\nu}) - s_{\bar{x}_2}(\frac{i}{\nu})| \leq (1 + \frac{A_{D_t}}{\nu})^{i-1} \epsilon, (*)$$

We have that  $((1 + \frac{A_{D_t}}{\nu})^{i-1}) = ((1 + \frac{A_{D_t}}{\nu})^\nu)^{\frac{i-1}{\nu}}$ , with  $(1 + \frac{A_{D_t}}{\nu})^\nu \simeq e^{A_{D_t}}$ , so that, as  $\frac{i-1}{\nu}$  is finite,  $(1 + \frac{A_{D_t}}{\nu})^{i-1} \simeq e^{A_{D_t}(\frac{i-1}{\nu})^\circ}$ , which is finite, so that, as  $\epsilon$  is infinitesimal,  $|\bar{x}_1(\frac{i}{\nu}) - \bar{x}_2(\frac{i}{\nu})| \simeq 0$ , as required.

For the final claim, we have by (\*), that;

$$\begin{aligned} |s_{\bar{x}_1}(t) - s_{\bar{x}_2}(t)| &= |s_{\bar{x}_1}(\frac{[t\nu]}{\nu}) - s_{\bar{x}_2}(\frac{[t\nu]}{\nu})| \\ &\leq (1 + \frac{A_{D_t}}{\nu})^{[t\nu]-1} |\bar{x}_1 - \bar{x}_2| \\ &\simeq e^{A_{D_t} \frac{[t\nu]-1}{\nu}} |\bar{x}_1 - \bar{x}_2| \end{aligned}$$

$$\simeq e^{A_{D_t} \circ t} |\bar{x}_1 - \bar{x}_2|$$

It follows that;

$$|s_{\bar{x}_1}(t) - s_{\bar{x}_2}(t)| \leq (e^{A_{D_t} \circ t} + 1) |\bar{x}_1 - \bar{x}_2| \text{ as required.}$$

□

**Definition 1.10.** For  $\bar{x} \in \mathcal{R}_\eta^3$  finite, with  $\frac{[x_i \sqrt{\eta}]}{\sqrt{\eta}} = x_i$ , for  $1 \leq i \leq 3$ ,  $t_0 \in \mathcal{T}_\nu$  finite, define a reverse trajectory  $\bar{s}_{\bar{x}, t_0}$  inductively on  $[0, \frac{[t_0 \nu]}{\nu}]$  by;

$$\bar{s}_{\bar{x}, t_0}(\frac{j+1}{\nu}) - \bar{s}_{\bar{x}, t_0}(\frac{j}{\nu}) = \frac{1}{\nu} \frac{\bar{J}_\eta(\bar{s}_{\bar{x}}(\frac{j}{\nu}), \frac{[t_0 \nu] - j}{\nu})}{\rho_\eta(\bar{s}_{\bar{x}}(\frac{j}{\nu}), \frac{[t_0 \nu] - j}{\nu})}, \text{ for } 0 \leq j \leq [t_0 \nu] - 1$$

$$\bar{s}_{\bar{x}, t_0}(0) = \bar{x}$$

**Lemma 1.11.** Trajectories are essentially onto For  $t_0 \in \mathcal{T}_\nu$  finite,  $\bar{y} \in \mathcal{R}_\eta^3$  finite, there exists  $\bar{x} \in \mathcal{R}_\eta^3$ , with  $\bar{x} \simeq \bar{y}$ , and  $\bar{x}_1 \in \mathcal{R}_\eta^3$  finite, with  $s_{\bar{x}_1}(t_0) = \bar{x}$ . Moreover;

$$|\bar{y} - \bar{x}| \leq (e^{\circ t_0} + 2) \left( \frac{3}{\sqrt{\eta}} + \frac{F}{\nu} \right) + \frac{F e^{\circ t_0}}{\nu}$$

*Proof.* Let  $\bar{s}_{\bar{y}', t_0}$  be the reverse trajectory defined by  $\bar{y}'$  and  $t_0$ , where  $\frac{[y_i \sqrt{\eta}]}{\sqrt{\eta}} = y'_i$ , for  $1 \leq i \leq 3$ , so that  $\bar{y}' \simeq \bar{y}$  and  $\bar{y}'$  is finite. Let  $\bar{x}_1 = \bar{s}_{\bar{y}', t_0}(\frac{[t_0 \nu]}{\nu})$ . By the proof of Lemma 1.6,  $\bar{x}_1$  is finite. Let  $\bar{x}_0$  be defined by  $\frac{[x_{1,i} \sqrt{\eta}]}{\sqrt{\eta}} = x_{0,i}$ , so that  $\bar{x}_1 \simeq \bar{x}_0$ , and  $\bar{x}_0$  is finite. Let  $\bar{x} = \bar{s}_{\bar{x}_0}(\frac{[t_0 \nu]}{\nu})$ , where  $\bar{s}_{\bar{x}_0}$  is the forward trajectory defined by  $\bar{x}_0$ , see Definition 1.3, so that, again by Lemma 1.6,  $\bar{x}$  is finite. We claim that  $\bar{x} \simeq \bar{y}'$ , ( $\dagger$ ), so that  $\bar{x} \simeq \bar{y}$  by transitivity. We have, using Definitions 1.3, 1.10 and Lemma 1.4 that;

$$\begin{aligned} & |\bar{s}_{\bar{y}', t_0}(\frac{[t_0 \nu]}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{1}{\nu})| \\ & \leq |\bar{x}_1 - \bar{x}_0| + |\bar{s}_{\bar{x}_0}(0) - \bar{s}_{\bar{x}_0}(\frac{1}{\nu})| \\ & \leq \frac{3}{\sqrt{\eta}} + \frac{F}{\nu} \simeq 0 \quad (*) \end{aligned}$$

Then, by Definitions 1.3 and 1.10, using the bound  $D_{t_0, \bar{y}', \bar{x}_0}$  from Lemma 1.6 and the constant  $N = A_{D_{t_0, \bar{y}', \bar{x}_0}}$  from Lemma 1.7, we have, for  $0 \leq j \leq [t_0 \nu] - 2$ ;

$$\begin{aligned}
& |\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+2}{\nu})| \\
&= |(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu})) + \frac{1}{\nu}(\frac{\bar{J}_\eta(\bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}), \frac{j+1}{\nu})}{\rho_\eta(\bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}), \frac{j+1}{\nu})} - \frac{\bar{J}_\eta(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}), \frac{j+1}{\nu})}{\rho_\eta(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}), \frac{j+1}{\nu})})| \\
&\leq |(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| + \frac{N}{\nu}|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| \\
&\leq |(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| + \frac{N}{\nu}|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| \\
&\leq |(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| + \frac{N}{\nu}(|\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}) \\
&\quad - \bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu})| + |\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu})|) \\
&= (1 + \frac{N}{\nu})|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| + \frac{N}{\nu}|\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j+1}{\nu}) \\
&\quad - \bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu})| \\
&\leq (1 + \frac{N}{\nu})|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu} - \frac{j}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{j+1}{\nu}))| + \frac{NF}{\nu^2}
\end{aligned}$$

It follows, by a simple induction, and the formula for a geometric progression, that;

$$\begin{aligned}
& |\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{[t_0\nu]}{\nu})| \\
&\leq (1 + \frac{N}{\nu})^{[t_0\nu]-2}|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{1}{\nu}))| + \delta(* \sum_{k=0}^{[t_0\nu]-3} (1 + \frac{N}{\nu})^k) \\
&= (1 + \frac{N}{\nu})^{[t_0\nu]-2}|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{1}{\nu}))| + \frac{\delta\nu}{N}((1 + \frac{N}{\nu})^{[t_0\nu]-2} - 1) (** )
\end{aligned}$$

where  $\delta = \frac{NF}{\nu^2}$ . We have that  $\frac{\delta\nu}{N} = \frac{F}{\nu} \simeq 0$ , and by the proof of Lemma 1.9,  $(1 + \frac{N}{\nu})^{[t_0\nu]-2} \simeq e^{\circ \frac{[t_0\nu]}{\nu} - \frac{2}{\nu}} = e^{\circ t_0}$ . It follows, using (\*) and (\*\*), that;

$$\begin{aligned}
& \frac{\delta\nu}{N}((1 + \frac{N}{\nu})^{[t_0\nu]-2} - 1) \simeq \frac{\delta\nu}{N}(e^{\circ t_0} - 1) \simeq 0 \\
& (1 + \frac{N}{\nu})^{[t_0\nu]-2}|(\bar{s}_{\bar{y}', t_0}(\frac{[t_0\nu]}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{1}{\nu}))| \leq (e^{\circ t_0} + 1)(\frac{3}{\sqrt{\eta}} + \frac{F}{\nu}) \simeq 0 \\
& |\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{[t_0\nu]}{\nu})| \simeq 0 (***)
\end{aligned}$$

We have, using Lemma 1.6 and Definition 1.10, that  $\bar{y}' = \bar{s}_{\bar{y}', t_0}(0) \simeq \bar{s}_{\bar{y}', t_0}(\frac{1}{\nu})$  and  $\bar{x} = \bar{s}_{\bar{x}_0}(\frac{[t_0\nu]}{\nu})$ , so that, by (\*\*\*),  $\bar{y}' \simeq \bar{x}$ , proving (†). For the final claim, using the above calculation;

$$\begin{aligned}
|\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{x}| &= |\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{s}_{\bar{x}_0}(\frac{[t_0\nu]}{\nu})| \\
&\leq (e^{\circ t_0} + 1)(\frac{3}{\sqrt{\eta}} + \frac{F}{\nu}) + \frac{F e^{\circ t_0}}{\nu} \\
|\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{y}'| &= |\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{s}_{\bar{y}', t_0}(0)| \leq \frac{F}{\nu} \\
|\bar{y}' - \bar{y}| &\leq \frac{3}{\sqrt{\eta}}
\end{aligned}$$

so that;

$$\begin{aligned}
|\bar{y} - \bar{x}| &\leq |\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{x}| + |\bar{s}_{\bar{y}', t_0}(\frac{1}{\nu}) - \bar{y}'| + |\bar{y}' - \bar{y}| \\
&\leq (e^{\circ t_0} + 1)(\frac{3}{\sqrt{\eta}} + \frac{F}{\nu}) + \frac{F e^{\circ t_0}}{\nu} + \frac{F}{\nu} + \frac{3}{\sqrt{\eta}} \\
&= (e^{\circ t_0} + 2)(\frac{3}{\sqrt{\eta}} + \frac{F}{\nu}) + \frac{F e^{\circ t_0}}{\nu}
\end{aligned}$$

□

**Lemma 1.12.** *Let  $\{\bar{x}, \bar{y}\} \subset R \cap {}^*B(\bar{0}, r)$ ,  $\bar{x} \neq \bar{y}$ , with  $\nu \geq \eta$  and  $r \in \mathcal{R}_{>0}$ , then for finite  $t \in \mathcal{T}_\nu$ , there exists a constant  $C_{t,r} \in \mathcal{R}_{>0}$ , uniform in  $\{\bar{x}, \bar{y}\}$ , such that  $|s_{\bar{x}}(t) - s_{\bar{y}}(t)| \geq \frac{C_{t,r}}{\sqrt{\eta}}$ .*

*Proof.* First, fix  $\{\bar{x}, \bar{y}\} \subset R \cap {}^*B(\bar{0}, r)$ ,  $\bar{x} \neq \bar{y}$ , we claim that for any infinitesimal  $\epsilon > 0$ ,  $|s_{\bar{x}}(t) - s_{\bar{y}}(t)| \geq \frac{\epsilon}{\sqrt{\eta}}$ , (\*). Given this, we have that  $|s_{\bar{x}}(t) - s_{\bar{y}}(t)|\sqrt{\eta} \geq \frac{1}{n}$  for all infinite  $n \in {}^*\mathcal{N}$ . By underflow, the internal set  $\{n \in {}^*\mathcal{N} : |s_{\bar{x}}(t) - s_{\bar{y}}(t)|\sqrt{\eta} \geq \frac{1}{n}\}$ , must contain a finite element  $C_{t,\bar{x},\bar{y}} \in \mathcal{N}$ , for which  $|s_{\bar{x}}(t) - s_{\bar{y}}(t)| \geq \frac{C_{t,\bar{x},\bar{y}}}{\sqrt{\eta}}$ . To prove (\*), suppose that  $|s_{\bar{x}}(t) - s_{\bar{y}}(t)| < \frac{\epsilon}{\sqrt{\eta}}$ , (†), for some  $\epsilon > 0$  infinitesimal. Let  $\bar{x}_1 = s_{\bar{x}}(t)$  and  $\bar{x}_2 = s_{\bar{y}}(t)$ . By Lemma 1.6,  $\bar{x}_1$  and  $\bar{x}_2$  are finite. We extend the definition of a reverse trajectory in Definition 1.10 to include arbitrary finite starting points in  $\mathcal{R}_\eta^3$ . By a simple adaptation of Lemma 1.9;

$$\begin{aligned}
|s_{\bar{x}_1,t}(t) - s_{\bar{x}_2,t}(t)| &\leq (e^{A_{D_t} \circ t} + 1)|\bar{x}_1 - \bar{x}_2| \\
&< (e^{A_{D_t} \circ t} + 1)\frac{\epsilon}{\sqrt{\eta}}
\end{aligned}$$

By the definition of a trajectory, a reverse trajectory and  $\bar{x}_1$ ;

$$|\bar{x} - s_{\bar{x}_1,t}(t)| = |s_{\bar{x}}(0) - s_{\bar{x}_1,t}(t)|$$

$$|s_{\bar{x}}(t) - s_{\bar{x}_1,t}(0)| = |\bar{x}_1 - \bar{x}_1| = 0$$

By a straightforward adaptation of the calculation in Lemma 1.11, replacing trajectories by reversed trajectories, we have that that;

$$\begin{aligned} & |s_{\bar{x}}(\frac{1}{\nu}) - s_{\bar{x}_1,t}(t)| \\ & \leq (1 + \frac{N}{\nu})^{[t_0\nu]-2} |s_{\bar{x}}(t) - s_{\bar{x}_1,t}(\frac{1}{\nu})| + \frac{F}{N} ((1 + \frac{N}{\nu})^{[t_0\nu]-2} - 1) \\ & \leq (e^{N^\circ t} + 1) |s_{\bar{x}}(t) - s_{\bar{x}_1,t}(\frac{1}{\nu})| + \frac{F}{\nu} (e^{N^\circ t} + 1) \end{aligned}$$

where  $N = A_{D_{t,\bar{x},\bar{x}_1}}$

$$\begin{aligned} & |s_{\bar{x}}(0) - s_{\bar{x}_1,t}(t)| \leq |s_{\bar{x}}(0) - s_{\bar{x}}(\frac{1}{\nu})| + |s_{\bar{x}}(\frac{1}{\nu}) - s_{\bar{x}_1,t}(t)| \\ & \leq \frac{2F}{\nu} + (e^{N^\circ t} + 1) |s_{\bar{x}}(t) - s_{\bar{x}_1,t}(\frac{1}{\nu})| + \frac{F}{\nu} (e^{N^\circ t} + 1) \\ & \leq \frac{2F}{\nu} + (e^{N^\circ t} + 1) (|s_{\bar{x}}(t) - s_{\bar{x}_1,t}(0)| + |s_{\bar{x}_1,t}(0) - s_{\bar{x}_1,t}(\frac{1}{\nu})|) + \frac{F}{\nu} (e^{N^\circ t} + 1) \\ & \leq \frac{2F}{\nu} + (e^{N^\circ t} + 1) (\frac{2F}{\nu}) + \frac{F}{\nu} (e^{N^\circ t} + 1) \\ & = \frac{F}{\nu} (5 + 3e^{N^\circ t}) \end{aligned}$$

to obtain that;

$$|\bar{x} - s_{\bar{x}_1,t}(t)| \leq \frac{F}{\nu} (5 + 3e^{N^\circ t}) \quad (\dagger\dagger)$$

Similarly, we obtain that;

$$|\bar{y} - s_{\bar{x}_2,t}(t)| \leq \frac{F}{\nu} (5 + 3e^{N^\circ t}) \quad (\dagger\dagger\dagger)$$

so that, using  $(\dagger)$ ,  $(\dagger\dagger)$ ,  $(\dagger\dagger\dagger)$ , using the fact that  $2F(5 + 3e^{N^\circ t}) \leq \eta^{\frac{1}{3}}$ , and with the choice  $\nu \geq \eta$ ;

$$\begin{aligned} & |\bar{x} - \bar{y}| \leq |\bar{x} - s_{\bar{x}_1,t}(t)| + |\bar{y} - s_{\bar{x}_2,t}(t)| + |s_{\bar{x}_1,t}(t) - s_{\bar{x}_2,t}(t)| \\ & \leq 2\frac{F}{\nu} (5 + 3e^{N^\circ t}) + \frac{\epsilon}{\sqrt{\eta}} \\ & \leq 2\frac{F}{\eta} (5 + 3e^{N^\circ t}) + \frac{\epsilon}{\sqrt{\eta}} \\ & \leq \eta^{-\frac{2}{3}} + \frac{\epsilon}{\sqrt{\eta}} \end{aligned}$$

$$\leq \frac{1}{2\sqrt{\eta}}$$

which contradicts the fact that  $\{\bar{x}, \bar{y}\} \subset R$ , with  $\bar{x} \neq \bar{y}$ , so that  $|\bar{x} - \bar{y}| \geq \frac{1}{\sqrt{\eta}}$ , proving (\*). It follows that for the internal function  $f_{t,r} : (R \cap {}^*B(\bar{0}, r)) \times (R \cap {}^*B(\bar{0}, r)) \setminus \Delta \rightarrow {}^*\mathcal{R}$ , defined by;

$$f_{t,r}(\bar{x}, \bar{y}) = \sqrt{\eta} |s_{\bar{x}}(t) - s_{\bar{y}}(t)|$$

has the property that;

$f_{t,r} \geq \epsilon$  for all infinitesimals  $\epsilon$ . By underflow, similarly to the above argument, we can find a constant  $C_{t,r} \in \mathcal{R}_{>0}$ , uniform in  $\{\bar{x}, \bar{y}\}$ , with  $f_{t,r} \geq C_{t,r}$  as required.  $\square$

**Definition 1.13.** *We define;*

$$R_\eta = \{\bar{x} \in \mathcal{R}_\eta^3 : [x_i \sqrt{\eta}] = x_i \sqrt{\eta}, \text{ for } 1 \leq i \leq 3\}$$

$$\mathcal{T}_\nu = \bigcup_{0 \leq j \leq \nu^2 - 1} [\frac{j}{\nu}, \frac{j+1}{\nu})$$

*We extend the definition of trajectories, For  $\bar{x} \in \mathcal{R}_\eta$ , by defining;*

$$s_{\bar{x}}(\frac{j+1}{\nu}) = s_{\bar{x}}(\frac{j}{\nu}) + \frac{1}{\nu} \frac{\bar{J}_\eta(\bar{s}_{\bar{x}}(\frac{j}{\nu}), \frac{j}{\nu})}{\rho_\eta(\bar{s}_{\bar{x}}(\frac{j}{\nu}), \frac{j}{\nu})} \pmod{\mathcal{R}_\eta^3}, \text{ for } 0 \leq j \leq \nu^2 - 2$$

*where for  $x \in \mathcal{R}^*$ , we define  $x \pmod{\mathcal{R}_\eta}$ , to be the unique  $y \in \mathcal{R}_\eta$ , for which  $x - y = n\sqrt{\eta}$ , with  $n \in {}^*\mathcal{Z}$ , and for  $\bar{x} \in (\mathcal{R}^*)^3$ , we extend the definition coordinate wise. For  $t \in \mathcal{T}_\nu$ , we define  $s_{\bar{x}}(t) = s_{\bar{x}}(\frac{[t\nu]}{\nu})$ .*

*For  $t \in \mathcal{T}_\nu \setminus [\frac{\nu^2-1}{\nu}, \nu)$ , we define a point  $s_{\bar{x}}(t)$  to be critical, if;*

$$|s_{\bar{x}}(t) - s_{\bar{x}}(t + \frac{1}{\nu})| > 1$$

**Lemma 1.14.** *For  $\bar{x} \in \mathcal{R}_\eta^3$  infinite and finite  $t_0 \in \mathcal{T}_\nu$ , we have that  $s_{\bar{x}}(t_0)$  is infinite.*

*Proof.* Enumerate the \*-finite set of critical points in the trajectory  $s_{\bar{x}}$ , by  $\{s_i : 1 \leq i \leq \kappa\}$ . For any such critical point  $s_{i_0}$ , we must have, using Definition 1.13 and the bound  $G$  in Lemma 1.6, that  $|(s_{i_0})_i| > \frac{\eta-1}{2\sqrt{\eta}} - \frac{G}{\nu}$ , for some  $1 \leq i \leq 3$ . In particular, such critical points in the trajectory are infinite. If  $s_{\bar{x}}(\frac{[t_0\nu]}{\nu})$  is a critical point, the proof is complete. Otherwise, let  $s_{i_1}$  be the greatest element in the \*-finite set

consisting of critical points with the corresponding time  $t_{i_1} \leq \frac{[t_0\nu]}{\nu}$ , and use the proof of Lemma 1.7, to conclude that  $|s_{i_1} - s_{\bar{x}}(\frac{[t_0\nu]}{\nu})| \leq \frac{Gt_0}{\nu}$ . In particular, as  $\frac{Gt_0}{\nu}$  is finite,  $s_{\bar{x}}(t_0)$  is infinite.  $\square$

**Definition 1.15.** For  $t \in \mathcal{T}_\nu$ , we define the range of the trajectories  $R_t$  by;

$$R_t = \{s_{\bar{x}}(t) : \bar{x} \in \mathcal{R}_\eta^3\}$$

It is easy to see, as  $R_\eta$  is  $*$ -finite, and the definition of the trajectories is internal, that  $R_t$  is  $*$ -finite. Moreover, as  $\{t \in \mathcal{T}_\nu : [t\nu] = t\nu\}$  is  $*$ -finite, we have that  $\bigcup_{t \in \mathcal{T}_\nu} R_t$  is  $*$ -finite.

**Lemma 1.16.** For all  $S \subset {}^*\mathcal{R}^3$ , with  $S$   $*$ -finite and  $S$  infinite, there exists  $\kappa_S \in {}^*\mathcal{N}$ ,  $\kappa_S$  odd and infinite, and  $\bar{x}_S \in {}^*\mathcal{R}^3$ , with  $|\bar{x}_S| < \frac{1}{\kappa_S}$ , such that if;

$$V_{\kappa_S, \bar{x}_S} = \{\bar{x}_S + \bar{w} : w_i = \frac{j}{\sqrt{\kappa_S}}, -\frac{(\kappa_S-1)}{2} \leq j \leq \frac{(\kappa_S-1)}{2}, 1 \leq i \leq 3\}$$

then for all  $\bar{v} \in V_{\kappa_S, \bar{x}_S}$ , and  $\{\bar{x}_1, \bar{x}_2\} \subset S$ , with  $\bar{x}_1 \neq \bar{x}_2$ , we have that  $|\bar{v} - \bar{x}_1| \neq |\bar{v} - \bar{x}_2|$ , (\*).

*Proof.* For  $S = \{s_1, \dots, s_n\}$  finite, then  $\{\bar{w} \in \mathcal{R}^3 : |\bar{w} - s_i| = |\bar{w} - s_j|, \text{ for some } s_i \neq s_j\}$  forms a finite union of hyperplanes  $H_j \subset \mathcal{R}^3$ ,  $1 \leq j \leq r$ . Choose  $m \in \mathcal{N}$ , with  $m$  even,  $m > n$ , and let;

$$W_m = \{\bar{z} : z_i = \frac{j}{\sqrt{m}}, -\frac{(m-1)}{2} \leq j \leq \frac{(m-1)}{2}, 1 \leq i \leq 3\}$$

We have that  $W_m \cap \bigcup_{1 \leq j \leq r} H_j$  is a finite set of points  $\{p_1, \dots, p_s\}$ . Let  $t = \min\{d(\bar{z}, \bigcup_{1 \leq j \leq r} H_j) : \bar{z} \in W_m \setminus \{p_1, \dots, p_s\}\}$ , and choose  $\bar{y} \in B(\bar{0}, \lambda)$ , with  $\lambda < \min(t, \frac{1}{m})$ , such that  $\bar{y} - p_i \notin \bigcup_{1 \leq j \leq r} H_j$ , for  $1 \leq j \leq s$ . Then, it is clear that  $(\bar{y} + W_m) \cap \bigcup_{1 \leq j \leq r} H_j = \emptyset$ . By construction,  $\{m, \bar{y}, \bar{y} + W_m\}$  has the property (\*), relative to  $S$ . We can formulate the above property (\*) for finite sets in the language  $L_{\mathcal{R}}$ , using second order logic, developed in [8]. By transfer to  ${}^*\mathcal{R}$ , for any given  $*$ -finite set  $S$ , we can find  $\{\kappa_S, \bar{x}_S, V_{\kappa_S, \bar{x}_S}\}$ , with the same property (\*) relative to  $S$ . As  $m > n$ , with  $m \in \mathcal{N}$  odd, where  $n = \text{Card}(S)$ , in the property (\*) for finite sets,  $\kappa_S > {}^*\text{Card}S$ , with  $\kappa_S \in {}^*\mathcal{N}$  odd and infinite.  $\square$

**Lemma 1.17.** *Let  $S = \bigcup_{t \in \mathcal{T}_\nu} R_t$ , then there exists  $\kappa \in {}^*\mathcal{N}$ ,  $\kappa$  odd and infinite, with  $\bar{x}_\kappa \in {}^*\mathcal{R}^3$ ,  $|\bar{x}_\kappa| < \frac{1}{\kappa} \simeq 0$ , such that if, as in Definitions 1.3 and 1.13;*

$$\mathcal{R}_\kappa = \bigcup_{-\frac{(\kappa-1)}{2} \leq i \leq \frac{(\kappa-1)}{2}} \left[ \frac{i}{\sqrt{\kappa}}, \frac{i+1}{\sqrt{\kappa}} \right)$$

$$R_\kappa = \{ \bar{x} \in \mathcal{R}_\kappa^3 : [x_i \sqrt{\kappa}] = x_i \sqrt{\kappa}, \text{ for } 1 \leq i \leq 3 \}$$

*then every  $\bar{y} \in \bar{x}_\kappa + R_\kappa$  has a nearest neighbour in  $R_t$ , that is there exists  $\bar{w}_t \in R_t$ , with the property that;*

$$|\bar{y} - \bar{w}_t| < |\bar{y} - \bar{z}_t| \text{ for every } \bar{z}_t \in R_t.$$

*Proof.* By Lemma 1.15,  $S = \bigcup_{t \in \mathcal{T}_\nu} R_t$  is  $*$ -finite, and by Lemma 1.16, we can find  $\kappa \in {}^*\mathcal{N}$ ,  $\kappa$  odd and infinite, with  $\bar{x}_\kappa \in {}^*\mathcal{R}^3$ ,  $|\bar{x}_\kappa| < \frac{1}{\kappa}$ , and the property that there do not exist  $\{\bar{x}_1, \bar{x}_2\} \subset S$ , with  $|\bar{y} - \bar{x}_1| = |\bar{y} - \bar{x}_2|$ , for a given  $\bar{y} \in \bar{x}_\kappa + R_\kappa$ , (\*). In particular, the property (\*) holds for  $R_t$ , when  $t \in \mathcal{T}_\nu$ . Let  $D_{t, \bar{y}} = \{|\bar{y} - \bar{z}_t| : \bar{z}_t \in R_t\}$ , then  $D_t$  is  $*$ -finite, hence has a least element  $\mu$ . Choosing  $\bar{w}_t \in R_t$ , with  $|\bar{y} - \bar{w}_t| = \mu$ , it is clear, by the property (\*), that  $\bar{w}_t$  has the required property.  $\square$

**Definition 1.18.** *Let notation be as in the previous lemma. Let  $\bar{w} \in R_t$ , then we define the  $*$ -finite sets  $I_{\bar{w}, t}$  and  $\bar{p}_{\bar{w}, t}$ , by;*

$$I_{\bar{w}, t} = \{ \bar{x} \in R : s_{\bar{x}}(t) = \bar{w} \}$$

$$\bar{p}_{\bar{w}, t} \in I_{\bar{w}, t}, \bar{p}_{\bar{w}, t} \ll \bar{x}, \bar{x} \in I_{\bar{w}, t}$$

$$J_{\bar{w}, t} = \{ \bar{z} \in \bar{x}_\kappa + R_\kappa, \bar{z} \text{ is a nearest neighbour to } \bar{w} \}$$

*where  $\ll$  is the lexicographic total well order on  $I_{\bar{w}, t}$ .*

*For  $\bar{w} \in R_t$ , we define  $\lambda_{\bar{w}, t} = {}^*\text{Card}(J_{\bar{w}, t})$ . We define  $\rho_\kappa$  on  $\bar{x}_\kappa + \mathcal{R}_\kappa^3 \times \mathcal{T}_\nu$  by;*

$$\rho_\kappa(\bar{x}_\kappa + \bar{x}, \frac{j}{\nu}) = \frac{\kappa^{\frac{3}{2}} \rho_\eta(\bar{p}_{\bar{w}_{\bar{x}}, t}, 0)}{\eta^{\frac{3}{2}} \lambda_{\bar{w}, t}}$$

*where  $\bar{w}_{\bar{x}}$  is the nearest neighbour to  $\bar{x}_\kappa + \bar{x}$  in  $R_t$ , and  $0 \leq j \leq \nu^2 - 1$ ,  $[x_i \sqrt{\kappa}] = x_i \sqrt{\kappa}$ , for  $1 \leq i \leq 3$ .*

$$\rho_\kappa(\bar{x}_\kappa + \bar{x}, t) = \rho_\kappa(\bar{x}_\kappa + \bar{y}, \frac{[t\nu]}{\nu})$$

$$\text{where } y_i = \frac{[x_i\sqrt{\kappa}]}{\sqrt{\eta}}, \bar{x} \in \bar{R}_\kappa^3.$$

**Lemma 1.19.** *For  $\bar{y} \in \bar{R}_\kappa^3$  and  $t \in \mathcal{T}_\nu$  finite, with  $\nu \geq \eta$ , then  $|I_{\bar{w}_\bar{x}, t}| = 1$ , where  $x_i = \frac{[y_i\sqrt{\kappa}]}{\sqrt{\kappa}}$ , for  $1 \leq i \leq 3$ , and  $\rho_\kappa(\bar{x}_\kappa + \bar{y}, t)$  is finite and not infinitesimal.*

*Proof.* Using Definition 1.18, if  $\bar{z} \in I_{\bar{w}_\bar{x}, t}$ , then  $\bar{z} \in R$  and  $s_{\bar{z}}(t) = \bar{w}_\bar{x}$ . If  $\bar{z}$  is infinite, then, by Lemma 1.14,  $s_{\bar{z}}(t)$  is infinite, so that  $\bar{w}_\bar{x}$  is infinite, (\*). By Lemma 1.11, as  $\bar{x}_\kappa + \bar{x}$  is finite, there exists  $\bar{x}_1 \in R$ , with  $s_{\bar{x}_1}(t) \simeq \bar{x}_\kappa + \bar{x}$ , in particular  $s_{\bar{x}_1}(t) \in R_t$ , and the nearest neighbour  $\bar{w}_\bar{x}$  to  $\bar{x}_\kappa + \bar{x}$  must be finite, contradicting (\*). It follows that  $\bar{z}$  is finite. Now suppose that  $I_{\bar{w}_\bar{x}, t}$  contains two distinct finite elements  $\{\bar{z}_1, \bar{z}_2\} \subset R$ , then, using Lemma 1.8,  $s_{\bar{z}_1}(t) \neq s_{\bar{z}_2}(t)$  and define distinct elements of  $R_t$ , so that, by the definition of  $I_{\bar{w}_\bar{x}, t}$ ,  $\bar{w}_\bar{x} = s_{\bar{z}_1}(t) \neq s_{\bar{z}_2}(t) = \bar{w}_\bar{x}$ , a contradiction. It follows that  $|I_{\bar{w}_\bar{x}, t}| = 1$ , and, using the definition of  $R_t$ , and the first argument again, that  $\bar{p}_{\bar{w}_\bar{x}, t}$  is finite, (‡). It follows, by Definitions 1.3 and 1.18, that  $\rho_\eta(\bar{p}_{\bar{w}_\bar{x}, t}, 0)$  is finite. Let  $\bar{w} \in R_t$ , and let  $C = (3 + 2F)e^{\circ t_0} + 6$ . We claim that if  $\bar{x} \in \bar{x}_\kappa + R_\kappa$ , and  $\bar{x}$  is a nearest neighbour to  $\bar{w}$ , then  $\bar{x} \in {}^*B(\bar{w}, \frac{C}{\sqrt{\eta}})$ , (†). Suppose not, then there exists  $\bar{x} \in \bar{x}_\kappa + R_\kappa$ , such that  $|\bar{x} - \bar{w}_\bar{x}| > \frac{C}{\sqrt{\eta}}$ , where  $\bar{w}_\bar{x}$  is the nearest neighbour to  $\bar{x}$ . By Lemma 1.11, using the fact that  $\nu \geq \sqrt{\eta}$ , we can find  $\bar{x}' \in R_t$  with;

$$\begin{aligned} |\bar{x} - \bar{x}'| &\leq (e^{\circ t_0} + 2)\left(\frac{3}{\sqrt{\eta}} + \frac{F}{\nu}\right) + \frac{Fe^{\circ t_0}}{\nu} \\ &\leq \frac{1}{\sqrt{\eta}}((3 + 2F)e^{\circ t_0} + 6) \end{aligned}$$

which implies that  $|\bar{x} - \bar{x}'| < |\bar{x} - \bar{w}_\bar{x}|$ , contradicting the fact that  $\bar{w}_\bar{x}$  is the nearest neighbour to  $\bar{x}$ , proving (†). Using Lemmas 1.6 and 1.12, we can find  $C_{t,r} \in \mathcal{R}_{>0}$ , such that if  $\bar{w}' \in R_t$ , and  $|\bar{w} - \bar{w}'| < 1$ , then  $|\bar{w} - \bar{w}'| \geq \frac{C_{t,r}}{\sqrt{\eta}}$ . It follows that if  $\bar{x} \in \bar{x}_\kappa + R_\kappa$ , and  $\bar{x} \in {}^*B(\bar{w}, \frac{C_{t,r}}{2\sqrt{\eta}})$ , then  $\bar{x}$  is a nearest neighbour to  $\bar{w}$ , (††). Suppose that  $C_{t,r} \geq C$ , then;

$$\begin{aligned} &{}^*Card(\{\bar{x} \in \bar{x}_\kappa + R_\kappa : \bar{w} \text{ is a nearest neighbour to } \bar{x}\}) \\ &= {}^*Card((\bar{x}_\kappa + R_\kappa) \cap {}^*B(\bar{w}, \frac{C}{\sqrt{\eta}})) \quad (\dagger\dagger\dagger) \end{aligned}$$

If  $C_{t,r} < C$ , then;

$$\begin{aligned}
& *Card(\bar{x}_\kappa + R_\kappa \cap *B(\bar{w}, \frac{C_{t,r}}{\sqrt{\eta}})) \\
& \leq *Card(\{\bar{x} \in \bar{x}_\kappa + R_\kappa : \bar{w} \text{ is a nearest neighbour to } \bar{x}\}) \\
& \leq *Card((\bar{x}_\kappa + R_\kappa) \cap *B(\bar{w}, \frac{C}{\sqrt{\eta}})) \quad (\dagger\dagger\dagger)
\end{aligned}$$

Observe that  $*D(\bar{w}, \frac{r}{\sqrt{2}}) \subset *B(\bar{w}, r) \subset *D(\bar{w}, r)$ , where  $*D(\bar{w}, r)$  is the  $*$ -cube of diameter  $2r$  centred at  $\bar{w}$ .

We have, with the usual definition of  $[, ]$  the binomial theorem, and  $r^2\kappa$  infinite, that;

$$\begin{aligned}
& *Card((\bar{x}_\kappa + R_\kappa) \cap *D(\bar{w}, r)) \\
& = *Card((\bar{v} + R_\kappa) \cap *D(\bar{0}, r)) \\
& = (2[r\sqrt{\kappa}] + 1 - \delta_1)(2[r\sqrt{\kappa}] + 1 - \delta_2)(2[r\sqrt{\kappa}] + 1 - \delta_3) \\
& = (2r\sqrt{\kappa} + \gamma_1)(2r\sqrt{\kappa} + \gamma_2)(2r\sqrt{\kappa} + \gamma_3) \\
& = 8r^3\kappa^{\frac{3}{2}} + \theta_r
\end{aligned}$$

where  $v_i = x_{i,\kappa} - \frac{[x_{i,\kappa}\sqrt{\kappa}]}{\sqrt{\kappa}}$ , for  $1 \leq i \leq 3$ , so that  $0 \leq v_i < \frac{1}{\sqrt{\kappa}}$ , and  $\delta_i \in \{0, 1\}$ ,  $1 \leq i \leq 3$ ,  $0 \leq \gamma_i < 2$ ,  $|\theta_r| \leq 50r^2\kappa$ . It follows, from  $(\dagger\dagger\dagger)$ , that;

$$\frac{8C^3\kappa^{\frac{3}{2}}}{(2\eta)^{\frac{3}{2}}} + \theta \frac{C}{\sqrt{2\eta}} \leq \lambda_{\bar{w},t} \leq 8 \frac{C^3\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}} + \theta \frac{C}{\sqrt{\eta}} \quad (*)$$

and, from  $(\dagger\dagger\dagger)$ , that;

$$\frac{8C_{t,r}^3\kappa^{\frac{3}{2}}}{(2\eta)^{\frac{3}{2}}} + \theta \frac{C_{t,r}}{\sqrt{2\eta}} \leq \lambda_{\bar{w},t} \leq 8 \frac{C^3\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}} + \theta \frac{C}{\sqrt{\eta}} \quad (**)$$

It follows from  $(*)$ , that;

$$\begin{aligned}
\frac{1}{8C^3 + \frac{\eta^{\frac{3}{2}}}{\kappa^{\frac{3}{2}}}\theta \frac{C}{\sqrt{\eta}}} &= \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}(8 \frac{C^3\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}} + \theta \frac{C}{\sqrt{\eta}})} \leq \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}\lambda_{\bar{w},t}} \leq \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}(8 \frac{C^3\kappa^{\frac{3}{2}}}{(2\eta)^{\frac{3}{2}}} + \theta \frac{C}{\sqrt{2\eta}})} = \frac{1}{8 \frac{C^3}{2^{\frac{3}{2}}} + \frac{\eta^{\frac{3}{2}}}{\kappa^{\frac{3}{2}}}\theta \frac{C}{\sqrt{2\eta}}} \\
\frac{1}{8C^3 + \alpha} &\leq \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}\lambda_{\bar{w},t}} \leq \frac{1}{8 \frac{C^3}{2^{\frac{3}{2}}} + \beta}
\end{aligned}$$

$$\text{where } |\alpha| = \left| \frac{\eta^{\frac{3}{2}}}{\kappa^{\frac{3}{2}}}\theta \frac{C}{\sqrt{\eta}} \right| \leq 50C^2 \frac{\eta^{\frac{3}{2}}}{\kappa^{\frac{3}{2}}} \frac{\kappa}{\eta} = 50C^2 \frac{\eta^{\frac{1}{2}}}{\kappa^{\frac{1}{2}}} \simeq 0$$

$$|\beta| = \left| \frac{\eta^{\frac{3}{2}} \theta_{\frac{C}{\sqrt{2\eta}}}}{\kappa^{\frac{3}{2}}} \right| \leq 25C^2 \frac{\eta^{\frac{3}{2}} \kappa}{\kappa^{\frac{3}{2}} \eta} = 25C^2 \frac{\eta^{\frac{1}{2}}}{\kappa^{\frac{1}{2}}} \simeq 0$$

and, from (\*\*), that;

$$\begin{aligned} \frac{1}{8C^3 + \frac{\eta^{\frac{3}{2}} \theta_{\frac{C}{\sqrt{2\eta}}}}{\kappa^{\frac{3}{2}}}} &= \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}} (8\frac{C^3 \kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}}} + \theta_{\frac{C}{\sqrt{2\eta}}})} \leq \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}} \lambda_{\bar{w},t}} \leq \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}} \left( \frac{8C_{t,r}^3 \kappa^{\frac{3}{2}}}{(2\eta)^{\frac{3}{2}}} + \theta_{\frac{C_{t,r}}{\sqrt{2\eta}}} \right)} = \frac{1}{8\frac{C_{t,r}^3}{2^{\frac{3}{2}}} + \frac{\eta^{\frac{3}{2}} \theta_{\frac{C_{t,r}}{\sqrt{2\eta}}}}{\kappa^{\frac{3}{2}}}} \\ \frac{1}{8C^3 + \alpha} &\leq \frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}} \lambda_{\bar{w},t}} \leq \frac{1}{8\frac{C_{t,r}^3}{2^{\frac{3}{2}}} + \gamma} \end{aligned}$$

$$\text{where } |\gamma| = \left| \frac{\eta^{\frac{3}{2}} \theta_{\frac{C_{t,r}}{\sqrt{2\eta}}}}{\kappa^{\frac{3}{2}}} \right| \leq 25C_{t,r}^2 \frac{\eta^{\frac{3}{2}} \kappa}{\kappa^{\frac{3}{2}} \eta} = 25C_{t,r}^2 \frac{\eta^{\frac{1}{2}}}{\kappa^{\frac{1}{2}}} \simeq 0$$

As  $\{C, C_{t,r}\} \subset \mathcal{R}_{>0}$ , and  $\{\alpha, \beta, \gamma\}$  are infinitesimal, in both cases, we have that  $\frac{\kappa^{\frac{3}{2}}}{\eta^{\frac{3}{2}} \lambda_{\bar{w},t}}$  is finite and not infinitesimal. It follows, using (#), Definition 1.18 and the fact that  $1 < \rho_{\eta,0} \leq H$ , for some  $H \in \mathcal{R}_{>0}$ , that  $\rho_{\kappa}(\bar{x}_{\kappa} + \bar{y}, t)$  is finite and not infinitesimal as well.  $\square$

**Definition 1.20.** We say that an internal set  $V \subset \bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3$  is finitely bounded, if there exists a constant  $r \in \mathcal{R}_{>0}$  such that  $V \subset \bar{x}_{\kappa} + {}^*B(\bar{0}, r)$ . We say that  $f$  is measurable on  $\bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3$  if;

$$f(\bar{x}_{\kappa} + \bar{z}) = f(\bar{x}_{\kappa} + \bar{y})$$

for  $\bar{z} \in \mathcal{R}_{\kappa}^3$  and  $y_i = \frac{[z_i \sqrt{\kappa}]}{\sqrt{\kappa}}$ , for  $1 \leq i \leq 3$ . We let  $\mu_{\kappa}$  be the internal counting measure on  $\mathcal{R}_{\kappa}^3$ . By Definition 1.18, we have that  $\rho_{\kappa,t}$  is measurable on  $\bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3$  for  $t \in \mathcal{T}_{\nu}$ . For  $f$  measurable on  $\bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3$ , and  $V \subset \bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3$  finitely bounded, we define the internal integral;

$$\int_V f d\mu_{\kappa} = \frac{1}{\kappa^{\frac{3}{2}}} {}^* \sum_{\bar{x} \in (\bar{x}_{\kappa} + R_{\kappa}) \cap V} f(\bar{x})$$

We let  $\mathfrak{C}_{\kappa}$  be the internal  $*$ -sigma algebra generated by the internal sets;

$$\{V_{\bar{x}} : \bar{x} \in \bar{x}_{\kappa} + R_{\kappa}\}$$

where  $V_{\bar{x}} = \{\bar{y} \in \bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3 : 0 \leq y_i - x_i < \frac{1}{\sqrt{\kappa}}, 1 \leq i \leq 3\}$ , and let  $L(\mathfrak{C}_{\kappa})$  be the  $\sigma$ -algebra generated by  $\mathfrak{C}_{\kappa}$ . We let  $L(\mu_{\kappa})$  be the Loeb measure on  $L(\mathfrak{C}_{\kappa})$ , associated to  $\mu_{\kappa}$ , see [5] or [4]. For  $V \in L(\mathfrak{C}_{\kappa})$ , we say that  $V$  is finitely bounded if there exists a constant  $r \in \mathcal{R}_{>0}$  such that  $V \subset \bar{x}_{\kappa} + {}^*B(\bar{0}, r)$ . For  $f$  measurable and finite on  $\bar{x}_{\kappa} + \mathcal{R}_{\kappa}^3$ , with

standard part  $\circ f$ , we denote the Loeb integral by;

$$\int_V \circ f dL(\mu_\kappa)$$

Note that the Loeb measure  $L(\mu_\kappa)$  specialises to Lebesgue measure  $d\mu$  on  $\mathcal{R}^3$ , see [1]. This implies that, for  $V \in L(\mathfrak{C}_\kappa)$ , finitely bounded, with  $V = st^{-1}(W)$ , for  $W$  Lebesgue measurable in  $\mathcal{R}^3$  and  $g$   $S$ -continuous;

$$\int_V g(\bar{x}) dL(\mu_\kappa)(\bar{x}) = \int_W g(\circ \bar{x}) d\mu(\circ \bar{x})$$

**Lemma 1.21.** *The  $*$ -intersection  $W$  of an internal collection  $\{V_i : i \in I\}$ ,  $V_i \subset \bar{x}_\kappa + \mathcal{R}_\kappa^3$ , indexed by internal  $I$ , of  $*$ -closed sets, is internal and  $*$ -closed. If the collection contains a finitely bounded set, then  $W$  is finitely bounded. If  $f$  is measurable on  $\bar{x}_\kappa + \mathcal{R}_\kappa^3$  and finite on  $V$ , with  $V \subset \bar{x}_\kappa + \mathcal{R}_\kappa^3$  finitely bounded and internal, then  $\int_V f d\mu_\kappa$  is finite. Similarly, if  $f$  is finite on  $V \in L(\mathfrak{C}_\kappa)$ , with  $|f| \leq D$ ,  $D \in \mathcal{R}_{>0}$ , then  $\int_V \circ f dL(\mu_\kappa)$  is finite.*

*Proof.* The claim that  $W$  is  $*$ -closed and internal, follows by transfer from the fact that a finite intersection of closed sets  $\{V_1, \dots, V_n\}$ ,  $V_i \subset \mathcal{R}^3$ ,  $1 \leq i \leq n$  is a closed set. If some  $V_1$  is finitely bounded, the fact that  $W$  is finitely bounded follows from the fact that  $W \subset \bar{x}_\kappa + V_1$ , with  $V_1 \subset {}^*B(\bar{0}, r)$ , and  $r \in \mathcal{R}_{>0}$ . For the final claim, choose  ${}^*B(\bar{0}, r)$ , with  $V \subset \bar{x}_\kappa + {}^*B(\bar{0}, r)$ , and let  $S(r) = {}^*[-r, r]^3 \supset {}^*B(\bar{0}, r)$ . There exists a constant  $C \in \mathcal{R}_{>0}$ , with  $|f|_V \leq C$ , otherwise, by compactness and the fact that  $V$  is internal,  $f$  would not be finite on  $V$ . Then;

$$\begin{aligned} |\int_V f d\mu_\kappa| &\leq \int_V |f| d\mu_\kappa \leq C \int_V d\mu_\kappa \\ &= \frac{C}{\kappa^{\frac{3}{2}}} {}^*Card(\{\bar{x} \in (\bar{x}_\kappa + R_\kappa) \cap V\}) \\ &\leq \frac{C}{\kappa^{\frac{3}{2}}} {}^*Card(\{\bar{x} \in (\bar{x}_\kappa + R_\kappa) \cap \bar{x}_\kappa + S(r)\}) \\ &= \frac{C}{\kappa^{\frac{3}{2}}} (2[r\sqrt{\kappa}] + 1)^3 \\ &\leq \frac{C}{\kappa^{\frac{3}{2}}} (2r\sqrt{\kappa} + 3)^3 \\ &\leq \frac{C}{\kappa^{\frac{3}{2}}} ((2r + 1)\sqrt{\kappa})^3 \\ &= C(2r + 1)^3 \end{aligned}$$

which is finite, as required. The second claim, using the definition of finitely bounded, the remark in Definition 1.20 and the fact that  $*B(0, r) \subset st^{-1}(B(0, r + 1))$ , for  $r \in \mathcal{R}_{>0}$ , follows from;

$$\begin{aligned}
& \left| \int_V \circ f dL(\mu_\kappa) \right| \\
& \leq \int_V |\circ f| dL(\mu_\kappa) \\
& \leq D \int_V dL(\mu_\kappa) \\
& \leq D \int_{\bar{x}_\kappa + st^{-1}(B(\bar{0}, r+1))} dL(\mu_\kappa) \\
& = D \int_{B(\bar{0}, r+1)} d\mu = \frac{4D\pi(r+1)^3}{3}
\end{aligned}$$

where  $d\mu$  is Lebesgue measure. □

**Definition 1.22.** *Transport of Open Balls*

Let  $V \subset \bar{x}_\kappa + \mathcal{R}_\kappa^3$ ,  $V = \bar{x}_\kappa + W$ ,  $W = st^{-1}(B(\bar{0}, r))$ , for  $r \in \mathcal{R}_{>0}$ . Let  $\{t_1, t_2\} \subset \mathcal{T}_\nu$ , with  $t_1 \leq t_2$ . Let  $S_V = R_{t_1} \cap V$  and let;

$$S_{V, t_1, t_2} = \{s_{\bar{x}}(t_2) : \bar{x} \in R, s_{\bar{x}}(t_1) \in S_V\}$$

We define the transport  $V_{t_1, t_2}$  of  $V$  to be the intersection of all  $V_i \subset \bar{x}_\kappa + \mathcal{R}_\kappa^3$ , where  $V_i \supset S_{V, t_1, t_2}$ ,  $V_i = \bar{x}_\kappa + W_i$ ,  $W_i = st^{-1}(Z_i)$ , for  $Z_i$  open in  $\mathcal{R}^3$ . Note that, by the remark in Definition 1.20,  $V_i \in L(\mathfrak{C}_\kappa)$ , and, as  $L(\mathfrak{C}_\kappa)$  is a  $\sigma$ -algebra,  $V_{t_1, t_2} \in L(\mathfrak{C}_\kappa)$  as well.

**Lemma 1.23.** *For any  $\mu_\kappa$  measurable  $A \subset \mathcal{R}_\kappa^3$ , which is finitely bounded, we have that;*

$$\int_A \rho_\kappa d\mu_\kappa \simeq \int_A \circ \rho_\kappa dL(\mu_\kappa)$$

*Proof.* We claim, that  $\rho_\kappa$  is  $S$ -integrable on  $A$ , for the  $*$ - $\sigma$  algebra  $\mathcal{C}_{A, \kappa}$ , generated by intersecting  $\mathcal{C}_\kappa$  with  $A$ . As  $\mu_\kappa(A)$  is finite, by Definition 3.17 of [5], we have to check that;

- (i).  $\int_A |\rho_\kappa| d\mu_\kappa$  is finite.
- (ii). For  $B \in \mathcal{C}_{A, \kappa}$ , with  $\mu_\kappa(B) \simeq 0$ ;

$$\int_B |\rho_\kappa| d\mu_\kappa \simeq 0$$

For (i), we have, by Lemma 1.19, that  $\rho_\kappa|_A$  is finite. It follows, using overflow to  $\{n \in {}^*\mathcal{N} : \exists x \in A, \rho_\kappa(x) > n\}$ , that there exists a constant  $D \in \mathcal{R}_{>0}$ , with  $|\rho_\kappa|_A| \leq D$ . Then;

$$\begin{aligned} \int_A |\rho_\kappa| d\mu_\kappa &\leq C \int_A d\mu_\kappa \\ &= C\mu_\kappa(A) \end{aligned}$$

which is finite, and;

$$\begin{aligned} \int_B |\rho_\kappa| d\mu_\kappa &\leq C \int_B d\mu_\kappa \\ &= C\mu_\kappa(B) \simeq 0 \end{aligned}$$

It then follows, by Theorem 3.20 of [5], that;

$${}^\circ \int_A \rho_\kappa d\mu_\kappa = \int_A {}^\circ \rho_\kappa dL(\mu_\kappa)$$

as required. □

**Lemma 1.24.** *Let notation be as in Definition 1.22, then the transport  $V_{t_1, t_2}$  of  $V$ , is finitely bounded. Moreover,  $\int_V {}^\circ \rho_{\kappa, t_1} dL(\mu_\kappa)$  is finite and;*

$$\int_V {}^\circ \rho_{\kappa, t_1} dL(\mu_\kappa) = \int_{V_{t_1, t_2}} {}^\circ \rho_{\kappa, t_2} dL(\mu_\kappa)$$

*Proof.* For the first part, choose  $\epsilon > 0$ , with  $V \subset {}^*B(\bar{0}, r + \epsilon)$ . Using the proof of Lemma 1.7, we have that, for  $s_{\bar{x}}$ , with  $s_{\bar{x}}(t_1) \in S_V$ ;

$$\begin{aligned} |s_{\bar{x}}(t_2)| &\leq |s_{\bar{x}}(t_1)| + G(t_2 - t_1) \\ &\leq r + \epsilon + G(t_2 - t_1) = E \end{aligned}$$

We then have that, as  $\bar{x}_\kappa$  is infinitesimal, that  $\bar{x}_\kappa + st^{-1}(B(\bar{0}, E + 1)) \supset S_{V_{t_1, t_2}}$ , so that the transport  $V_{t_1, t_2} \subset \bar{x}_\kappa + st^{-1}(B(\bar{0}, E + 1)) \subset \bar{x}_\kappa + {}^*B(\bar{0}, {}^\circ(E + 2))$ . As  ${}^\circ(E + 2) \in \mathcal{R}_{>0}$ , we obtain the result by Definition 1.20. For the second part, keeping the same notation as above, let  $W = (\bar{x}_\kappa + R_\kappa) \cap {}^*B(\bar{0}, r + \epsilon)$ , then  $W \supset V$  is internal and, by Lemma 1.19,  $\rho_{\kappa, t_1}$  is finite on  $W$ . By the argument in Lemma 1.21, we can find a constant  $C \in \mathcal{R}_{>0}$ , with  $|\rho_{\kappa, t_1}|_W| < C$ . In particular,

$|\rho_{\kappa,t_1}|_V < C$ , so that  $|\circ\rho_{\kappa,t_1}|_V < C$ . It follows that;

$$\begin{aligned} & \left| \int_V \circ\rho_{\kappa,t_1} dL(\mu_\kappa) \right| \leq \int_V |\circ\rho_{\kappa,t_1}| dL(\mu_\kappa) \\ & \leq C \int_V dL(\mu_\kappa) \\ & = C \int_{\bar{x}_\kappa + st^{-1}(B(\bar{0},r))} dL(\mu_\kappa) \\ & = C \int_{B(\bar{0},r)} d\mu = \frac{4C\pi r^3}{3} \end{aligned}$$

as required. We claim that if  $\bar{y} \in R_{t_2} \cap V_{t_1,t_2}$ , with  $s_{\bar{x}}(t_2) = \bar{y}$ , then, for all  $\epsilon \in \mathcal{R}_{>0}$ , there exists  $\bar{z} \in R_{t_1} \cap (\bar{x}_\kappa + {}^*B(\bar{0}, r + \epsilon))$  with  $s_{\bar{x}'}(t_1) = \bar{z}$  and  $s_{\bar{x}'}(t_2) \simeq \bar{y}$ , ( $\dagger$ ). Suppose not, then for every  $\bar{w} \in R_{t_1} \cap {}^*B(\bar{0}, r + \epsilon)$ , we have that  $|s_{\bar{p}_{\bar{w}}}(t_2) - \bar{y}| \geq \delta$ , for all infinitesimals  $\delta$ . Let;

$$Z = \{\bar{p}_{\bar{w}} : \bar{w} \in R_{t_1} \cap {}^*B(\bar{0}, r + \epsilon)\}$$

Then, as  $R_{t_1} \cap {}^*B(\bar{0}, r + \epsilon)$  is \*-finite, so is  $Z$ . Define  $g$  on  $Z$ , by  $g(\bar{p}_{\bar{w}}) = |\bar{y} - s_{\bar{p}_{\bar{w}}}(t_2)|$ . Then, using the fact that  $|g| \geq \delta$ , for all infinitesimals  $\delta$ , by the usual underflow argument, we can find  $\gamma \in \mathcal{R}_{>0}$ , with  $|g| \geq \gamma$ . Let  $K = (\bar{x}_\kappa + st^{-1}(A^c))$  where  $A$  is the closed ball;

$$A = \{\bar{z} \in \mathcal{R}^3 : |\bar{z} - \circ\bar{y}| \leq \frac{\gamma}{2}\}$$

Then,  $A^c$  is an open set such that  $\bar{y} \notin \bar{x}_\kappa + st^{-1}(A^c)$ . As  $\bar{x}_\kappa + st^{-1}(A^c) \supset S_{V,t_1,t_2}$ ,  $V_{t_1,t_2} \subset \bar{x}_\kappa + st^{-1}(A^c) \cap V_{t_1,t_2}$ , so that  $\bar{y} \in \bar{x}_\kappa + st^{-1}(A^c)$ , a contradiction. For the next claim, let  $r_n = r(1 - \frac{1}{n})$ , then, using the notation in Definition 1.25, we claim that  $W_{t_1,r_n} \subset V$  and  $W_{t_1,t_2,r_n} \subset V_{t_1,t_2}$ , ( $\ddagger$ ). The first part of ( $\ddagger$ ) is clear. If  $W_i = st^{-1}(Z_i)$ , with  $\bar{x}_\kappa + W_i \supset S_{V,t_1,t_2}$ , we have that  $\bar{x}_\kappa + W_i \supset U_{t_1,t_2,r_n}$ , and, if  $\bar{x} \in W_{t_1,t_2,r_n}$ , we can find  $\bar{w} \in U_{t_1,t_2,r_n}$ , with  $\bar{x} \simeq \bar{w}$ . Using the fact that  $Z_i$  is open, choose  $\epsilon \in \mathcal{R}_{>0}$ , such that  $\bar{x}_\kappa + {}^*B(\bar{w}, \epsilon)$ , with  $\bar{x}_\kappa + {}^*B(\bar{w}, \epsilon) \subset \bar{x}_\kappa + W_i$ , so that  $\bar{x} \in \bar{x}_\kappa + W_i$ . It follows that  $\bar{x} \in V_{t_1,t_2}$ , proving ( $\ddagger$ ). Using the result of Lemma 1.26, we have that;

$$\int_{W_{t_1,r_n}} \rho_\kappa d\mu_\kappa = \int_{W_{t_1,t_2,r_n}} \rho_\kappa d\mu_\kappa \quad (\#\#)$$

so that, using ( $\#\#$ ), Lemma 1.23, for all  $r_n$ , with  $n \in \mathcal{N}$ ;

$$\int_V \circ\rho_\kappa dL(\mu_\kappa) = \int_{W_{t_1,r_n}} \circ\rho_\kappa dL(\mu_\kappa) + \int_{(V \setminus W_{t_1,r_n})} \circ\rho_\kappa dL(\mu_\kappa)$$

$$\begin{aligned}
&\simeq \int_{W_{t_1, r_n}} \rho_\kappa d\mu_\kappa + \int_{(V \setminus W_{t_1, r_n})} \circ \rho_\kappa dL(\mu_\kappa) \\
&\int_{V_{t_1, t_2}} \circ \rho_\kappa dL(\mu_\kappa) = \int_{W_{t_1, t_2, r_n}} \circ \rho_\kappa dL(\mu_\kappa) + \int_{(V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})} \circ \rho_\kappa dL(\mu_\kappa) \\
&\simeq \int_{W_{t_1, t_2, r_n}} \rho_\kappa d\mu_\kappa + \int_{(V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})} \circ \rho_\kappa dL(\mu_\kappa) \\
&= \int_{W_{t_1, r_n}} \rho_\kappa d\mu_\kappa + \int_{(V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})} \circ \rho_\kappa dL(\mu_\kappa) \\
&\simeq \int_V \circ \rho_\kappa dL(\mu_\kappa) - \int_{(V \setminus W_{t_1, r_n})} \circ \rho_\kappa dL(\mu_\kappa) + \int_{(V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})} \circ \rho_\kappa dL(\mu_\kappa) \\
&(\#\#\#)
\end{aligned}$$

As  $\max(|\rho_\kappa|_V, |\rho_\kappa|_{V_{t_1, t_2}}) \leq H$  for some  $H \in \mathcal{R}_{>0}$ , we have that;

$$\begin{aligned}
&|\int_{(V \setminus W_{t_1, r_n})} \circ \rho_\kappa dL(\mu_\kappa)| \leq (H+1)L(\mu_\kappa)((V \setminus W_{t_1, r_n})) \\
&|\int_{(V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})} \circ \rho_\kappa dL(\mu_\kappa)| \leq (H+1)L(\mu_\kappa)((V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})) (***)
\end{aligned}$$

By the proof above, we have, for any  $\delta \in \mathcal{R}_{>0}$ , that if  $\bar{x} \in V_{t_1, t_2}$ , we can find  $\bar{y} \in \bar{x}_\kappa + {}^*B(\bar{0}, r + \delta)$  with  $s_{\bar{p}_{\bar{y}, t_1}}(t_2) \simeq \bar{x}$ . Suppose that  $\bar{x} \notin W_{t_1, t_2, r_n}$ , then if  $\bar{y} \in \bar{x}_\kappa + {}^*B(\bar{0}, r_{n-1})$ , we claim that we can find  $\bar{y}' \in \bar{x}_\kappa + {}^*B(\bar{0}, r_n)$ ,  $\bar{y}' \simeq \bar{y}$ , such that  $\bar{x}' \in J_{s_{\bar{p}_{\bar{y}', t_1}}(t_2), t_2}$ , where  $x'_i \sqrt{\kappa} = [x_i \sqrt{\kappa}]$ , for  $1 \leq i \leq 3$ , ( $\#\#\#\#$ ), contradicting the fact that  $\bar{x} \notin W_{t_1, t_2, r_n}$ , so that  $\bar{y} \notin \bar{x}_\kappa + {}^*B(\bar{0}, r_{n-1})$ ,  $\bar{y}' \notin \bar{x}_\kappa + {}^*B(\bar{0}, r_{n-2})$ ,  $\bar{x} \in W_{t_1, t_2, r_{n-2}, 1+\delta}$ . To prove ( $\#\#\#\#$ ), let  $\bar{w}_{\bar{x}'}$  be the nearest neighbour to  $\bar{x}'$ , and let  $\bar{y}' \in R_{t_1}$ , with  $s_{\bar{p}_{\bar{y}'}}(t_2) = \bar{w}_{\bar{x}'}$ . Suppose that  $|\bar{y}' - \bar{y}| > \epsilon$ , for some  $\epsilon \in \mathcal{R}_{>0}$ , then, using the proof of Lemma 1.12, we have that  $|\bar{w}_{\bar{x}'} - s_{\bar{p}_{\bar{y}, t_1}}(t_2)| \geq D\epsilon$ , for some  $D \in \mathcal{R}_{>0}$ . As  $\bar{w}_{\bar{x}'} \simeq \bar{x}' \simeq \bar{x} \simeq s_{\bar{p}_{\bar{y}, t_1}}(t_2)$ , we obtain a contradiction, so that  $\bar{y}' \simeq \bar{y}$ . As  $\bar{y} \in \bar{x}_\kappa + {}^*B(\bar{0}, r_{n-1})$ , and  $\bar{y}' \simeq \bar{y}$ , clearly  $\bar{y}' \in \bar{x}_\kappa + {}^*B(\bar{0}, r_n)$ , as required. It follows that  $(V_{t_1, t_2} \setminus W_{t_1, t_2, r_n}) \subset W_{t_1, t_2, r_{n-2}, 1+\delta}$ . It follows, using the final result of Lemma 1.26, that;

$$\begin{aligned}
&L(\mu_\kappa)((V_{t_1, t_2} \setminus W_{t_1, t_2, r_n})) \leq L(\mu_\kappa)(W_{t_1, t_2, r_{n-2}, 1+\delta}) \\
&\simeq \mu_\kappa(W_{t_1, t_2, r_{n-2}, 1+\delta}) \\
&\leq D(1 + \delta - r_{n-2})
\end{aligned}$$

An easier calculation shows that, for arbitrary  $\delta \in \mathcal{R}_{>0}$ , we can find a constant  $E \in \mathcal{R}_{>0}$ , with;

$$\begin{aligned} & L(\mu_\kappa)((V \setminus W_{t_1, r_n})) \\ & \leq E(1 + \delta - r_n) \text{ (***)} \end{aligned}$$

Letting  $n \rightarrow \infty$ , and using the fact that  $\delta \in \mathcal{R}_{>0}$ , was arbitrary, the results of (###), (\*\*\*) and (\*\*\*), we conclude;

$$\int_{V_{t_1, t_2}} \circ \rho_\kappa dL(\mu_\kappa) \simeq \int_V \circ \rho_\kappa dL(\mu_\kappa)$$

and, therefore;

$$\int_{V_{t_1, t_2}} \circ \rho_\kappa dL(\mu_\kappa) = \int_V \circ \rho_\kappa dL(\mu_\kappa)$$

as required. □

**Definition 1.25.** Let  $\{r, r_1, r_2\} \subset \mathcal{R}_{>0}$ , with  $r_1 < r_2$  and  $\{t_1, t_2\} \subset \mathcal{T}_\nu$  finite, with  $t_1 < t_2$ , then we define;

$$U_{t_1, r} = R_t \cap (\bar{x}_K + {}^*B(\bar{0}, r))$$

$$V_{t_1, r} = {}^*\bigcup_{\bar{w} \in U_{t_1, r}} J_{\bar{w}, t_1}$$

$$W_{t_1, r} = {}^*\bigcup_{\bar{x} \in V_{t_1, r}} V_{\bar{x}}$$

$$U_{t_1, t_2, r} = \{s_{\bar{p}_{\bar{w}, t_1}}(t_2) : \bar{w} \in U_{t_1, r}\}$$

$$V_{t_1, t_2, r} = {}^*\bigcup_{\bar{w} \in U_{t_1, t_2, r}} J_{\bar{w}, t_2}$$

$$W_{t_1, t_2, r} = {}^*\bigcup_{\bar{x} \in V_{t_1, t_2, r}} V_{\bar{x}}$$

$$U_{t_1, r_1, r_2} = R_t \cap (\bar{x}_K + {}^*A(\bar{0}, r_1, r_2))$$

$$V_{t_1, r_1, r_2} = {}^*\bigcup_{\bar{w} \in U_{t_1, r}} J_{\bar{w}, t_1}$$

$$W_{t_1, r_1, r_2} = {}^*\bigcup_{\bar{x} \in V_{t_1, r}} V_{\bar{x}}$$

$$U_{t_1, t_2, r_1, r_2} = \{s_{\bar{p}_{\bar{w}, t_1}}(t_2) : \bar{w} \in U_{t_1, r}\}$$

$$V_{t_1, t_2, r_1, r_2} = {}^* \bigcup_{\bar{w} \in U_{t_1, t_2, r}} J_{\bar{w}, t_2}$$

$$W_{t_1, t_2, r_1, r_2} = {}^* \bigcup_{\bar{x} \in V_{t_1, t_2, r}} V_{\bar{x}}$$

where  $B(\bar{0}, r) = \{\bar{z} \in \mathcal{R}^3 : |\bar{z}| < r\}$  and  $A(\bar{0}, r_1, r_2) = \{\bar{z} \in \mathcal{R}^3 : r_1 < |\bar{z}| < r_2\}$

**Lemma 1.26.** *We have that  $\{U_{t_1, r}, V_{t_1, r}, U_{t_1, t_2, r}, V_{t_1, t_2, r}\}$  are  $*$ -finite,  $\{W_{t_1, r}, W_{t_1, t_2, r}\}$  are  $\mu_\kappa$  measurable, and there exists a constant  $C \in \mathcal{R}_{>0}$ , with  $\mu_\kappa(W_{t_1, t_2, r}) \leq Cr^3$ . Moreover;*

$$\int_{W_{t_1, r}} \rho_\kappa d\mu_\kappa = \int_{W_{t_1, t_2, r}} \rho_\kappa d\mu_\kappa$$

*We have that  $\{U_{t_1, r_1, r_2}, V_{t_1, r_1, r_2}, U_{t_1, t_2, r_1, r_2}, V_{t_1, t_2, r_1, r_2}\}$  are  $*$ -finite,  $\{W_{t_1, r_1, r_2}, W_{t_1, t_2, r_1, r_2}\}$  are  $\mu_\kappa$  measurable, and there exists a constant  $D \in \mathcal{R}_{>0}$ , with  $\mu_\kappa(W_{t_1, t_2, r}) \leq D(r_2 - r_1)$ ,  $r_1 > \frac{r_2}{3}$ ,  $r_2 < 1$ .*

*Proof.* The fact that  $U_{t_1, r}$  is  $*$ -finite follows from the facts that  $(\bar{x}_K + {}^*B(\bar{0}, r))$  is internal,  $R_t$  is  $*$ -finite, and an internal subset of a  $*$ -finite set is  $*$ -finite.  $V_{t_1, r}$  is  $*$ -finite as each  $J_{\bar{w}, t_1}$  is  $*$ -finite,  $U_{t_1, r}$  is  $*$ -finite and a  $*$ -finite union of  $*$ -finite sets is  $*$ -finite.  $U_{t_1, t_2, r}$  is  $*$ -finite as it is in internal bijection with  $U_{t_1, r}$ , and  $V_{t_1, t_2, r}$  is  $*$ -finite, for the same reason as  $V_{t_1, r}$ .  $W_{t_1, r}$  and  $W_{t_1, t_2, r}$  are  $\mu_\kappa$  measurable, using Definition 1.20. For the next claim, choose  $\bar{w}_0 \in U_{t_1, r}$ , then, for any  $\bar{w} \in U_{t_1, r}$ , we have that  $|\bar{w} - \bar{w}_0| < 2r$ . It follows, from the proof of Lemma 1.9, that there exists a constant  $E \in \mathcal{R}_{>0}$ , with  $|\bar{w}' - s_{\bar{p}_{\bar{w}_0, t_1}}(t_2)| < Er$ , for any  $\bar{w}' \in U_{t_1, t_2, r}$ . If  $\bar{x} \in V_{t_1, t_2, r}$ , then there exists  $\bar{w}' \in U_{t_1, t_2, r}$  with  $\bar{x} \in J_{\bar{w}', t_2}$ , and, by  $(\dagger)$  in Lemma 1.19, we have there exists a constant  $F \in \mathcal{R}_{>0}$ , with  $|\bar{x} - \bar{w}'| < \frac{F}{\sqrt{\eta}}$ . It follows, by the triangle inequality, and the fact that  $\frac{1}{\sqrt{\eta}}$  is infinitesimal, that  $|\bar{x} - s_{\bar{p}_{\bar{w}_0, t_1}}(t_2)| < Er + \frac{F}{\sqrt{\eta}} < (E+1)r$ . In particular;

$$W_{t_1, t_2, r} \subset {}^*B(s_{\bar{p}_{\bar{w}_0, t_1}}(t_2), (E+2)r)$$

$$\subset {}^*B({}^\circ s_{\bar{p}_{\bar{w}_0, t_1}}(t_2), (E+3)r)$$

$$\subset st^{-1}(B({}^\circ s_{\bar{p}_{\bar{w}_0, t_1}}(t_2), (E+4)r))$$

It follows that;

$$\begin{aligned} \mu_\kappa(W_{t_1, t_2, r}) &\simeq L(\mu_\kappa)(W_{t_1, t_2, r}) \\ &\leq L(\mu_\kappa)(st^{-1}(B({}^\circ s_{\bar{p}_{\bar{w}_0, t_1}}(t_2), (E+4)r))) \\ &= \frac{4\pi(E+1)^3 r^3}{3} \end{aligned}$$

and;

$$\mu_\kappa(W_{t_1, t_2, r}) \leq Cr^3, \text{ where } C = \frac{4\pi(E+2)^3}{3}.$$

We have, using Definition 1.18, that;

$$\begin{aligned} \int_{W_{t_1, r}} \rho_\kappa d\mu_\kappa &= \frac{1}{\kappa^{\frac{3}{2}}} * \sum_{\bar{x} \in V_{t_1, r}} \frac{\kappa^{\frac{3}{2}} \rho_\eta(\bar{p}_{\bar{x}, t_1}, 0)}{\eta^{\frac{3}{2}} \lambda_{\bar{w}, t_1}} \\ &= * \sum_{\bar{x} \in V_{t_1, r}} \frac{\rho_\eta(\bar{p}_{\bar{x}, t_1}, 0)}{\eta^{\frac{3}{2}} \lambda_{\bar{w}, t_1}} \\ &= * \sum_{\bar{w} \in U_{t_1, r}} \frac{\rho_\eta(\bar{p}_{\bar{w}, t_1}, 0)}{\eta^{\frac{3}{2}}} \\ &= * \sum_{\bar{w} \in U_{t_1, t_2, r}} \frac{\rho_\eta(\bar{p}_{\bar{w}, t_2}, 0)}{\eta^{\frac{3}{2}}} \\ &= * \sum_{\bar{x} \in V_{t_1, t_2, r}} \frac{\rho_\eta(\bar{p}_{\bar{x}, t_2}, 0)}{\eta^{\frac{3}{2}} \lambda_{\bar{w}, t_2}} \\ &= \frac{1}{\kappa^{\frac{3}{2}}} * \sum_{\bar{x} \in V_{t_1, t_2, r}} \frac{\kappa^{\frac{3}{2}} \rho_\eta(\bar{p}_{\bar{x}, t_2}, 0)}{\eta^{\frac{3}{2}} \lambda_{\bar{w}, t_2}} = \int_{W_{t_1, t_2, r}} \rho_\kappa d\mu_\kappa \end{aligned}$$

Similarly,  $\{U_{t_1, r_1, r_2}, V_{t_1, r_1, r_2}, U_{t_1, t_2, r_1, r_2}, V_{t_1, t_2, r_1, r_2}\}$  are \*-finite, replacing  $*B(\bar{0}, r)$  with  $*A(\bar{0}, r_1, r_2)$ , and the same for  $\{W_{t_1, r_1, r_2}, W_{t_1, t_2, r_1, r_2}\}$  being  $\mu_\kappa$  measurable. For the last claim, we can cover  $A(\bar{0}, r_1, r_2)$  with  $N_{r_1, r_2}$  open balls  $B(\bar{x}_i, \frac{r_2 - r_1}{2})$ , where each  $\bar{x}_i \in A(\bar{0}, r_1, r_2)$ , such that, for  $r_1 > \frac{r_2}{3}$ , each  $B(\bar{x}_i, \frac{r_2 - r_1}{2}) \subset A(\bar{0}, \frac{3r_1 - r_2}{2}, \frac{3r_2 - r_1}{2})$ , We have that;

$$N_{r_1, r_2} \text{vol}(B(\bar{x}_i, \frac{r_2 - r_1}{2})) - \text{vol}(T) \leq \text{vol}(A(\bar{0}, \frac{3r_1 - r_2}{2}, \frac{3r_2 - r_1}{2}))$$

where  $T$  is the set of overlaps, and clearly  $\text{vol}(T) \leq \text{vol}(A(\bar{0}, \frac{3r_1 - r_2}{2}, \frac{3r_2 - r_1}{2}))$ , so that;

$$\text{vol}(B(\bar{x}_i, \frac{r_2 - r_1}{2})) \leq \frac{2\text{vol}(A(\bar{0}, \frac{3r_1 - r_2}{2}, \frac{3r_2 - r_1}{2}))}{N_{r_1, r_2}} (*)$$

and a corresponding result holds for the transfers. We have that;

$$V_{t_1, r_1, r_2} \subset \bigcup_{1 \leq i \leq N_{r_1, r_2}} V_{\bar{x}_i, t_1, r'}$$

where  $V_{\bar{x}_i, t_1, r'}$  is defined by  $\bar{x}_\kappa + {}^*B(\bar{x}_i, r')$  and, by  $(*)$ ,  $r' \leq \left(\frac{3\text{vol}(A(\bar{0}, \frac{3r_1-r_2}{2}, \frac{3r_2-r_1}{2}))}{2\pi N_{r_1, r_2}}\right)^{\frac{1}{3}}$ ,  $(**)$ . By the previous result, using  $(**)$ , we have that;

$$\begin{aligned} V_{t_1, t_2, r_1, r_2} &\subset \bigcup_{1 \leq i \leq N_{r_1, r_2}} V_{\bar{x}_i, t_1, t_2, r'} \\ \mu_\kappa(W_{t_1, t_2, r_1, r_2}) &\leq \mu_\kappa\left(\bigcup_{1 \leq i \leq N_{r_1, r_2}} V_{\bar{x}_i, t_1, t_2, r'}\right) \\ &\leq N_{r_1, r_2} \max_i(\mu_\kappa(V_{\bar{x}_i, t_1, t_2, r'})) \\ &\leq CN_{r_1, r_2} \frac{4\pi r'^3}{3} \\ &\leq \frac{4\pi CN_{r_1, r_2}}{3} \frac{3\text{vol}(A(\bar{0}, \frac{3r_1-r_2}{2}, \frac{3r_2-r_1}{2}))}{2\pi N_{r_1, r_2}} \\ &= 6C\text{vol}(A(\bar{0}, \frac{3r_1-r_2}{2}, \frac{3r_2-r_1}{2})) \\ &\leq 6\pi C\left(\frac{r_2-r_1}{2}\right)(4(r_2+r_1)^2 + 16(r_2-r_1)^2) \\ &= 3\pi C(r_2-r_1)(4(r_2+r_1)^2 + 16(r_2-r_1)^2) \\ &\leq 3\pi C(r_2-r_1)(16+16) \\ &= D(r_2-r_1) \end{aligned}$$

where  $D = 96C$ .

□

**Lemma 1.27.** *Let  $\{\bar{x}, \bar{y}\} \subset R$  be finite, and  $t \in \mathcal{T}_\nu$  finite, with  $\bar{x} \neq \bar{y}$ , then, for  $\nu \geq \eta^2$ , there exists  $C_{t, \bar{x}, \bar{y}} \in \mathcal{R}_{>0}$ , with;*

$$\left((1 + \frac{C_t}{\nu})^{-[\nu t]} - \frac{1}{\nu^{\frac{1}{2}}}\right)|\bar{x} - \bar{y}| \leq |s_{\bar{x}}(t) - s_{\bar{y}}(t)| \leq (1 + \frac{C_t}{\nu})^{[\nu t]}|\bar{x} - \bar{y}|$$

*Proof.* The second inequality follows by a simple adaptation of the proof of the last part of Lemma 1.9, noting that  $((1 + \frac{C_t}{\nu})^{[\nu t]-1} \leq ((1 + \frac{C_t}{\nu})^{[\nu t]}$ . For the second part, suppose that;

$$|s_{\bar{x}}(t) - s_{\bar{y}}(t)| < \left( \left(1 + \frac{C_t}{\nu}\right)^{-[\nu t]} - \frac{1}{\nu^{\frac{1}{2}}} \right) |\bar{x} - \bar{y}| \quad (*)$$

As in Lemma 1.12, using (\*), letting  $\bar{x}' = s_{\bar{x}}(t)$  and  $\bar{y}' = s_{\bar{y}}(t)$ , we have that;

$$\begin{aligned} |s_{\bar{x}',t}(t) - s_{\bar{y}',t}(t)| &\leq \left(1 + \frac{C_t}{\nu}\right)^{[\nu t]} |s_{\bar{x}}(t) - s_{\bar{y}}(t)| \\ &< \left(1 + \frac{C_t}{\nu}\right)^{[\nu t]} \left( \left(1 + \frac{C_t}{\nu}\right)^{-[\nu t]} - \frac{1}{\nu^{\frac{1}{2}}} \right) |\bar{x} - \bar{y}| \\ &= \left(1 - \frac{\left(1 + \frac{C_t}{\nu}\right)^{[\nu t]}}{\nu^{\frac{1}{2}}}\right) |\bar{x} - \bar{y}| \quad (**)$$

By (††), (†††), of Lemma 1.12, there exists  $E_{t,\bar{x},\bar{y}} \in \mathcal{R}_{>0}$ , with;

$$\max(|\bar{x} - s_{\bar{x}',t}(t)|, |\bar{y} - s_{\bar{y}',t}(t)|) \leq \frac{E_{t,\bar{x},\bar{y}}}{\nu} \quad (***)$$

so that, by the triangle inequality, using (\*\*), (\*\*\*) ;

$$|\bar{x} - \bar{y}| < \frac{2E_{t,\bar{x},\bar{y}}}{\nu} + \left(1 - \frac{\left(1 + \frac{C_t}{\nu}\right)^{[\nu t]}}{\nu^{\frac{1}{2}}}\right) |\bar{x} - \bar{y}| \quad (\dagger)$$

If  $\nu \geq \eta^2$ , then, using the fact that  $|\bar{x} - \bar{y}| \geq \frac{1}{\sqrt{\eta}}$ , so that  $\frac{1}{|\bar{x} - \bar{y}|} \leq \sqrt{\eta}$ ;

$$\nu > \frac{2E_{t,\bar{x},\bar{y}} \nu^{\frac{1}{2}} \eta^{\frac{1}{2}}}{\left(1 + \frac{C_t}{\nu}\right)^{[\nu t]}} \geq \frac{2E_{t,\bar{x},\bar{y}} \nu^{\frac{1}{2}}}{\left(1 + \frac{C_t}{\nu}\right)^{[\nu t]} |\bar{x} - \bar{y}|}$$

so that;

$$\frac{2E_{t,\bar{x},\bar{y}}}{\nu} - \frac{\left(1 + \frac{C_t}{\nu}\right)^{[\nu t]}}{\nu^{\frac{1}{2}}} |\bar{x} - \bar{y}| < 0$$

and, from (†), we obtain the contradiction, that  $|\bar{x} - \bar{y}| < |\bar{x} - \bar{y}|$  so (\*) fails, and we obtain the result.

□

**Lemma 1.28.** *With  $\nu \geq \eta^2$ , let  $\bar{x} \in \bar{x}_\kappa + R_\kappa$  be finite and let  $\{t_1, t_2\} \subset \mathcal{T}_\nu$  be finite, with  $t_1 < t_2$ , and  $t_1 \simeq t_2$ . Let  $\bar{w}_1 \in R_{t_1}$  and  $\bar{w}_2 \in R_{t_2}$  be the nearest neighbours to  $\bar{x}$  at  $t_1$  and  $t_2$ , then  $\bar{p}_{\bar{w}_1, t_1} \simeq \bar{p}_{\bar{w}_2, t_2}$ .*

*Proof.* Suppose, for contradiction, that  $|\bar{p}_{\bar{w}_1, t_1} - \bar{p}_{\bar{w}_2, t_2}| \geq \epsilon$ , for some  $\epsilon \in \mathcal{R}_{>0}$ . We have, using (†) of Lemma 1.19 that;

$$|s_{\bar{p}_{\bar{w}_1, t_1}}(t_1) - s_{\bar{p}_{\bar{w}_2, t_2}}(t_2)|$$

$$\begin{aligned}
&= |\bar{w}_1 - \bar{w}_2| \\
&\leq |\bar{w}_1 - \bar{x}| + |\bar{w}_2 - \bar{x}| \\
&\leq \frac{2D}{\sqrt{\eta}} \simeq 0 \text{ (*) (can be improved for small } t_2 - t_1)
\end{aligned}$$

where  $D = (3 + 2F)e^{\circ t_1} + 6$ . We have that, using Lemma 1.7, that;

$$|s_{\bar{p}_{\bar{w}_1, t_1}}(t_2) - s_{\bar{p}_{\bar{w}_1, t_1}}(t_1)| \leq \frac{G([t_2\nu] - [t_1\nu])}{\nu} \simeq 0 (**)$$

as  $G \in \mathcal{R}_{>0}$  and  $t_1 \simeq t_2$ , so that, using (\*), (\*\*);

$$\begin{aligned}
&|s_{\bar{p}_{\bar{w}_1, t_1}}(t_2) - s_{\bar{p}_{\bar{w}_2, t_2}}(t_2)| \\
&\leq |s_{\bar{p}_{\bar{w}_1, t_1}}(t_2) - s_{\bar{p}_{\bar{w}_1, t_1}}(t_1)| + |s_{\bar{p}_{\bar{w}_1, t_1}}(t_1) - s_{\bar{p}_{\bar{w}_2, t_2}}(t_2)| \simeq 0 (***)
\end{aligned}$$

By Lemma 1.27, we have that;

$$\begin{aligned}
&|s_{\bar{p}_{\bar{w}_1, t_1}}(t_2) - s_{\bar{p}_{\bar{w}_2, t_2}}(t_2)| \geq \left( \left(1 + \frac{C_{t_2}}{\nu}\right)^{-[\nu t]} - \frac{1}{\nu^{\frac{1}{2}}} \right) |\bar{p}_{\bar{w}_1, t_1} - \bar{p}_{\bar{w}_2, t_2}| \\
&= \left( \left(1 + \frac{C_{t_2}}{\nu}\right)^{-[\nu t]} - \frac{1}{\nu^{\frac{1}{2}}} \right) \epsilon \\
&> \frac{(1 + \frac{C_{t_2}}{\nu})^{-[\nu t]} \epsilon}{2} > \frac{e^{-C_{t_2} \epsilon}}{4}
\end{aligned}$$

which contradicts (\*\*\*), as required. □

**Definition 1.29.** For  $\{n, m\} \subset \mathcal{N}$ , with the rounding down convention on  $[\cdot]$ , we define  $R_n = \{\bar{x} : \frac{[x_i \sqrt{n}]}{\sqrt{n}} = x_i, \text{ for } 1 \leq i \leq 3\}$ . For  $\bar{x} \in R_n$  finite, we define a path  $\bar{s}_{n, m, \bar{x}}$  inductively on  $\mathcal{T}_m$  by;

$$\bar{s}_{n, m, \bar{x}}\left(\frac{j+1}{m}\right) - \bar{s}_{n, m, \bar{x}}\left(\frac{j}{m}\right) = \frac{1}{m} \frac{\bar{J}(\bar{s}_{n, m, \bar{x}}(\frac{j}{m}), \frac{j}{m})}{\rho(\bar{s}_{n, m, \bar{x}}(\frac{j}{m}), \frac{j}{m})}, \text{ for } 0 \leq j \leq m^2 - 1$$

For infinite  $\{\eta, \nu\} \subset {}^*\mathcal{N}$ , let  $\bar{s}_{\eta, \nu} : R_\eta \times \mathcal{T}_\nu \rightarrow \mathcal{R}_\eta^3$ , be defined by  $\bar{s}_{\eta, \nu}(\bar{x}, \frac{j}{\nu}) = \bar{s}_{\eta, \nu, \bar{x}}(\frac{j}{\nu})$ ,  $0 \leq j \leq \nu^2$ , where  $\bar{s}_{\eta, \nu, \bar{x}}$  is obtained by transfer from  $\bar{s}_{n, m, \bar{x}}$ , using  $\frac{{}^*\bar{J}}{{}^*\rho}$  and the usual wrap around convention. Denote also by  $\bar{s}_{\eta, \nu}$ , the  $\mu_\eta \times \mu_\nu$  measurable version.

**Lemma 1.30.** For all infinite  $\{\eta, \nu\} \subset {}^*\mathcal{N}$ , finite  $(\bar{x}, t) \in \mathcal{R}_\eta^3 \times \mathcal{T}_\nu$ ;

$$\left(\frac{\partial \bar{s}_{\eta,\nu}}{\partial x_i}\right)_\eta|_{\bar{x},t}, \left(\frac{\partial \bar{s}_{\eta,\nu}}{\partial t}\right)_\nu|_{\bar{x},t} \text{ and } \left(\frac{\partial^2 \bar{s}_{\eta,\nu}}{\partial x_i \partial t}\right)_{\eta,\nu}|_{\bar{x},t}$$

are finite, for  $1 \leq i \leq 3$ , where, for  $\bar{f} : \mathcal{R}_\eta \times \mathcal{T}_\nu \rightarrow \mathcal{R}_\eta^3$ ,  $\mu_\eta \times \mu_\nu$  measurable,  $\left(\frac{\partial \bar{f}}{\partial x_i}\right)_\eta$  and  $\left(\frac{\partial \bar{f}}{\partial t}\right)_\nu$  are the nonstandard derivatives. In particular,  $\bar{s}_{\eta,\nu}$  is  $S$ -continuous, in the sense that, for  $\{(\bar{x}, t), (\bar{x}', t')\} \subset \mathcal{R}_\eta^3 \times \mathcal{T}_\nu$ , with  $(\bar{x}, t) \simeq (\bar{x}', t')$ ;

$$\bar{s}_{\eta,\nu}(\bar{x}, t) \simeq \bar{s}_{\eta,\nu}(\bar{x}', t')$$

*Proof.* For  $\{n, m\} \subset \mathcal{N}$ , letting the initial condition  $\bar{x}$  vary over  $\mathcal{R}_n^3$ , we have, using the chain rule and the inductive definition, that, for  $0 \leq k \leq m^2 - 1$ , the functions  $\bar{s}_{k,n,m}(\bar{x}) = \bar{s}_{n,m,\bar{x}}(\bar{x}, \frac{k}{m})$  are smooth. We have, for  $1 \leq i \leq 3$  that;

$$\begin{aligned} \frac{\partial \bar{s}_{k+1,n,m}}{\partial x_i} &= \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \\ &+ \frac{1}{m} \left( \text{grad}\left(\frac{j_1}{\rho}\right) \cdot \frac{\partial \bar{s}_{k,n,m}}{\partial x_i}, \text{grad}\left(\frac{j_2}{\rho}\right) \cdot \frac{\partial \bar{s}_{k,n,m}}{\partial x_i}, \text{grad}\left(\frac{j_3}{\rho}\right) \cdot \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right) \\ \left| \frac{\partial \bar{s}_{k+1,n,m}}{\partial x_i} \right| &\leq \left| \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| \\ &+ \frac{1}{m} \left( \left| \text{grad}\left(\frac{j_1}{\rho}\right) \cdot \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| + \left| \text{grad}\left(\frac{j_2}{\rho}\right) \cdot \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| + \left| \text{grad}\left(\frac{j_3}{\rho}\right) \cdot \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| \right) \\ &\leq \left| \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| + \frac{1}{m} \max\left(\left| \text{grad}\left(\frac{j_1}{\rho}\right) \right|, \left| \text{grad}\left(\frac{j_2}{\rho}\right) \right|, \left| \text{grad}\left(\frac{j_3}{\rho}\right) \right|\right) \left| \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| \\ &= \left(1 + \frac{G_i}{m}\right) \left| \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| \end{aligned}$$

It follows, by a simple induction, that, for  $1 \leq i \leq 3$ ;

$$\begin{aligned} \left| \frac{\partial \bar{s}_{k,n,m}}{\partial x_i} \right| &\leq \left(1 + \frac{G_i}{m}\right)^k \left| \frac{\partial \bar{s}_{0,n,m}}{\partial x_i} \right| \\ &= \left(1 + \frac{G_i}{m}\right)^k \\ &= \left(\left(1 + \frac{G_i}{m}\right)^m\right)^{\frac{k}{m}} \\ &\leq (e^{G_i} + 1)^t, \text{ for } \frac{k}{m} \leq t, t \in \mathcal{R}_{>0}. \end{aligned}$$

In particular,  $\left| \frac{\partial (\bar{s}_{k,n,m})_j}{\partial x_i} \right| \leq (e^{G_i} + 1)^t$ , for  $1 \leq j \leq 3$ .

Using the proof of Lemma 1.7, we have that, for  $\{\bar{x}, \bar{x}'\} \subset \mathcal{R}_n^3$ ,  $1 \leq j \leq 3$ , that;

$$\begin{aligned}
& |(\bar{s}_{k,n,m})_j(\bar{x}') - (\bar{s}_{k,n,m})_j(\bar{x})| \leq |\text{grad}((\bar{s}_{k,n,m})_j)| |\bar{x}' - \bar{x}| \\
& \leq (|\frac{\partial(\bar{s}_{k,n,m})_j}{\partial x_1}| + |\frac{\partial(\bar{s}_{k,n,m})_j}{\partial x_2}| + |\frac{\partial(\bar{s}_{k,n,m})_j}{\partial x_3}|) |\bar{x}' - \bar{x}| \\
& \leq ((e^{G_1} + 1)^t + (e^{G_2} + 1)^t + (e^{G_3} + 1)^t) |\bar{x}' - \bar{x}|
\end{aligned}$$

and;

$$\begin{aligned}
& |\bar{s}_{k,n,m}(\bar{x}') - \bar{s}_{k,n,m}(\bar{x})| \leq \sum_{j=1}^3 |(\bar{s}_{k,n,m})_j(\bar{x}') - (\bar{s}_{k,n,m})_j(\bar{x})| \\
& \leq 3((e^{G_1} + 1)^t + (e^{G_2} + 1)^t + (e^{G_3} + 1)^t) |\bar{x}' - \bar{x}|
\end{aligned}$$

In particular, it follows that, for  $1 \leq i \leq 3$ ;

$$\begin{aligned}
& |\sqrt{m}(\bar{s}_{k,n,m})(\bar{x} + \frac{\bar{e}_i}{\sqrt{m}}) - \sqrt{m}(\bar{s}_{k,n,m})(\bar{x})| \\
& \leq 3((e^{G_1} + 1)^t + (e^{G_2} + 1)^t + (e^{G_3} + 1)^t) \sqrt{m} |\frac{\bar{e}_i}{\sqrt{m}}| \\
& = 3((e^{G_1} + 1)^t + (e^{G_2} + 1)^t + (e^{G_3} + 1)^t)
\end{aligned}$$

It follows that for fixed  $\{r, s\} \subset \mathcal{R}_{>0}$

$$\begin{aligned}
\mathcal{R} \models & (\forall n \in \mathcal{N})(\forall m \in \mathcal{N})(\forall t, 0 \leq t \leq s, t \in \mathcal{T}_m)(\forall \bar{x}, |\bar{x}| \leq r \\
& , \bar{x} \in \mathcal{R}_n^3) |(\frac{\partial \bar{s}_{n,m}}{\partial x_i})_m |_{\bar{x}, t}| \leq C_s
\end{aligned}$$

for  $1 \leq i \leq 3$ , where  $C_s = 3((e^{G_1} + 1)^s + (e^{G_2} + 1)^s + (e^{G_3} + 1)^s)$

By transfer, the above statement holds in  $^*\mathcal{R}$ , in particular, For all infinite  $\{\eta, \nu\} \subset ^*\mathcal{N}$ , finite  $(\bar{x}, t) \in \mathcal{R}_\eta^3 \times \mathcal{T}_\nu$ , with  $|t| \leq s$ ;

$$|(\frac{\partial \bar{s}_{\eta,\nu}}{\partial x_i})_\eta |_{\bar{x}, t}| \leq C_s, (*)$$

and is finite, for  $1 \leq i \leq 3$ . We have, by the definition of a trajectory, the nonstandard partial derivative, using the bound  $G_t$  from Lemma 1.7 that;

$$\begin{aligned}
& |(\frac{\partial \bar{s}_{\eta,\nu}}{\partial t})_\nu |_{\bar{x}, t}| \\
& = |\nu(\bar{s}_{\eta,\nu}(\bar{x}, t + \frac{1}{\nu}) - \bar{s}_{\eta,\nu}(\bar{x}, t))|
\end{aligned}$$

$$= \left| \frac{*\bar{J}(\bar{s}_{\eta,\nu}(\bar{x},t), \frac{[t\nu]}{\nu})}{*\rho(\bar{s}_{\eta,\nu}(\bar{x},t), \frac{[t\nu]}{\nu})} \right| \leq G_t$$

For the next claim, we have, using the definition of a trajectory, that;

$$\begin{aligned} \left( \frac{\partial^2 \bar{s}_{\eta,\nu}}{\partial x_i \partial t} \right)_{\eta,\nu} |_{\bar{x},t} &= \left( \frac{\partial(\bar{f}_t(\bar{s}_{\eta,\nu}(\bar{x},t)))}{\partial x_i} \right)_{\eta} \\ &= \left( \left( \frac{\partial(f_{1t}(\bar{s}_{\eta,\nu}(\bar{x},t)))}{\partial x_i} \right)_{\eta}, \left( \frac{\partial(f_{2t}(\bar{s}_{\eta,\nu}(\bar{x},t)))}{\partial x_i} \right)_{\eta}, \left( \frac{\partial(f_{3t}(\bar{s}_{\eta,\nu}(\bar{x},t)))}{\partial x_i} \right)_{\eta} \right) \end{aligned}$$

$$\text{where } \bar{f}_t(\bar{y}) = \frac{*\bar{J}(\bar{y}, \frac{[t\nu]}{\nu})}{*\rho(\bar{y}, \frac{[t\nu]}{\nu})}, \text{ and } \bar{f}_t = (f_{1t}, f_{2t}, f_{3t})$$

We have, using Taylor's theorem in the final step, see [7], and (\*), that, for  $1 \leq j \leq 3$ ;

$$\begin{aligned} & \left| \left( \frac{\partial(f_{jt}(\bar{s}_{\eta,\nu}(\bar{x},t)))}{\partial x_i} \right)_{\eta} \right| = \left| \sqrt{\eta} (f_{jt}(\bar{s}_{\eta,\nu}(\bar{x} + \frac{\bar{e}_i}{\sqrt{\eta}}, t)) - f_{jt}(\bar{s}_{\eta,\nu}(\bar{x}, t))) \right| \\ &= \left| \sqrt{\eta} \left( \left( \frac{f_{jt}(y_1+\delta_1, y_2+\delta_2, y_3+\delta_3) - f_{jt}(y_1, y_2+\delta_2, y_3+\delta_3)}{\delta_1}, \frac{f_{jt}(y_1, y_2+\delta_2, y_3+\delta_3) - f_{jt}(y_1, y_2, y_3+\delta_3)}{\delta_2}, \right. \right. \right. \\ & \quad \left. \left. \frac{f_{jt}(y_1, y_2, y_3+\delta_3) - f_{jt}(y_1, y_2, y_3)}{\delta_3} \right) \cdot \bar{\delta} \right| \\ &\leq \left| \left( \frac{f_{jt}(y_1+\delta_1, y_2+\delta_2, y_3+\delta_3) - f_{jt}(y_1, y_2+\delta_2, y_3+\delta_3)}{\delta_1}, \frac{f_{jt}(y_1, y_2+\delta_2, y_3+\delta_3) - f_{jt}(y_1, y_2, y_3+\delta_3)}{\delta_2}, \right. \right. \\ & \quad \left. \left. \frac{f_{jt}(y_1, y_2, y_3+\delta_3) - f_{jt}(y_1, y_2, y_3)}{\delta_3} \right) \right| \left| \sqrt{\eta} \bar{\delta} \right| \\ &\simeq \left| \left( *\frac{\partial g_j}{\partial x_1} \Big|_{(y_1, y_2+\delta_2, y_3+\delta_3, t)}, *\frac{\partial g_j}{\partial x_2} \Big|_{(y_1, y_2, y_3+\delta_3, t)}, *\frac{\partial g_j}{\partial x_3} \Big|_{(y_1, y_2, y_3, t)} \right) \right| \left| \left( \frac{\partial \bar{s}_{\eta,\nu}}{\partial x_i} \right)_{\eta} \Big|_{\bar{x},t} \right| \\ &\leq C_s \left( \left| *\frac{\partial g_j}{\partial x_1} \Big|_{(y_1, y_2+\delta_2, y_3+\delta_3, t)} \right| + \left| *\frac{\partial g_j}{\partial x_2} \Big|_{(y_1, y_2, y_3+\delta_3, t)} \right| + \left| *\frac{\partial g_j}{\partial x_3} \Big|_{(y_1, y_2, y_3, t)} \right| \right) \\ &\leq C_s \left( \sum_{k=1}^3 H_{j,k,t''} \right) \end{aligned}$$

where;

$$\bar{\delta} = (\delta_1, \delta_2, \delta_3) = \bar{s}_{\eta,\nu}(\bar{x} + \frac{\bar{e}_i}{\sqrt{\eta}}, t) - \bar{s}_{\eta,\nu}(\bar{x}, t), \bar{y} = (y_1, y_2, y_3) = \bar{s}_{\eta,\nu}(\bar{x}, t)$$

$$g_j(\bar{x}, t) = \left( \frac{\bar{J}(\bar{x}, t)}{\rho(\bar{x}, t)} \right)_j$$

and we have transferred the fact that, for  $1 \leq k \leq 3$  and  $\{t', t''\} \subset \mathcal{R}_{>0}$ ,  $t' \leq t''$ ,  $t < t''$ ;

$$\left| \frac{\partial g_{j,t'}}{\partial x_k} \right| = \left| \frac{\partial \left( \frac{j_{j,t'}}{\rho_{t'}} \right)}{\partial x_k} \right|$$

$$\begin{aligned}
&= \left| \frac{\frac{\partial j_{j,t'}}{\partial x_k} \rho_{t'} - j_{j,t'} \frac{\partial \rho}{\partial x_k}}{\rho_{t'}^2} \right| \\
&\leq \left| \frac{\partial j_{j,t'}}{\partial x_k} \right| |\rho_{t'}| + |j_{j,t'}| \left| \frac{\partial \rho_{t'}}{\partial x_k} \right| \\
&\leq D_{j,k,t''} E_{j,t''} + F_{j,t''} G_{j,k,t''} = H_{j,k,t''}
\end{aligned}$$

for some  $\{D_{j,k,t''}, E_{j,t''}, F_{j,t''}, G_{j,k,t''}, H_{j,k,t''}\} \subset \mathcal{R}_{>0}$

It follows that;

$$\begin{aligned}
&\left| \left( \frac{\partial^2 \bar{s}_{\eta,\nu}}{\partial x_i \partial t} \right)_{\eta,\nu} |_{\bar{x},t} \right| \leq \sum_{j=1}^3 \left| \frac{\partial (f_{jt}(\bar{s}_{\eta,\nu}(\bar{x},t)))}{\partial x_i} \right|_{\eta} \\
&\leq C_s \sum_{j,k=1}^3 H_{j,k,t''}
\end{aligned}$$

as required.

To verify the last claim, it is enough, by transitivity of  $\simeq$ , to show that, for  $\bar{x}$  finite, and  $\epsilon \simeq 0$ ,  $\delta \simeq 0$ ;

$$\bar{s}_{\eta,\nu}(\bar{x} + \epsilon \bar{e}_1, t) \simeq \bar{s}_{\eta,\nu}(\bar{x}, t)$$

$$\bar{s}_{\eta,\nu}(\bar{x}, t + \delta) \simeq \bar{s}_{\eta,\nu}(\bar{x}, t)$$

We can write  $\epsilon = \frac{r}{\sqrt{\eta}} + r'$ , where  $0 \leq r' < \frac{1}{\sqrt{\eta}}$ , then, using measurability, (\*), and without loss of generality, assuming that  $r \geq 0$ ;

$$\begin{aligned}
&|\bar{s}_{\eta,\nu}(\bar{x} + \epsilon \bar{e}_1, t) - \bar{s}_{\eta,\nu}(\bar{x}, t)| \\
&= |\bar{s}_{\eta,\nu}(\bar{x} + \frac{r}{\sqrt{\eta}} \bar{e}_1, t) - \bar{s}_{\eta,\nu}(\bar{x}, t)| \\
&\leq \sum_{k=0}^{r-1} |\bar{s}_{\eta,\nu}(\bar{x} + \frac{k+1}{\sqrt{\eta}} \bar{e}_1, t) - \bar{s}_{\eta,\nu}(\bar{x} + \frac{k}{\sqrt{\eta}} \bar{e}_1, t)| \\
&\leq \sum_{k=0}^{r-1} \frac{C_t}{\sqrt{\eta}} \\
&= \frac{C_t r}{\sqrt{\eta}} \simeq 0
\end{aligned}$$

Similarly, writing  $\delta = \frac{s}{\nu} + s'$ , where  $0 \leq s' < \frac{1}{\nu}$ , and without loss of generality, assuming  $s \geq 0$ , we can show that;

$$|\bar{s}_{\eta,\nu}(\bar{x}, t + \delta) - \bar{s}_{\eta,\nu}(\bar{x}, t)|$$

$$\leq \frac{G_{t+\delta+1^s}}{\nu} \simeq 0$$

as required.  $\square$

**Definition 1.31.** Let  $\{\bar{w}_i, \bar{w}_j\} \subset R_t$ , with  $\bar{w}_i \neq \bar{w}_j$ , then we let  $H_{ij,t} \subset {}^*\mathcal{R}^3$  be the internal hyperplane, defined by;

$$H_{ij,t} = \{\bar{z} \in {}^*\mathcal{R}^3 : |\bar{z} - \bar{w}_i| = |\bar{z} - \bar{w}_j|\}$$

Fixing  $\bar{w}_i$ , we let  $H_{i,t} \subset {}^*\mathcal{R}^3$  be the internal hyperplane arrangement, defined by;

$$H_{i,t} = {}^*\bigcup_{\bar{x}_j \in R_t, \bar{x}_j \neq \bar{x}_i} H_{ij,t}$$

We define an alcove  $Z \subset {}^*\mathcal{R}^3 \setminus H_{i,t}$  to be a  $*$ -connected component of  ${}^*\mathcal{R}^3 \setminus H_{i,t}$ , and, noting that  $\bar{w}_i \notin H_{i,t}$ , let  $Z_i$  be the unique alcove containing  $\bar{w}_i$ .

**Lemma 1.32.** Let notation be as in Definition 1.31 and 1.18, then, for  $\bar{x}_i$  finite;

$$J_{\bar{x}_i,t} = Z_i \cap (\bar{x}_\kappa + R_\kappa)$$

In particular,  $Z_i$  is  $*$ -bounded.

*Proof.* By Lemma 1.17, we have that  $\bar{x}_\kappa + R_\kappa \subset {}^*\mathcal{R}^3 \setminus H_{i,t}$ . We claim, firstly, that  $Z_i \cap (\bar{x}_\kappa + R_\kappa) \subset J_{\bar{x}_i,t}$ . Suppose, for contradiction, that we can find  $\bar{w} \in Z_i \cap (\bar{x}_\kappa + R_\kappa)$ , with  $\bar{w} \notin J_{\bar{x}_i,t}$ , then we can find  $\bar{w}_j \in R_t$ , with  $|\bar{w} - \bar{w}_j| < |\bar{w} - \bar{w}_i|$ . It follows that any  $*$ -continuous path from  $\bar{w}$  to  $\bar{w}_i$ , would pass through  $H_{ij,t}$ , in particular  $Z_i$  cannot be  $*$ -connected. We claim, second, that  $J_{\bar{x}_i,t} \subset Z_i \cap (\bar{x}_\kappa + R_\kappa)$ . Suppose, for contradiction, that we can find  $\bar{w} \in J_{\bar{x}_i,t}$ , with  $\bar{w} \notin Z_i \cap (\bar{x}_\kappa + R_\kappa)$ , then, letting  $l_{\bar{w},\bar{x}_i}$ , be the line connecting  $\bar{w}$  and  $\bar{x}_i$ , we must have that  $l_{\bar{w},\bar{x}_i}$  intersects one of the hyperplanes  $H_{ij,t}$ , otherwise we could construct a  $*$ -connected space  $Z_i \cup l_{\bar{w},\bar{x}_i} \supset Z_i$ , contradicting the definition of a component  $Z_i$  as maximal. It follows that  $|\bar{w} - \bar{x}_j| < |\bar{w} - \bar{x}_i|$  and  $\bar{x}_i$  is not the nearest neighbour. For the final claim, if  $Z_i$  is not  $*$ -bounded, we can, using the transfer principle, for a given  $\epsilon \in \mathcal{R}_{>0}$ , find  $\bar{w}_{\eta,\epsilon}$  such that  ${}^*B(\bar{w}_{\eta,\epsilon}, \epsilon) \subset Z_i$ , and  $\bar{w}_{\eta,\epsilon} \in \bar{x}_\kappa + \mathcal{R}_\kappa^3$  is infinite. Clearly, we can find  $\bar{s} \in (\bar{x}_\kappa + R_\kappa) \cap {}^*B(\bar{w}_{\eta,\epsilon}, \epsilon) \subset Z_i$ , so that  $\bar{s}$  is infinite as well. By the previous result, we have that  $\bar{s} \in J_{\bar{x}_i,t}$ , contradicting the fact in (†) of Lemma 1.19, that  $\bar{s} \in {}^*B(\bar{x}_i, \frac{C}{\sqrt{\eta}})$ , so  $\bar{s}$  is finite.  $\square$

**Definition 1.33.** Let notation be as in Definition 1.31, let  $Z_i$  be an alcove corresponding to a finite  $\bar{w}_i \in R_t$ , let  $S_{i,t}$  be the  $*$ -finite set enumerating the  $*$ -convex components  $V_{ik,t}$  of  $\delta Z_i$ , with  $\dim(V_{ik,t}) = 2$ , corresponding to hyperplanes  $H_{ik,t}$ , so that  $\delta Z_i = * \bigcup_{k \in S_{i,t}} V_{ik,t}$ . We define  $T_{i,t}$ , the relevant points to  $\bar{w}_i$ , by;

$$T_{i,t} = \{\bar{w}_k \in R_t : k \in S_{i,t}, V_{ik,t} \subset H_{ik,t}\}$$

We have that  $\bar{Z}_i$  is convex, so that, by the Krein-Milman Theorem,  $\bar{Z}_i$  is the  $*$ -convex hull of the  $*$ -finite set  $V_{i,t}$  of extreme points, which we refer to as vertices.

**Lemma 1.34.** If  $\bar{v} \in V_{i,t}$ , then there exist  $\{\bar{w}_{k_1}, \bar{w}_{k_2}, \bar{w}_{k_3}\} \subset T_{i,t}$ , such that;

$$\bar{v} = H_{ik_1,t} \cap H_{ik_2,t} \cap H_{ik_3,t}$$

*Proof.* As  $\bar{v} \in V_{i,t} \subset \delta Z_i$ , we can find  $\bar{w}_{k_1} \in T_{i,t}$ , with  $\bar{v} \in H_{ik_1,t}$ . Choose a point  $\bar{p} \in V_{ik_1,t}$  with  $\bar{p} \neq \bar{v}$ , Suppose that  $\bar{v} \notin * \bigcup_{k \in S_{i,t}, k \neq k_1} H_{ik,t}$ , so that, as  $* \bigcup_{k \in S_{i,t}, k \neq k_1} H_{ik,t}$  is  $*$ -closed, there exists  $\epsilon \in * \mathcal{R}_{>0}$ , with  $*B(\bar{v}, \epsilon) \cap * \bigcup_{k \in S_{i,t}, k \neq k_1} H_{ik,t} = \emptyset$ , (\*). As  $\bar{v} \in \delta Z_i$ , we can find a  $*$ -sequence  $(\bar{v}_n) \subset *B(\bar{v}, \epsilon) \cap Z_i$ , with  $n \in * \mathcal{N}$ , such that  $|\bar{v}_n - \bar{v}| \rightarrow 0$  and  $l_{\bar{v}_n, \bar{p}} \cap H_{ik_1,t} = \bar{p}$ , (\*\*). As  $\bar{Z}_i$  is bounded, and using (\*\*), there exist points  $\bar{p}_n = l_{\bar{v}_n, \bar{p}} \cap * \bigcup_{k \in S_{i,t}, k \neq k_1} H_{ik,t}$ , and, without loss of generality, we can assume that  $\bar{p}_n = l_{\bar{v}_n, \bar{p}} \cap V_{ik,t}$  for some  $k \neq k_1$ . By (\*), we have that  $|\bar{p}_n - \bar{v}| \geq \epsilon$ , so the same is true for the  $*$ -limit  $\bar{p}' \in V_{ik,t}$ . Clearly,  $\bar{v} \in l_{\bar{p}', \bar{p}}$ , which contradicts the fact that  $\bar{v}$  is a vertex. Suppose that  $\bar{v} \in H_{ik_1,t} \cap H_{ik_2,t}$ , for  $k_1 \neq k_2$ , but  $\bar{v} \notin * \bigcup_{k \in S_{i,t}, k \notin \{k_1, k_2\}} H_{ik,t}$ , then, as  $\dim(V_{ik_1,t} \cap V_{ik_2,t}) = 1$ , we can choose  $\bar{p} \in V_{ik_1,t} \cap V_{ik_2,t}$ , with  $\bar{p} \neq \bar{v}$  again. Repeating the same argument, we can choose  $(\bar{v}_n)$  again such that  $l_{\bar{v}_n, \bar{p}} \cap H_{ik_1,t} = l_{\bar{v}_n, \bar{p}} \cap H_{ik_2,t} = \bar{p}$ , and obtain similarly a contradiction.  $\square$

**Definition 1.35.** Suppose that  $\{t_1, t_2\} \subset \mathcal{T}_\nu$ , with  $t_1 < t_2$ , and let  $Z_{i,t_1}$  be an alcove, corresponding to  $\bar{w}_i \in R_{t_1}$ , with  $\bar{w}_i$  finite. For  $t_1 \leq t \leq t_2$ , we define the shifted alcove  $Z_{i,t}$  to be the alcove corresponding to  $s_{\bar{p}_{\bar{w}_i}}(t)$ , which we later abbreviate to  $\bar{w}_{it}$ . We define the shift to be simple, if;

$$(i). T_{i,t} = \{s_{\bar{p}_{\bar{w}_i}}(t) : \bar{w} \in T_{i,t_1}\}$$

(ii). For every pair of vertices  $(\bar{v}, \bar{v}') \in V_{i,t_1} \times V_{i,t}$  there exist unique triples  $(k_1, k_2, k_3)$  and  $(k'_1, k'_2, k'_3)$  with;

$$\bar{v} = H_{ik_1,t_1} \cap H_{ik_2,t_1} \cap H_{ik_3,t_1}$$

$$\bar{v}' = H_{ik'_1,t_2} \cap H_{ik'_2,t} \cap H_{ik'_3,t}$$

(iii).  $\bar{v} \in V_{i,t_1}$  iff the intersection  $H_{ik_1,t} \cap H_{ik_2,t} \cap H_{ik_3,t}$  of the shifted planes defines a vertex  $\bar{v}' \in V_{i,t}$ .

If  $\bar{w}_{jt} \in T_{i,t}$ , we define the centre  $\bar{c}_{ij,t} = \frac{\bar{w}_{it} + \bar{w}_{jt}}{2}$ .

**Lemma 1.36.** We have that there exists  $S_{it_1} \in \mathcal{R}_{>0}$ , with;

$$|(\bar{v}_t)'|_{t_1}| \leq S_{it_1}$$

*Proof.* We have that  $\bar{c}_{ij,t} \in H_{ij,t}$  and for  $\bar{z} \in H_{ij,t}$ , we have;

$$\langle \bar{z} - \bar{c}_{ij,t}, \bar{w}_{jt} - \bar{w}_{it} \rangle = 0$$

so that with coordinates  $\bar{z} = (z_1, z_2, z_3)$ ,  $\bar{w}_{jt} = (w_{j1t}, w_{j2t}, w_{j3t})$ ,  $\bar{w}_{it} = (w_{i1t}, w_{i2t}, w_{i3t})$ ,  $\bar{c}_{ij,t} = (c_{ijt1}, c_{ijt2}, c_{ijt3}) = (\frac{w_{i1t} + w_{j1t}}{2}, \frac{w_{i2t} + w_{j2t}}{2}, \frac{w_{i3t} + w_{j3t}}{2})$ , the equation of the plane  $H_{ij,t}$  is given by;

$$\sum_{k=1}^3 (w_{jkt} - w_{ikt}) z_k = \sum_{k=1}^3 c_{ijtk} (w_{jkt} - w_{ikt})$$

By (ii), a vertex  $\bar{v}_t \in H_{i,t}$ , with coordinates  $\bar{v}_t = (v_{1t}, v_{2t}, v_{3t})$ , given by the intersection  $H_{ik_1,t} \cap H_{ik_2,t} \cap H_{ik_3,t}$ , is the intersection of the planes;

$$\sum_{k=1}^3 (w_{k_1kt} - w_{ikt}) v_{kt} = \sum_{k=1}^3 c_{ik_1tk} (w_{k_1kt} - w_{ikt})$$

$$= \sum_{k=1}^3 \frac{(w_{ikt} + w_{k_1kt})(w_{k_1kt} - w_{ikt})}{2} = \sum_{k=1}^3 \frac{w_{k_1kt}^2 - w_{ikt}^2}{2}$$

$$\sum_{k=1}^3 (w_{k_2kt} - w_{ikt}) v_{kt} = \sum_{k=1}^3 c_{ik_2tk} (w_{k_2kt} - w_{ikt})$$

$$= \sum_{k=1}^3 \frac{(w_{ikt} + w_{k_2kt})(w_{k_2kt} - w_{ikt})}{2} = \sum_{k=1}^3 \frac{w_{k_2kt}^2 - w_{ikt}^2}{2}$$

$$\sum_{k=1}^3 (w_{k_3kt} - w_{ikt}) v_{kt} = \sum_{k=1}^3 c_{ik_3tk} (w_{k_3kt} - w_{ikt})$$

$$= \sum_{k=1}^3 \frac{(w_{ikt} + w_{k_3kt})(w_{k_3kt} - w_{ikt})}{2} = \sum_{k=1}^3 \frac{w_{k_3kt}^2 - w_{ikt}^2}{2}$$

which we can write in matrix form, as;

$$A_{ik_1k_2k_3t}\bar{v}_t = \bar{d}_{ik_1k_2k_3t}$$

where, for  $1 \leq l \leq 3$ ,  $1 \leq k \leq 3$ ;

$$(A_{ik_1k_2k_3t})_{lk} = w_{k_lkt} - w_{ikt}, (\bar{d}_{ik_1k_2k_3t})_l = \sum_{k=1}^3 \frac{w_{k_lkt}^2 - w_{ikt}^2}{2}$$

By the intersection property, we have that;

$$\bar{v}_t = (A_{ik_1k_2k_3t})^{-1}\bar{d}_{ik_1k_2k_3t}$$

and, letting  $(\cdot)'_\nu$  denote the nonstandard derivative;

$$\begin{aligned} (\bar{v}_t)'_\nu|_{t_1} &= ((A_{ik_1k_2k_3t})^{-1}\bar{d}_{ik_1k_2k_3t})'_\nu|_{t_1} \\ &= \nu((A_{ik_1k_2k_3(t+\frac{1}{\nu})})^{-1}\bar{d}_{ik_1k_2k_3(t+\frac{1}{\nu})} - (A_{ik_1k_2k_3t})^{-1}\bar{d}_{ik_1k_2k_3t})|_{t_1} \\ &= \nu((A_{ik_1k_2k_3(t+\frac{1}{\nu})})^{-1} - (A_{ik_1k_2k_3t})^{-1})\bar{d}_{ik_1k_2k_3(t+\frac{1}{\nu})} \\ &\quad + (A_{ik_1k_2k_3t})^{-1}(\bar{d}_{ik_1k_2k_3(t+\frac{1}{\nu})} - \bar{d}_{ik_1k_2k_3t})|_{t_1} \\ &= (((A_{ik_1k_2k_3t})^{-1})'_\nu\bar{d}_{ik_1k_2k_3(t+\frac{1}{\nu})} + (A_{ik_1k_2k_3t})^{-1}(\bar{d}_{ik_1k_2k_3t})'_\nu)|_{t_1} \\ &= \left(\frac{\text{cof}(A_{ik_1k_2k_3t})^T}{\det(A_{ik_1k_2k_3t})}\right)'_\nu|_{t_1}\bar{d}_{ik_1k_2k_3(t+\frac{1}{\nu})} + \left(\frac{\text{cof}(A_{ik_1k_2k_3t_1})^T}{\det(A_{ik_1k_2k_3t_1})}\right)(\bar{d}_{ik_1k_2k_3t})'_\nu|_{t_1} \quad (E) \end{aligned}$$

We have, using Lemmas 1.7 and 1.9, that for  $1 \leq l \leq 3$ , there exists  $D_{t_1+\frac{1}{\nu},l} \in \mathcal{R}_{>0}$ , with;

$$\begin{aligned} |(\bar{d}_{ik_1k_2k_3(t_1+\frac{1}{\nu})})_l| &= \left|\sum_{k=1}^3 \frac{w_{k_lk(t_1+\frac{1}{\nu})}^2 - w_{ik(t_1+\frac{1}{\nu})}^2}{2}\right| \\ &= \frac{1}{2}||\sigma_{\bar{p}_{\bar{w}_i}}(t_1 + \frac{1}{\nu})|^2 - |\sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1 + \frac{1}{\nu})|^2|| \\ &= \frac{1}{2}(|\sigma_{\bar{p}_{\bar{w}_i}}(t_1 + \frac{1}{\nu})| + |\sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1 + \frac{1}{\nu})|)(|\sigma_{\bar{p}_{\bar{w}_i}}(t_1 + \frac{1}{\nu})| - |\sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1 + \frac{1}{\nu})|)| \\ &\leq D_{t_1+\frac{1}{\nu},l}||\sigma_{\bar{p}_{\bar{w}_i}}(t_1 + \frac{1}{\nu})| - |\sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1 + \frac{1}{\nu})|| \\ &\leq D_{t_1+\frac{1}{\nu},l}|\sigma_{\bar{p}_{\bar{w}_i}}(t_1 + \frac{1}{\nu}) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1 + \frac{1}{\nu})| \\ &\leq D_{t_1+\frac{1}{\nu},l}e^{A_{D_{t_1+\frac{1}{\nu}}}\circ(t_1+\frac{1}{\nu})}|\bar{p}_{\bar{w}_i} - \bar{p}_{\bar{w}_{k_l}}| \end{aligned}$$

Choosing  $\kappa$  sufficiently large, using Lemma 1.32, (†) of Lemma 1.19 and the fact that  $\bar{w}_{k_l,t} \in T_{i,t}$ , we can find  $\bar{z} \in H_{ik_l,t}$ , with  $|\bar{z} - \bar{w}_{i,t}| < \frac{C+1}{\sqrt{\eta}}$ . By the definition of  $H_{ik_l,t}$  and the fact that  $|c_{ik_l,t} - \bar{w}_{i,t}| \leq |\bar{z} - \bar{w}_{i,t}| < \frac{C+1}{\sqrt{\eta}}$ , we must have that  $|\bar{w}_{k_l,t} - \bar{w}_{i,t}| < \frac{2(C+1)}{\sqrt{\eta}}$ , and, by Lemma 1.27;

$$|\bar{p}_{\bar{w}_i} - \bar{p}_{\bar{w}_{k_l}}| < \frac{2(C+1)}{\sqrt{\eta}} 2(1 + \frac{C_{t_l}}{\nu})^{\lfloor \nu t_l \rfloor} < \frac{4(C+1)}{\sqrt{\eta}} (e^{C_{t_l} \circ t_l} + 1) = \frac{L_{ilt_l}}{\sqrt{\eta}}$$

It follows that there exists  $K_{ilt_l} \in \mathcal{R}_{>0}$ , with;

$$|(\bar{d}_{ik_1k_2k_3(t_1+\frac{1}{\nu})})_l| < D_{t_1+\frac{1}{\nu},l} e^{A_{D_{t_1+\frac{1}{\nu}} \circ (t_1+\frac{1}{\nu})}} \frac{4(C+1)}{\sqrt{\eta}} (e^{C_t \circ t} + 1) = \frac{K_{ilt_l}}{\sqrt{\eta}} \quad (\#)$$

We have, using Lemma 1.30, that;

$$\begin{aligned} & |((\bar{d}_{ik_1k_2k_3t})'_\nu)_l|_{t_1} = |(\sum_{k=1}^3 \frac{w_{k_lkt}^2 - w_{ikt}^2}{2})'_\nu|_{t_1}| \\ &= |\sum_{k=1}^3 \frac{(w_{ikt}^2)_\nu - (w_{k_lkt}^2)_\nu}{2}|_{t_1}| \\ &= |\nu \sum_{k=1}^3 \frac{(w_{ik(t+\frac{1}{\nu})}^2 - w_{ikt}^2) - (w_{k_lk(t+\frac{1}{\nu})}^2 - w_{k_lkt}^2)}{2}|_{t_1}| \\ &= |\nu \sum_{k=1}^3 \frac{(w_{ik(t+\frac{1}{\nu})} + w_{ikt})(w_{ik(t+\frac{1}{\nu})} - w_{ikt}) - (w_{k_lk(t+\frac{1}{\nu})} + w_{k_lkt})(w_{k_lk(t+\frac{1}{\nu})} - w_{k_lkt})}{2}|_{t_1}| \\ &= |\sum_{k=1}^3 \frac{(w_{ik(t_1+\frac{1}{\nu})} + w_{ikt_1})(w_{ikt_1})'_\nu|_{t_1} - (w_{k_lk(t_1+\frac{1}{\nu})} + w_{k_lkt_1})(w_{k_lkt_1})'_\nu|_{t_1}| \\ &\simeq |\sum_{k=1}^3 w_{ikt_1} (w_{ikt})'_\nu|_{t_1} - w_{k_lkt_1} (w_{k_lkt})'_\nu|_{t_1}| \\ &= |\sum_{k=1}^3 (w_{ikt_1} (w_{ikt})'_\nu|_{t_1} - w_{k_lkt_1} (w_{ikt})'_\nu|_{t_1} + w_{k_lkt_1} (w_{ikt})'_\nu|_{t_1} - w_{k_lkt_1} (w_{k_lkt})'_\nu|_{t_1})| \\ &= |(\sigma_{\bar{p}_{\bar{w}_i}}(t_1) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1))(\sigma_{\bar{p}_{\bar{w}_i}}(t))'_\nu|_{t_1} + (\sigma_{\bar{p}_{\bar{w}_i}}(t) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t))'_\nu|_{t_1} \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1)| \\ &\leq |\sigma_{\bar{p}_{\bar{w}_i}}(t_1) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1)| |(\sigma_{\bar{p}_{\bar{w}_i}}(t))'_\nu|_{t_1}| + |(\sigma_{\bar{p}_{\bar{w}_i}}(t) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t))'_\nu|_{t_1}| |\sigma_{\bar{p}_{\bar{w}_{k_l}}}(t_1)| \\ &\leq e^{A_{D_{t_1} \circ (t_1)}} |\bar{p}_{\bar{w}_i} - \bar{p}_{\bar{w}_{k_l}}| |(\sigma_{\bar{p}_{\bar{w}_i}}(t))'_\nu|_{t_1}| + G_{l,t_1} |(\sigma_{\bar{p}_{\bar{w}_i}}(t) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t))'_\nu|_{t_1}| \\ &\leq e^{A_{D_{t_1} \circ (t_1)}} \frac{L_{ilt_l}}{\sqrt{\eta}} |(\sigma_{\bar{p}_{\bar{w}_i}}(t))'_\nu|_{t_1}| + G_{l,t_1} |(\sigma_{\bar{p}_{\bar{w}_i}}(t) - \sigma_{\bar{p}_{\bar{w}_{k_l}}}(t))'_\nu|_{t_1}| \\ &\leq \frac{M_{ilt_l}}{\sqrt{\eta}} N_{it_l} + \sqrt{3} G_{l,t_1} \max_{\bar{z} \in *B(\bar{p}_{\bar{w}_{k_l}}, r_{il}+1)} (|(\frac{\partial^2 \bar{s}_{\eta,\nu}}{\partial x_i \partial t})_{\eta,\nu}|_{\bar{z},t_1}) |\bar{p}_{\bar{w}_i} - \bar{p}_{\bar{w}_{k_l}}| \\ &\leq \frac{M_{ilt_l}}{\sqrt{\eta}} N_{it_l} + \sqrt{3} G_{l,t_1} H_{lr_{il}} \frac{L_{ilt_l}}{\sqrt{\eta}} = \frac{K_{ilt_l}}{\sqrt{\eta}} \quad (\dagger) \end{aligned}$$

so that;

$$|((\bar{d}_{ik_1k_2k_3t})'_\nu)|_{t_1} \leq \sum_{l=1}^3 \frac{K_{ilt_l}}{\sqrt{\eta}} = \frac{P_{it_l}}{\sqrt{\eta}} \quad (\dagger\dagger)$$

where  $\{H_{lr_{il}}, M_{ilt_l}, K_{ilt_l}, N_{it_l}, G_{l,t_1}\} \subset \mathcal{R}_{>0}$ ,  $r_{il} = |\bar{p}_{\bar{w}_i} - \bar{p}_{\bar{w}_{k_l}}|$ .  
We have that;

$$\det(A_{ik_1k_2k_3t_1}) = \bar{\Delta}_{ik_1t_1} \cdot (\bar{\Delta}_{ik_2t_1} \times \bar{\Delta}_{ik_3t_1})$$

where  $\bar{\Delta}_{ik_jt} = \bar{w}_{k_j,t} - \bar{w}_{i,t}$ , for  $1 \leq j \leq 3$

We have that;

$$\begin{aligned} & \bar{\Delta}_{ik_2(t+\frac{1}{\nu})} \times \bar{\Delta}_{ik_3(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_2t} \times \bar{\Delta}_{ik_3t} \\ &= (\bar{\Delta}_{ik_2(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_2t}) \times \bar{\Delta}_{ik_3(t+\frac{1}{\nu})} + (\bar{\Delta}_{ik_3(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_3t}) \times \bar{\Delta}_{ik_2t} \quad (*) \end{aligned}$$

so that, using (\*);

$$\begin{aligned} & \det(A_{ik_1k_2k_3(t+\frac{1}{\nu})}) - \det(A_{ik_1k_2k_3t}) \\ &= \bar{\Delta}_{ik_1(t+\frac{1}{\nu})} \cdot (\bar{\Delta}_{ik_2(t+\frac{1}{\nu})} \times \bar{\Delta}_{ik_3(t+\frac{1}{\nu})}) - \bar{\Delta}_{ik_1t} \cdot (\bar{\Delta}_{ik_2t} \times \bar{\Delta}_{ik_3t}) \\ &= (\bar{\Delta}_{ik_1(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_1t}) \cdot (\bar{\Delta}_{ik_2(t+\frac{1}{\nu})} \times \bar{\Delta}_{ik_3(t+\frac{1}{\nu})}) - \bar{\Delta}_{ik_1t} \cdot (\bar{\Delta}_{ik_2(t+\frac{1}{\nu})} \times \\ & \quad \bar{\Delta}_{ik_3(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_2t} \times \bar{\Delta}_{ik_3t}) \\ &= (\bar{\Delta}_{ik_1(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_1t}) \cdot (\bar{\Delta}_{ik_2(t+\frac{1}{\nu})} \times \bar{\Delta}_{ik_3(t+\frac{1}{\nu})}) - \bar{\Delta}_{ik_1t} \cdot ((\bar{\Delta}_{ik_2(t+\frac{1}{\nu})} - \\ & \quad \bar{\Delta}_{ik_2t}) \times \bar{\Delta}_{ik_3(t+\frac{1}{\nu})} + (\bar{\Delta}_{ik_3(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_3t}) \times \bar{\Delta}_{ik_2t}) \quad (\dagger\dagger) \end{aligned}$$

By Lemma 1.30, and using the argument ( $\dagger$ ) above, we have, for  $1 \leq j \leq 3$ , that;

$$|\bar{\Delta}_{ik_jt}| \leq \frac{M_{ik_jt}}{\sqrt{\eta}}$$

$$|\bar{\Delta}_{ik_j(t+\frac{1}{\nu})} - \bar{\Delta}_{ik_jt}| = \frac{1}{\nu} |(\bar{s}_{\bar{p}_{\bar{w}_{k_j}}} - \bar{s}_{\bar{p}_{\bar{w}_i}})'_\nu|_t \leq \frac{L_{ik_jt}}{\nu\sqrt{\eta}}$$

where  $\{M_{ik_jt}, L_{ik_jt}\} \subset \mathcal{R}_{>0}$ . It follows, from ( $\dagger\dagger$ ), that;

$$\begin{aligned}
& |det(A_{ik_1k_2k_3(t+\frac{1}{\nu})}) - det(A_{ik_1k_2k_3t})| \\
& \leq |(\overline{\Delta}_{ik_1(t+\frac{1}{\nu})} - \overline{\Delta}_{ik_1t})| |\overline{\Delta}_{ik_2(t+\frac{1}{\nu})}| |\overline{\Delta}_{ik_3(t+\frac{1}{\nu})}| + |\overline{\Delta}_{ik_1t}| |(\overline{\Delta}_{ik_2(t+\frac{1}{\nu})} - \overline{\Delta}_{ik_2t})| |\overline{\Delta}_{ik_3(t+\frac{1}{\nu})}| \\
& + |\overline{\Delta}_{ik_1t}| |(\overline{\Delta}_{ik_3(t+\frac{1}{\nu})} - \overline{\Delta}_{ik_3t})| |\overline{\Delta}_{ik_2t}| \\
& \leq \frac{1}{\nu\eta^{\frac{3}{2}}} (M_{ik_2t}M_{ik_3t}Lik_1t + M_{ik_1t}M_{ik_3t}Lik_2t + M_{ik_1t}M_{ik_2t}Lik_3t) \\
& = \frac{O_{ik_1k_2k_3t}}{\nu\eta^{\frac{3}{2}}} (!)
\end{aligned}$$

where  $O_{ik_1k_2k_3t} \in \mathcal{R}_{>0}$ . In particular;

$$|(det(A_{ik_1k_2k_3t}))'_\nu| \leq \frac{O_{ik_1k_2k_3t}}{\eta^{\frac{3}{2}}}, (!!)$$

It follows that if  $t = \frac{s}{\nu} \leq t_1$ ;

$$\begin{aligned}
& |det(A_{ik_1k_2k_3(t_0+t)}) - det(A_{ik_1k_2k_3t_0})| \\
& = |^* \sum_{i=0}^{s-1} (det(A_{ik_1k_2k_3(t_0+\frac{i+1}{\nu})}) - det(A_{ik_1k_2k_3(t_0+\frac{i}{\nu})}))| \\
& \leq ^* \sum_{i=0}^{s-1} |det(A_{ik_1k_2k_3(t_0+\frac{i+1}{\nu})}) - det(A_{ik_1k_2k_3(t_0+\frac{i}{\nu})})| \\
& \leq ^* \sum_{i=0}^{s-1} \frac{O_{ik_1k_2k_3(t_0+\frac{i}{\nu})}}{\nu\eta^{\frac{3}{2}}} \\
& \leq \frac{sO_{ik_1k_2k_3t_0t_1}}{\nu\eta^{\frac{3}{2}}} = \frac{t\nu O_{ik_1k_2k_3t_0t_1}}{\nu\eta^{\frac{3}{2}}} = \frac{tO_{ik_1k_2k_3t_0t_1}}{\eta^{\frac{3}{2}}}
\end{aligned}$$

for  $max_{0 \leq i \leq s-1} (O_{ik_1k_2k_3(t_0+\frac{i}{\nu})}) \leq O_{ik_1k_2k_3t_0t_1}$  and  $O_{ik_1k_2k_3t_0t_1} \in \mathcal{R}_{>0}$ .

so that;

$$|det(A_{ik_1k_2k_3t}) - det(A_{ik_1k_2k_30})| \leq \frac{tO_{ik_1k_2k_30t_1}}{\eta^{\frac{3}{2}}}$$

...Correction.  $det(A_{ik_1k_2k_30}) \neq 0$ , with coordinates  $\frac{i}{\sqrt{\eta}}$ ,  $i \in \mathcal{Z}$ , implies that  $|det(A_{ik_1k_2k_30})| \geq \frac{M}{\eta^{\frac{3}{2}}}$ , where  $M \in \mathcal{Z}_{>0}$ . Same argument as below, find  $t_0 \in \mathcal{R}_{>0}$ , such that all triples  $(k_1k_2k_3)$  with  $det(A_{ik_1k_2k_30}) \neq 0$  remain vertices for  $0 \leq t < t_0$ . Continuity argument, to exclude contribution of vertices with  $det(A_{ik_1k_2k_3t}) = 0$ , for some  $0 \leq t < t_0$ ....

We have that  $|det(A_{ik_1k_2k_30})| = \frac{1}{\eta^{\frac{3}{2}}}$ , so that, choosing some  $t_3 \in \mathcal{R}_{>0}$  for  $0 < t \leq \min(\frac{1}{2O_{ik_1k_2k_30t_3}}, t_3)$ , with  $t \in \mathcal{R}_{>0}$ , we have that;

$$\frac{1}{2\eta^{\frac{3}{2}}} \leq |\det(A_{ik_1k_2k_3t})| \leq \frac{3}{2\eta^{\frac{3}{2}}} (\#\#\#)$$

We restrict attention to  $0 \leq t_1 < t_2 \leq \min(\frac{1}{2O_{ik_1k_2k_3}t_3}, t_3)$ .

We have, letting  $\|\cdot\|$  denote the maximum modulus matrix norm, and using the property, by transfer, from Lemma 1.37, that in  ${}^*\mathcal{R}^3$ ;

$$\begin{aligned} |A\bar{v}| &\leq 3\sqrt{3}\|A\|\|\bar{v}\| \\ \|(\text{cof}(A_{ik_1k_2k_3t_1})^T|_{t_1})\| &= \|\text{cof}(A_{ik_1k_2k_3t_1}|_{t_1})\| \\ &= \max_{1 \leq l, k \leq 3} (|\text{cof}(A_{ik_1k_2k_3t_1})_{lk}|) \\ &= \max_{1 \leq l, k \leq 3} (|(-1)^{l+k}(w_{k_l'k't} - w_{ik't})(w_{k_{\sigma(l')\sigma(k')t}} - w_{i\sigma(k')t}) - (w_{\sigma(k_l)kt} \\ &\quad - w_{ikt})(w_{k_l\sigma(k)t} - w_{i\sigma(k)t})|) \end{aligned}$$

where for  $1 \leq k, l \leq 3$ ,  $k' = \mu k (k \in A_k)$ ,  $A_k = \{1, 2, 3\} \setminus \{k\}$ ,  $l' = \mu l (l \in A_l)$ ,  $A_l = \{1, 2, 3\} \setminus \{l\}$ ,  $\sigma : A_k \rightarrow A_k$ ,  $\sigma \neq Id$ ,  $\sigma^2 = Id$ , and similarly for  $A_l$ .

We have, using the calculation in  $(\#)$ , for some  $C \in \mathcal{R}_{>0}$ ,  $1 \leq k, l \leq 3$ , that;

$$\begin{aligned} &|(-1)^{l+k}(w_{k_l'k't} - w_{ik't})(w_{k_{\sigma(l')\sigma(k')t}} - w_{i\sigma(k')t}) - (w_{\sigma(k_l)kt} - w_{ikt})(w_{k_l\sigma(k)t} \\ &\quad - w_{i\sigma(k)t})| \\ &\leq |w_{k_l'k't} - w_{ik't}| |w_{k_{\sigma(l')\sigma(k')t}} - w_{i\sigma(k')t}| + |w_{\sigma(k_l)kt} - w_{ikt}| |w_{k_l\sigma(k)t} - w_{i\sigma(k)t}| \\ &\leq 2\left(\frac{2\sqrt{3}(C+1)}{\sqrt{\eta}}\right)^2 = \frac{24(C+1)^2}{\eta} \end{aligned}$$

so that;

$$\|(\text{cof}(A_{ik_1k_2k_3t_1})^T|_{t_1})\| \leq \frac{24(C+1)^2}{\eta} (\#\#\#\#)$$

It follows from  $(\#\#)$ ,  $(\#\#\#)$ ,  $(\#\#\#\#)$ , and using Lemma 1.37, that;

$$\left| \frac{\text{cof}(A_{ik_1k_2k_3t_1})^T(\bar{d}_{ik_1k_2k_3t})_{\nu}|_{t_1}}{\det(A_{ik_1k_2k_3t_1})} \right| \leq \frac{2.24(C+1)^2 P_{it_1} \eta^{\frac{3}{2}}}{\eta\sqrt{\eta}}$$

$$= 48(C+1)^2 P_{it_1} = Q_{it_1} (D)$$

We have, using the chain rule, that;

$$\begin{aligned} & \left( \frac{\text{cof}(A_{ik_1 k_2 k_3 t})}{\det(A_{ik_1 k_2 k_3 t})} \right)'_{\nu} \Big|_{t_1} \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \\ &= \left( \left( \frac{1}{\det(A_{ik_1 k_2 k_3 t})} \right)'_{\nu} \Big|_{t_1} \text{cof}(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})^T + \frac{(\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1})} \right) \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \\ &= \left( - \frac{(\det(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1}) \det(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})} \right) \text{cof}(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})^T + \frac{(\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1})} \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \end{aligned} \quad (C)$$

By  $(\#)$ ,  $(\#\#\#)$ ,  $(\#\#\#\#)$ ,  $(!!)$ , and using Lemma 1.37 we have that;

$$\begin{aligned} & \left| \left( - \frac{(\det(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1}) \det(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})} \right) \text{cof}(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})^T \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \right| \\ & \leq 3\sqrt{3} \left| - \frac{(\det(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1}) \det(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})} \text{cof}(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})^T \right| \left| \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \right| \\ & = 3\sqrt{3} \left| \frac{(\det(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1}) \det(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})} \right| \left| \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \right| \left| \text{cof}(A_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})})^T \right| \\ & \leq 3\sqrt{3} \frac{O_{ik_1 k_2 k_3 t_1}}{\eta^{\frac{3}{2}}} 4\eta^3 \frac{K_{ilt_1}}{\sqrt{\eta}} \frac{24(C+1)^2}{\eta} \\ & = 288\sqrt{3}(C+1)^2 O_{ik_1 k_2 k_3 t_1} K_{ilt_1} (A) \end{aligned}$$

We have, using  $(\#)$ ,  $(\#\#\#)$ , Lemma 1.37, that;

$$\begin{aligned} & \left| \left( \frac{(\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1})} \right) \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \right| \\ & \leq 3\sqrt{3} \left| \frac{(\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1}}{\det(A_{ik_1 k_2 k_3 t_1})} \right| \left| \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \right| \\ & = 3\sqrt{3} \left| \frac{1}{\det(A_{ik_1 k_2 k_3 t_1})} \right| \left| \bar{d}_{ik_1 k_2 k_3(t_1 + \frac{1}{\nu})} \right| \left| (\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1} \right| \\ & \leq 3\sqrt{3} 2\eta^{\frac{3}{2}} K_{ilt_1} \frac{1}{\sqrt{\eta}} \left| (\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1} \right|, (B) \end{aligned}$$

We have that;

$$\begin{aligned} & \left| (\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1} \right| = \left| (\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1} \right| \\ & = \max_{1 \leq l, k \leq 3} \left( |(\text{cof}(A_{ik_1 k_2 k_3 t})'_{\nu})'_{\nu} \Big|_{t_1} \right) \\ & = \max_{1 \leq l, k \leq 3} \left( |(-1)^{l+k} (w_{k'l'k't} - w_{ik't}) (w_{k_{\sigma(l')\sigma(k')t} - w_{i\sigma(k')t}) - (w_{\sigma(k_l)kt} \right) \end{aligned}$$

$$-w_{ikt})(w_{k_l\sigma(k)t} - w_{i\sigma(k)t})]'|_{t_1}|$$

We have, using the product rule, the calculations in  $(\sharp)$  and  $(\dagger)$ , that for some  $\{C_l, D_l\} \subset \mathcal{R}_{>0}$ ,  $1 \leq k, l \leq 3$ ;

$$\begin{aligned} & | [(-1)^{l+k}(w_{k_l'k't} - w_{ik't})(w_{k_{\sigma(l')\sigma(k')t} - w_{i\sigma(k')t}) - (w_{\sigma(k_l)kt} - w_{ikt})(w_{k_l\sigma(k)t} \\ & - w_{i\sigma(k)t})]'|_{t_1}| \\ &= |(-1)^{l+k}(w_{k_l'k't} - w_{ik't})'|_{t_1}(w_{k_{\sigma(l')\sigma(k')(t_1+\frac{1}{\nu})} - w_{i\sigma(k')(t_1+\frac{1}{\nu})}) + (-1)^{l+k}(w_{k_l'k't_1} \\ & - w_{ik't_1})(w_{k_{\sigma(l')\sigma(k')t} - w_{i\sigma(k')t})'|_{t_1} - (w_{\sigma(k_l)kt} - w_{ikt})'|_{t_1}(w_{k_l\sigma(k)(t_1+\frac{1}{\nu})} - \\ & w_{i\sigma(k)(t_1+\frac{1}{\nu})}) - (w_{\sigma(k_l)kt_1} - w_{ikt_1})(w_{k_l\sigma(k)t} - w_{i\sigma(k)t})'|_{t_1}| \\ &\leq |(w_{k_l'k't} - w_{ik't})'|_{t_1}| |w_{k_{\sigma(l')\sigma(k')(t_1+\frac{1}{\nu})} - w_{i\sigma(k')(t_1+\frac{1}{\nu})}| + |w_{k_l'k't_1} \\ & - w_{ik't_1}| | (w_{k_{\sigma(l')\sigma(k')t} - w_{i\sigma(k')t})'|_{t_1}| + |(w_{\sigma(k_l)kt} - w_{ikt})'|_{t_1}| \\ & |w_{k_l\sigma(k)(t_1+\frac{1}{\nu})} - w_{i\sigma(k)(t_1+\frac{1}{\nu})}| + |w_{\sigma(k_l)kt_1} - w_{ikt_1}| | (w_{k_l\sigma(k)t} - w_{i\sigma(k)t})'|_{t_1}| \\ &\leq 4 \frac{C_1}{\sqrt{\eta}} \frac{D_1}{\sqrt{\eta}} = \frac{4C_1D_1}{\eta} \end{aligned}$$

so that;

$$|(\text{cof}(A_{ik_1k_2k_3t})^T)'_{\nu}|_{t_1}| \leq \frac{4C_1D_1}{\eta} (X)$$

and, from  $(B)$ ;

$$\begin{aligned} & |(\frac{\text{cof}(A_{ik_1k_2k_3t})^T)'_{\nu}|_{t_1}}{\det(A_{ik_1k_2k_3t_1})}) \bar{d}_{ik_1k_2k_3(t_1+\frac{1}{\nu})}| \leq 3\sqrt{3}2\eta^{\frac{3}{2}} \frac{K_{ilt_1}}{\sqrt{\eta}} \frac{4C_1D_1}{\eta} \\ &= 24\sqrt{3}K_{ilt_1}C_1D_1 \end{aligned}$$

so that, using  $(A)$ ,  $(C)$ ;

$$\begin{aligned} & |(\frac{\text{cof}(A_{ik_1k_2k_3t})^T)'_{\nu}|_{t_1}}{\det(A_{ik_1k_2k_3t})}) \bar{d}_{ik_1k_2k_3(t_1+\frac{1}{\nu})}| \\ &\leq 288\sqrt{3}(C+1)^2 O_{ik_1k_2k_3t_1} K_{ilt_1} + 24\sqrt{3}K_{ilt_1}C_1D_1 = R_{it_1} \end{aligned}$$

and from  $(D)$ ,  $(E)$ , that;

$$|(\bar{v}_t)'|_{t_1}| \leq Q_{it_1} + R_{it_1} = S_{it_1}$$

as required. □

**Lemma 1.37.** *Let  $\|\cdot\|$  be the maximum modulus norm on matrices and  $|\cdot|$  the usual vector modulus. Then, for  $\lambda \in \mathcal{R}$ ,  $\bar{v} \in \mathcal{R}^3$ , we have that;*

$$\|\lambda A\| = |\lambda| \|A\|$$

$$|A\bar{v}| \leq 3\sqrt{3} \|A\| |\bar{v}|$$

*Proof.* The first claim is clear. We have, for  $1 \leq i \leq 3$ , and  $\bar{w}_i = \sum_{j=1}^3 \beta_{ji} \bar{e}_j$ ,  $\bar{w} = \sum_{j=1}^3 \beta_j \bar{e}_j$ , where  $\beta_{ji} \in \{-1, 1\}$ ,  $\beta_j \in \{-1, 1\}$  noting that  $|\bar{w}_i| = |\bar{w}| = \sqrt{3}$ , that;

$$\begin{aligned} |A(\bar{e}_i)| &= (|a_{1i}|, |a_{2i}|, |a_{3i}|) \cdot \bar{w}_i \\ &\leq (|a_{1i}|, |a_{2i}|, |a_{3i}|) |\bar{w}_i| \\ &= \sqrt{3} (|a_{1i}|, |a_{2i}|, |a_{3i}|) \\ &= \sqrt{3} (\sum_{j=1}^3 |a_{ji}|^2)^{\frac{1}{2}} \\ &\leq \sqrt{3} (\sum_{j=1}^3 \|A\|^2)^{\frac{1}{2}} \\ &= 3 \|A\| \end{aligned}$$

so that;

$$\begin{aligned} |A\bar{v}| &= |\sum_{i=1}^3 v_i A(\bar{e}_i)| \\ &\leq \sum_{i=1}^3 |v_i| |A(\bar{e}_i)| \\ &\leq 3 \|A\| \sum_{i=1}^3 |v_i| \\ &= 3 \|A\| |\bar{v} \cdot \bar{w}| \\ &\leq 3 \|A\| |\bar{v}| |\bar{w}| \end{aligned}$$

$$= 3\sqrt{3}||A|||\bar{v}|$$

□

**Definition 1.38.** Let  $\bar{s}_{\eta,\nu} : \mathcal{R}_\eta^3 \times \mathcal{T}_\nu \rightarrow *R^3$  be as in Definition 2.29, and let  $\bar{s}_{\eta,\nu}(\bar{x}, t)$  be written in coordinates as  $(w_{\bar{x}1t}, w_{\bar{x}2t}, w_{\bar{x}3t})$ . We let  $\bar{d}_{\eta,\nu} : \mathcal{R}_\eta^{12} \times \mathcal{T}_\nu \rightarrow *R^3$  be defined by;

$$\bar{d}_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) = \left( \sum_{k=1}^3 \frac{w_{\bar{x}_1 kt}^2 - w_{\bar{x}_0 kt}^2}{2}, \sum_{k=1}^3 \frac{w_{\bar{x}_2 kt}^2 - w_{\bar{x}_0 kt}^2}{2}, \sum_{k=1}^3 \frac{w_{\bar{x}_3 kt}^2 - w_{\bar{x}_0 kt}^2}{2} \right)$$

and  $\bar{A}_{\eta,\nu} : \mathcal{R}_\eta^{12} \times \mathcal{T}_\nu \rightarrow *R^9$  be defined, for  $1 \leq l, k \leq 3$  by;

$$(\bar{A}_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t))_{lk} = w_{\bar{x}_l kt} - w_{\bar{x}_0 kt}$$

Let  $U \subset R_\eta^{12} \times \mathcal{T}_\nu$  be the  $*$ -finite set for which  $\det(\bar{A}_{\eta,\nu}) \neq 0$ , and  $V = \{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) : (\frac{[\bar{x}_0\sqrt{\eta}]}{\sqrt{\eta}}, \frac{[\bar{x}_1\sqrt{\eta}]}{\sqrt{\eta}}, \frac{[\bar{x}_2\sqrt{\eta}]}{\sqrt{\eta}}, \frac{[\bar{x}_3\sqrt{\eta}]}{\sqrt{\eta}}, \frac{[t\nu]}{\nu}) \in U\}$ . We define  $\bar{\Gamma}_{\eta,\nu} : V \rightarrow *R^3$  by;

$$\bar{\Gamma}_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) = (\bar{A}_{\eta,\nu})^{-1}\bar{d}_{\eta,\nu}$$

**Lemma 1.39.** Using similar notation to Lemma 2.30, we have, for  $1 \leq i, j \leq 3$ , that,

$$\left( \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_i \partial x_j} \right)_{\eta,\nu} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \bar{0}$$

*Proof.* Let  $\tau_{12} : \mathcal{R}_\eta^{12} \times \mathcal{T}_\nu \rightarrow \mathcal{R}_\eta^{12} \times \mathcal{T}_\nu$  be the elementary permutation which permutes  $\{\bar{x}_1, \bar{x}_2\}$ , and fixes  $\{\bar{x}_3, t\}$ . We have that;

$$(\bar{A}_{\eta,\nu}(\tau_{12}(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)))_{lk} = (\bar{A}_{\eta,\nu}(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t))_{lk}$$

$$= w_{\bar{x}_{\tau_{12}(l)} kt} - w_{\bar{x}_0 kt}$$

$$= (\bar{A}_{\eta,\nu}(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t))_{\tau_{12}(l)k}$$

and, swapping the first two rows;

$$\det(\bar{A}_{\eta,\nu})(\tau_{12}(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)) = \det(\bar{A}_{\eta,\nu})(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)$$

$$= -\det(\bar{A}_{\eta,\nu})(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)$$

It follows that  $\tau_{12}$  leaves  $U$  invariant. Moreover, as, for  $(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) \in U$ ,  $(\bar{A}_{\eta, \nu})^{-1} \bar{d}_{\eta, \nu}|_{\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t}$  defines the point intersection of the three hyperplanes  $\{H_{\bar{x}_0 \bar{x}_k t} : 1 \leq k \leq 3\}$ , which are permuted by  $\tau_{12}$ , we have that  $\Gamma_{\eta, \nu} \circ \tau_{12} = \Gamma_{\eta, \nu}$ . It follows that;

$$\begin{aligned} & \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu} \circ \tau_{12}|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \nu(\Gamma_{\eta, \nu}|_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, (t + \frac{1}{\nu}))} - \Gamma_{\eta, \nu}|_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)}) \\ & = \nu(\Gamma_{\eta, \nu} \circ \tau_{12}|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, (t + \frac{1}{\nu}))} - \Gamma_{\eta, \nu} \circ \tau_{12}|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)}) \\ & = (\Gamma_{\eta, \nu} \circ \tau_{12})'_{\nu}|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \\ & = \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu}|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \end{aligned}$$

so that  $(\frac{\partial \Gamma_{\eta, \nu}}{\partial t})_{\nu} \circ \tau_{12} = (\frac{\partial \Gamma_{\eta, \nu}}{\partial t})_{\nu}$  as well.

It follows that, for  $1 \leq j \leq 3$ ;

$$\begin{aligned} & \frac{\partial^2 \Gamma_{\eta, \nu}}{\partial x_{1j} \partial t}|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \frac{\partial}{\partial x_{1j}} \left( \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu} \circ \tau_{12} \right)|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \\ & = \sqrt{\eta} \left[ \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu} \circ \tau_{12} \right]|_{(\bar{x}_0, \frac{(\bar{x}_1 + \bar{e}_j)}{\sqrt{\eta}}, \bar{x}_2, \bar{x}_3, t)} - \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \Big] \\ & = \sqrt{\eta} \left[ \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu} \Big|_{(\bar{x}_0, \bar{x}_2, \frac{(\bar{x}_1 + \bar{e}_j)}{\sqrt{\eta}}, \bar{x}_3, t)} - \left(\frac{\partial \Gamma_{\eta, \nu}}{\partial t}\right)_{\nu} \Big|_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)} \right] \\ & = \left(\frac{\partial^2 \Gamma_{\eta, \nu}}{\partial x_{2j} \partial t}\right)_{\eta, \nu}|_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)} \\ & = \left(\frac{\partial^2 \Gamma_{\eta, \nu}}{\partial x_{2j} \partial t}\right)_{\eta, \nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \end{aligned}$$

so that  $\frac{\partial^2 \Gamma_{\eta, \nu}}{\partial x_{1j} \partial t} = \frac{\partial^2 \Gamma_{\eta, \nu}}{\partial x_{2j} \partial t} \circ \tau_{12}$ , (N).

As  $\bar{A}_{\eta, \nu} \bar{A}_{\eta, \nu}^{-1} = Id$ , we have that;

$$\frac{\partial(\bar{A}_{\eta, \nu} \bar{A}_{\eta, \nu}^{-1})}{\partial x_{1j}} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = O$$

so that;

$$\begin{aligned} & \left(\frac{\partial}{\partial x_{1j}}\right)_{\eta}(\bar{A}_{\eta, \nu}) \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \bar{A}_{\eta, \nu}^{-1} \Big|_{(\bar{x}_0, \bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_2, \bar{x}_3, t)} + \bar{A}_{\eta, \nu} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left(\frac{\partial}{\partial x_{1j}}\right)_{\eta}(\bar{A}_{\eta, \nu}^{-1}) \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \\ & = O \end{aligned}$$

and;

$$\begin{aligned}
\left(\frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j}}\right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} &= -\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)}\left(\frac{\partial \bar{A}_{\eta,\nu}}{\partial x_{1j}}\right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)}\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_2,\bar{x}_3,t)} \\
&= -\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)}\bar{B}_{\eta,\nu}|_{(\bar{x}_1,t)}\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_2,\bar{x}_3,t)} \quad (K)
\end{aligned}$$

where;

$$((\bar{B}_{\eta,\nu})|_{(\bar{x}_1,t)})_{1k} = \left(\frac{\partial \bar{s}}{\partial x_j}\right)|_{(\bar{x}_1,t)}_k$$

$$((\bar{B}_{\eta,\nu})|_{(\bar{x}_1,t)})_{lk} = 0, \text{ for } 2 \leq l \leq 3$$

Similarly;

$$\begin{aligned}
\left(\frac{\partial^2 \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j} \partial t}\right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \\
= \frac{\partial}{\partial t}(-\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)}\bar{B}_{\eta,\nu}|_{(\bar{x}_1,t)}\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_2,\bar{x}_3,t)})_\nu \quad (P)
\end{aligned}$$

We have, using the product rule, that;

$$\begin{aligned}
\left(\frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{1j} \partial t}\right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} &= \left(\frac{\partial^2 (\bar{A}_{\eta,\nu}^{-1} \bar{d}_{\eta,\nu})}{\partial x_{1j} \partial t}\right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \\
&= \frac{\partial}{\partial t} \left( \left( \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j}} \right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \bar{d}_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_2,\bar{x}_3,t)} + (\bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)}) \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial x_{1j}} \right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \right)_\nu \right) \\
&= \left( \frac{\partial^2 \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j} \partial t} \right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \bar{d}_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_2,\bar{x}_3,(t+\frac{1}{\nu}))} + \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j}} \right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial t} \right)_\nu|_{(\bar{x}_0,\bar{x}_1+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_2,\bar{x}_3,t)} \\
&\quad + \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial t} \right)_\nu|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial x_{1j}} \right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,(t+\frac{1}{\nu}))} + \bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \left( \frac{\partial^2 \bar{d}_{\eta,\nu}}{\partial x_{1j} \partial t} \right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \\
(L)
\end{aligned}$$

By symmetry, we have that;

$$\begin{aligned}
\left(\frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{2j} \partial t}\right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \\
= \left(\frac{\partial^2 \bar{A}_{\eta,\nu}^{-1}}{\partial x_{2j} \partial t}\right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \bar{d}_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_3,(t+\frac{1}{\nu}))} + \left(\frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{2j}}\right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \left(\frac{\partial \bar{d}_{\eta,\nu}}{\partial t}\right)_\nu|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2+\frac{\bar{e}_j}{\sqrt{\eta}},\bar{x}_3,t)} \\
+ \left(\frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial t}\right)_\nu|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \left(\frac{\partial \bar{d}_{\eta,\nu}}{\partial x_{2j}}\right)_\eta|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,(t+\frac{1}{\nu}))} + \bar{A}_{\eta,\nu}^{-1}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \left(\frac{\partial^2 \bar{d}_{\eta,\nu}}{\partial x_{2j} \partial t}\right)_{\eta,\nu}|_{(\bar{x}_0,\bar{x}_1,\bar{x}_2,\bar{x}_3,t)} \\
(M)
\end{aligned}$$

Combining (L), (M), (N), we obtain that;

$$\begin{aligned}
& \left( \frac{\partial^2 \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j} \partial t} \right)_{\eta,\nu} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \bar{d}_{\eta,\nu} \Big|_{(\bar{x}_0, \bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_2, \bar{x}_3, (t + \frac{1}{\nu}))} + \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{1j}} \right)_{\eta} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial t} \right)_{\nu} \Big|_{(\bar{x}_0, \bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_2, \bar{x}_3, t)} \\
& + \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial t} \right)_{\nu} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial x_{1j}} \right)_{\eta} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, (t + \frac{1}{\nu}))} + \bar{A}_{\eta,\nu}^{-1} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left( \frac{\partial^2 \bar{d}_{\eta,\nu}}{\partial x_{1j} \partial t} \right)_{\eta,\nu} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \\
& = \left( \frac{\partial^2 \bar{A}_{\eta,\nu}^{-1}}{\partial x_{2j} \partial t} \right)_{\eta,\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \bar{d}_{\eta,\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_3, (t + \frac{1}{\nu}))} + \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{2j}} \right)_{\eta} \circ \\
& \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial t} \right)_{\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_3, t)} \\
& + \left( \frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial t} \right)_{\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial x_{2j}} \right)_{\eta} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, (t + \frac{1}{\nu}))} + \bar{A}_{\eta,\nu}^{-1} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \left( \frac{\partial^2 \bar{d}_{\eta,\nu}}{\partial x_{2j} \partial t} \right)_{\eta,\nu} \circ \\
& \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \quad (O)
\end{aligned}$$

We have that;

$$\bar{d}_{\eta,\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \bar{d}_{\eta,\nu} \Big|_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3, t)} = \bar{d}_{\eta,\nu}^{\tau_{12}} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)}$$

where, for a vector  $\bar{v} \in {}^* \mathcal{R}^3$ ,  $\bar{v}^{\tau_{12}}$  is obtained from  $\bar{v}$  by swapping the first and second entries. A similar calculation shows that;

$$\left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial t} \right)_{\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \left( \frac{\partial \bar{d}_{\eta,\nu}}{\partial t} \right)_{\nu}^{\tau_{12}} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)}$$

We have that;

$$\bar{A}_{\eta,\nu} \circ \tau_{12} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \bar{A}_{\eta,\nu}^{\tau_{12}} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)}$$

where, for a matrix  $\bar{A} \in {}^* \mathcal{R}^9$ ,  $\bar{A}^{\tau_{12}}$  is obtained from  $\bar{A}$  by swapping the first and second rows.

It follows, as  $\bar{A}_{\eta,\nu} \bar{A}_{\eta,\nu}^{-1} = Id$ , that;

$$\begin{aligned}
I &= I \circ \tau_{12} = (\bar{A}_{\eta,\nu} \bar{A}_{\eta,\nu}^{-1}) \circ \tau_{12} \\
&= (\bar{A}_{\eta,\nu} \circ \tau_{12}) (\bar{A}_{\eta,\nu}^{-1} \circ \tau_{12}) \\
&= \bar{A}_{\eta,\nu}^{\tau_{12}} (\bar{A}_{\eta,\nu}^{-1} \circ \tau_{12})
\end{aligned}$$

so that;

$$\bar{A}_{\eta,\nu}^{-1} \circ \tau_{12} = (\bar{A}_{\eta,\nu}^{\tau_{12}})^{-1} = (\bar{A}_{\eta,\nu}^{-1})_{\tau_{12}}$$

where, for a matrix  $\bar{B} \in {}^*\mathcal{R}^9$ ,  $\bar{B}_{\tau_{12}}$  is obtained from  $\bar{B}$  by swapping the first and second columns.

Similarly, we have that;

$$\left(\frac{\partial \bar{A}_{\eta, \nu}^{-1}}{\partial t}\right)_{\nu} \circ \tau_{12} = \left(\left(\frac{\partial \bar{A}_{\eta, \nu}^{-1}}{\partial t}\right)_{\nu}\right)_{\tau_{12}}$$

We have, using Lemma 2.30, that;

$$\begin{aligned} & \frac{\partial}{\partial x_{1j}} (\bar{d}_{\eta, \nu})_{\eta} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3 t)} \\ &= \frac{\partial}{\partial x_{1j}} \left( \sum_{k=1}^3 \frac{w_{\bar{x}_1 kt}^2 - w_{\bar{x}_0 kt}^2}{2}, \sum_{k=1}^3 \frac{w_{\bar{x}_2 kt}^2 - w_{\bar{x}_0 kt}^2}{2}, \sum_{k=1}^3 \frac{w_{\bar{x}_3 kt}^2 - w_{\bar{x}_0 kt}^2}{2} \right) \\ &= \frac{\partial}{\partial x_j} \left( \sum_{k=1}^3 \frac{w_{\bar{x}_1 kt}^2}{2}, 0, 0 \right) |_{(\bar{x}_1, t)} \\ &= \left( \frac{\partial}{\partial x_j} \left( \frac{\sigma_{\bar{x}_1}(t) \cdot \sigma_{\bar{x}_1}(t)}{2} \right), 0, 0 \right) |_{\bar{x}_1} \\ &= \left( \frac{1}{2} \frac{\partial}{\partial x_j} (\sigma_{\bar{x}_1}(t)) \cdot (\sigma_{\bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}}(t) + \sigma_{\bar{x}_1}(t)), 0, 0 \right) |_{\bar{x}_1} \\ &\simeq \left( \left( \frac{\partial \bar{s}_{\eta, \nu}}{\partial x_j} \right)_{\eta} |_{(\bar{x}_1, t)} \cdot \bar{s}_{\eta, \nu} |_{(\bar{x}_1, t)}, 0, 0 \right) (F) \end{aligned}$$

Similarly, using (F), we have that;

$$\begin{aligned} & \frac{\partial^2}{\partial x_{1j} \partial t} (\bar{d}_{\eta, \nu})_{\eta, \nu} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3 t)} \\ &= \frac{\partial}{\partial t} \left( \left( \frac{1}{2} \frac{\partial}{\partial x_j} (\sigma_{\bar{x}_1}(t)) \cdot (\sigma_{\bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}}(t) + \sigma_{\bar{x}_1}(t + \frac{1}{\nu})), 0, 0 \right) |_{\bar{x}_1} \right) |_t \\ &\simeq \left( \left( \frac{\partial^2 \bar{s}_{\eta, \nu}}{\partial x_j \partial t} \right)_{\eta, \nu} |_{(\bar{x}_1, t)} \cdot \bar{s}_{\eta, \nu} |_{(\bar{x}_1, t)} + \left( \frac{\partial \bar{s}_{\eta, \nu}}{\partial x_j} \right)_{\eta} |_{(\bar{x}_1, t)} \cdot \left( \frac{\partial \bar{s}_{\eta, \nu}}{\partial t} \right)_{\nu} |_{(\bar{x}_1, t)}, 0, 0 \right) (G) \end{aligned}$$

It follows, by symmetry, that;

$$\begin{aligned} & \frac{\partial}{\partial x_{2j}} (\bar{d}_{\eta, \nu})_{\eta} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3 t)} \\ &= \left( 0, \frac{1}{2} \frac{\partial}{\partial x_j} (\sigma_{\bar{x}_2}(t)) \cdot (\sigma_{\bar{x}_2 + \frac{\bar{e}_j}{\sqrt{\eta}}}(t) + \sigma_{\bar{x}_2}(t)), 0 \right) |_{\bar{x}_2} \\ &\simeq \left( 0, \left( \frac{\partial \bar{s}_{\eta, \nu}}{\partial x_j} \right)_{\eta} |_{(\bar{x}_2, t)} \cdot \bar{s}_{\eta, \nu} |_{(\bar{x}_2, t)}, 0 \right) (F') \end{aligned}$$

and;

$$\frac{\partial}{\partial x_{2j}} (\bar{d}_{\eta, \nu})_{\eta} \circ \tau_{12} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3 t)}$$

$$\begin{aligned}
&= \frac{\partial}{\partial x_{2j}} (\bar{d}_{\eta,\nu}|)_{\eta} |_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3 t)} \\
&= (0, \frac{1}{2} \frac{\partial}{\partial x_j} (\sigma_{\bar{x}_1}(t)) \cdot (\sigma_{\bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}}(t) + \sigma_{\bar{x}_1}(t)), 0) |_{\bar{x}_1} \\
&\simeq (0, (\frac{\partial \bar{s}_{\eta,\nu}}{\partial x_j})_{\eta} |_{(\bar{x}_1, t)} \cdot \bar{s}_{\eta,\nu} |_{(\bar{x}_1, t)}, 0) (F'')
\end{aligned}$$

Similarly;

$$\begin{aligned}
&\frac{\partial^2}{\partial x_{2j} \partial t} (\bar{d}_{\eta,\nu}|)_{\eta,\nu} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3 t)} \\
&= \frac{\partial}{\partial t} ((0, \frac{1}{2} \frac{\partial}{\partial x_j} (\sigma_{\bar{x}_2}(t)) \cdot (\sigma_{\bar{x}_2 + \frac{\bar{e}_j}{\sqrt{\eta}}}(t) + \sigma_{\bar{x}_2}(t + \frac{1}{\nu})), 0) |_{\bar{x}_2}) |_t \\
&\simeq (0, ((\frac{\partial^2 \bar{s}_{\eta,\nu}}{\partial x_j \partial t})_{\eta,\nu} |_{(\bar{x}_2, t)} \cdot \bar{s}_{\eta,\nu} |_{(\bar{x}_2, t)} + (\frac{\partial \bar{s}_{\eta,\nu}}{\partial x_j})_{\eta} |_{(\bar{x}_2, t)} \cdot (\frac{\partial \bar{s}_{\eta,\nu}}{\partial t})_{\nu} |_{(\bar{x}_2, t)}), 0) (G')
\end{aligned}$$

and;

$$\begin{aligned}
&\frac{\partial^2}{\partial x_{2j} \partial t} (\bar{d}_{\eta,\nu}|)_{\eta,\nu} \circ \tau_{12} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3 t)} \\
&= \frac{\partial^2}{\partial x_{2j} \partial t} (\bar{d}_{\eta,\nu}|)_{\eta,\nu} |_{(\bar{x}_0, \bar{x}_2, \bar{x}_1, \bar{x}_3 t)} \\
&= \frac{\partial}{\partial t} ((0, \frac{1}{2} \frac{\partial}{\partial x_j} (\sigma_{\bar{x}_1}(t)) \cdot (\sigma_{\bar{x}_1 + \frac{\bar{e}_j}{\sqrt{\eta}}}(t) + \sigma_{\bar{x}_1}(t + \frac{1}{\nu})), 0) |_{\bar{x}_1}) |_t \\
&\simeq (0, ((\frac{\partial^2 \bar{s}_{\eta,\nu}}{\partial x_j \partial t})_{\eta,\nu} |_{(\bar{x}_1, t)} \cdot \bar{s}_{\eta,\nu} |_{(\bar{x}_1, t)} + (\frac{\partial \bar{s}_{\eta,\nu}}{\partial x_j})_{\eta} |_{(\bar{x}_1, t)} \cdot (\frac{\partial \bar{s}_{\eta,\nu}}{\partial t})_{\nu} |_{(\bar{x}_1, t)}), 0) (G'')
\end{aligned}$$

By symmetry, using  $(K)$ , we have that;

$$\begin{aligned}
&(\frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{2j}})_{\eta} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = -\bar{A}_{\eta,\nu}^{-1} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \bar{C}_{\eta,\nu} |_{(\bar{x}_2, t)} \bar{A}_{\eta,\nu}^{-1} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_3, t)} \\
&(K')
\end{aligned}$$

where;

$$\begin{aligned}
&((\bar{C}_{\eta,\nu}) |_{(\bar{x}_2, t)})_{2k} = (\frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_2, t)})_k \\
&((\bar{C}_{\eta,\nu}) |_{(\bar{x}_2, t)})_{lk} = 0, \text{ for } l = 1 \text{ or } l = 3
\end{aligned}$$

and;

$$\begin{aligned}
&(\frac{\partial \bar{A}_{\eta,\nu}^{-1}}{\partial x_{2j}})_{\eta} \circ \tau_{12} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \\
&= -\bar{A}_{\eta,\nu}^{-1} \circ \tau_{12} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} \bar{C}_{\eta,\nu} \circ \tau_{12} |_{(\bar{x}_2, t)} \bar{A}_{\eta,\nu}^{-1} \circ \tau_{12} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2 + \frac{\bar{e}_j}{\sqrt{\eta}}, \bar{x}_3, t)}
\end{aligned}$$







so that, using the product rule twice;

$$\begin{aligned}
& \frac{\partial}{\partial t} (\overline{A}_{\eta,\nu}^{-1}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \overline{B}_{\eta,\nu}|_{(\overline{x}_1,t)})_\nu [\overline{\Gamma}_{\eta,\nu}|_{(\overline{x}_0,\overline{x}_1+\frac{\overline{e}_j}{\sqrt{\eta}},\overline{x}_2,\overline{x}_3,(t+\frac{1}{\nu}))} - \overline{\Gamma}_{\eta,\nu}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2+\frac{\overline{e}_j}{\sqrt{\eta}},\overline{x}_3,(t+\frac{1}{\nu}))}] \\
& + \overline{A}_{\eta,\nu}^{-1}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \overline{B}_{\eta,\nu}|_{(\overline{x}_1,t)} [\frac{\partial}{\partial t} (\overline{\Gamma}_{\eta,\nu}|_{(\overline{x}_0,\overline{x}_1+\frac{\overline{e}_j}{\sqrt{\eta}},\overline{x}_2,\overline{x}_3,t)} - \overline{\Gamma}_{\eta,\nu}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2+\frac{\overline{e}_j}{\sqrt{\eta}},\overline{x}_3,t)})_\nu] \\
& = \frac{\partial}{\partial t} (\overline{A}_{\eta,\nu}^{-1}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \overline{B}_{\eta,\nu}|_{(\overline{x}_1,t)} [\overline{\Gamma}_{\eta,\nu}|_{(\overline{x}_0,\overline{x}_1+\frac{\overline{e}_j}{\sqrt{\eta}},\overline{x}_2,\overline{x}_3,t)} - \overline{\Gamma}_{\eta,\nu}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2+\frac{\overline{e}_j}{\sqrt{\eta}},\overline{x}_3,t)}])_\nu \\
& = \overline{0} \quad (O''''''''')
\end{aligned}$$

It follows that;

$$= \frac{\partial}{\partial t} (\overline{A}_{\eta,\nu}^{-1}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \overline{B}_{\eta,\nu}|_{(\overline{x}_1,t)} [(\frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{1j}} + \frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{2j}})|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)}])_\nu = \overline{0}$$

and, by the definition of  $\overline{B}_{\eta,\nu}|_{(\overline{x}_1,t)}$  and the fact that  $\overline{A}_{\eta,\nu}^{-1}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)}$  is invertible, we have that;

$$\frac{\partial}{\partial t} (\langle \frac{\partial \overline{s}}{\partial x_j}|_{(\overline{x}_1,t)}, (\frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{1j}} + \frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{2j}})|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \rangle)_\nu = 0, \quad (H)$$

so that;

$$\langle \frac{\partial \overline{s}}{\partial x_j}|_{(\overline{x}_1,t)}, (\frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{1j}} + \frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{2j}})|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \rangle = c \quad (I)$$

Applying  $\tau_{12}$ , we obtain that;

$$\langle \frac{\partial \overline{s}}{\partial x_j} \circ \tau_{12}|_{(\overline{x}_1,t)}, (\frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{1j}} + \frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{2j}}) \circ \tau_{12}|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \rangle = c$$

and, using the proof of (N), for  $1 \leq j \leq 3$ , we have that;

$$\frac{\partial \Gamma_{\eta,\nu}}{\partial x_{1j}} \circ \tau_{12} = \frac{\partial \Gamma_{\eta,\nu}}{\partial x_{2j}}$$

$$\frac{\partial \Gamma_{\eta,\nu}}{\partial x_{2j}} \circ \tau_{12} = \frac{\partial \Gamma_{\eta,\nu}}{\partial x_{1j}}$$

so that, for  $1 \leq j \leq 3$ ;

$$\langle \frac{\partial \overline{s}}{\partial x_j}|_{(\overline{x}_2,t)}, (\frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{1j}} + \frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{2j}})|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \rangle = c \quad (I')$$

Subtracting (I') from (I), we have, for  $1 \leq j \leq 3$ , that;

$$\langle \frac{\partial \overline{s}}{\partial x_j}|_{(\overline{x}_1,t)} - \frac{\partial \overline{s}}{\partial x_j}|_{(\overline{x}_2,t)}, (\frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{1j}} + \frac{\partial \overline{\Gamma}_{\eta,\nu}}{\partial x_{2j}})|_{(\overline{x}_0,\overline{x}_1,\overline{x}_2,\overline{x}_3,t)} \rangle = 0 \quad (I'')$$

Fix  $j$ , and choose  $(\bar{x}_{10}, \bar{x}_{20}, t_0)$  with;

$$\frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_{10}, t_0)} - \frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_{20}, t_0)} \neq \bar{0}$$

and a path  $\bar{x}_{20}(s)$ ,  $0 \leq s \leq s_0$ ,  $s_0 \in \mathcal{R}_{>0}$ , with  $\bar{x}_{20}(0) = \bar{x}_{20}$ , such that;

$$\frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_{10}, t_0)} - \frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_{20}(s), t_0+s)}$$

is fixed.

Let;

$$R(\bar{z}, s) = \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{1j}} + \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{2j}} \right) |_{(\bar{x}_0, \bar{x}_{10}, \bar{x}_{20}(s), \bar{z}, t_0+s)}$$

Then, for any  $\bar{x}_3$  with  $(\bar{x}_0, \bar{x}_{10}, \bar{x}_{20}, \bar{x}_3, t_0) \in U$ , on some  ${}^*B(\bar{x}_3, \epsilon)$ ,  $\epsilon \in \mathcal{R}_{>0}$ , we have that  $R$  maps  ${}^*B(\bar{x}_3, \epsilon) \times [0, s_0]$  into  $\mathcal{H}_{(\bar{x}_0, \bar{x}_{10}, \bar{x}_{20}, t_0)} \subset {}^*\mathcal{R}^3$ , where  $\mathcal{H}_{(\bar{x}_0, \bar{x}_{10}, \bar{x}_{20}, t_0)}$  is the  ${}^*$ -hyperplane defined by;

$$\mathcal{H}_{(\bar{x}_0, \bar{x}_{10}, \bar{x}_{20}, t_0)} = \{ \bar{v} : \langle \frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_{10}, t_0)} - \frac{\partial \bar{s}}{\partial x_j} |_{(\bar{x}_{20}, t_0)}, \bar{v} \rangle = 0 \}$$

As  $R|_{{}^*B(\bar{x}_3, \epsilon) \times [0, s_0]}$  is  ${}^*$ -analytic on  ${}^*B(\bar{x}_3, \epsilon) \times [0, s_0]$ , in particularly  ${}^*$ -smooth on  ${}^*B(\bar{x}_3, \epsilon) \times [0, s_0]$ , by the  ${}^*$ -inverse function theorem, there exists a smooth path  $\gamma_\epsilon : (-\delta, \delta) \rightarrow {}^*B(\bar{x}_3, \epsilon) \times [0, s_0]$ , such that  $R|_{Im(\gamma_\epsilon)} = \bar{f}$ . In particularly, transferring the property that the only real analytic functions with a limit point of zeros, are zero, we can conclude that  $R|_{{}^*B(\bar{x}_3, \epsilon) \times [0, s_0]} = \bar{f}$ , so that;

$$\left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{1j}} + \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{2j}} \right) |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2(s), \bar{z}, t_0+s)} = \bar{f}(\bar{x}_0, \bar{x}_1, \bar{x}_2, t_0), \bar{z} \in {}^*B(\bar{x}_3, \epsilon) \quad (I''')$$

In particular, differentiating with respect to  $t$ ;

$$\begin{aligned} & \left( \frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial t \partial x_{1j}} + \frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial t \partial x_{2j}} \right) |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2(s), \bar{x}_3, t_0+s)} \\ & + (\nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{1j}} \right)_1 \right) \bullet \bar{x}'_2(s), \nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{1j}} \right)_2 \right) \bullet \bar{x}'_2(s), \nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{1j}} \right)_3 \right) \bullet \bar{x}'_2(s)) |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2(s), \bar{x}_3, t_0+s)} \\ & + (\nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{2j}} \right)_1 \right) \bullet \bar{x}'_2(s), \nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{2j}} \right)_2 \right) \bullet \bar{x}'_2(s), \nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial x_{2j}} \right)_3 \right) \bullet \bar{x}'_2(s)) |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2(s), \bar{x}_3, t_0+s)} \\ & = \bar{0}, \quad (I'''' ) \end{aligned}$$

From  $(I''')$ , differentiating with respect to  $\bar{x}_3$ , taking  $s = 0$ , we have for  $1 \leq k \leq 3$ , that;

$$\left( \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{3k} \partial x_{1j}} + \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{3k} \partial x_{2j}} \right) \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t_0)} = \bar{0}$$

so that, taking the limit as  $\bar{x}_2 \rightarrow \bar{x}_1$ , we have, as below, that, for  $1 \leq k \leq 3$ ;

$$\frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{3k} \partial x_{1j}} = \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{3k} \partial x_{2j}} = \bar{0}$$

By symmetry we can conclude that, for  $1 \leq j \leq 3, 1 \leq k \leq 3, 1 \leq a < b \leq 3$ ;

$$\frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial x_{ak} \partial x_{bj}} = \bar{0}$$

By a similar argument to  $(N)$ , we have that;

$$\frac{\partial \Gamma_{\eta,\nu}}{\partial x_{1j}} = \frac{\partial \Gamma_{\eta,\nu}}{\partial x_{2j}} \circ \tau_{12}$$

so that, for  $1 \leq j \leq 3, 1 \leq k \leq 3$ ;

$$\frac{\partial^2 \Gamma_{\eta,\nu}}{\partial x_{1k} \partial x_{1j}} = \frac{\partial^2 \Gamma_{\eta,\nu}}{\partial x_{1k} \partial x_{2j}} \circ \tau_{12} = \bar{0}$$

and, by symmetry, we can conclude that, for  $1 \leq a \leq 3, 1 \leq j \leq 3, 1 \leq k \leq 3$ ;

$$\frac{\partial^2 \Gamma_{\eta,\nu}}{\partial x_{ak} \partial x_{aj}} = \bar{0}$$

as well. In particular, for  $1 \leq d \leq 3$ ;

$$(\nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta,\nu}}{\partial x_{1j}} \right)_d \right)) = (\nabla_{\bar{x}_2} \left( \left( \frac{\partial \bar{\Gamma}_{\eta,\nu}}{\partial x_{2j}} \right)_d \right)) = \bar{0} \quad (I''''')$$

From  $(I''''')$ ,  $(I''''''')$ , we obtain that;

$$\left( \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial t \partial x_{1j}} + \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial t \partial x_{2j}} \right) \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2(s), \bar{x}_3, t_0+s)} = \bar{0}, \quad (I''''''')$$

Using  $(N)$ , taking  $s = 0$ , we have that;

$$\frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial t \partial x_{1j}} \circ \tau_{1j} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t_0)} = - \frac{\partial^2 \bar{\Gamma}_{\eta,\nu}}{\partial t \partial x_{1j}} \Big|_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t_0)}$$

Taking the limit as  $\bar{x}_1 \rightarrow \bar{x}_2$ , we obtain that;

$$\frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial t \partial x_{1j}} \Big|_{\Delta} = \bar{0}$$

where  $\Delta = \{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) \in {}^* \mathcal{R}^{12} \times \mathcal{T}_\nu, \bar{x}_1 = \bar{x}_2\}$

and we extend  $\frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial t \partial x_{1j}}$   $*$ -analytically to  $\Delta$ . By the limiting property of  $*$ -analytic functions, we obtain that;

$$\frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial t \partial x_{1j}} = \bar{0}$$

and, by symmetry, we have, for  $1 \leq k \leq 3$ ,  $1 \leq j \leq 3$ , that;

$$\frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial t \partial x_{kj}} = \bar{0}$$

□

**Definition 1.40.** *We say that two vertices  $\{\bar{v}_1, \bar{v}_2\} \subset V_{it}$ , in the alcove  $Z_{i,t}$ , defined by  $\bar{x}_0$ , are connected if they are defined by relevant points  $\{(s_{\bar{x}_0}(t), s_{\bar{x}_{11}}(t), s_{\bar{x}_{21}}(t), s_{\bar{x}_{31}}(t))\}$  and  $\{(s_{\bar{x}_0}(t), s_{\bar{x}_{12}}(t), s_{\bar{x}_{22}}(t), s_{\bar{x}_{32}}(t))\}$ , for  $t \in (t_0, t_0 + \epsilon)$ , with  $\epsilon \in \mathcal{R}_{>0}$ , with the tuples  $(\bar{x}_0, \bar{x}_{11}, \bar{x}_{21}, \bar{x}_{31}, t)$  and  $(\bar{x}_0, \bar{x}_{12}, \bar{x}_{22}, \bar{x}_{32}, t)$  lying in the same  $*$ -connected component of  $U$ .*

Given an alcove  $Z_{i,t}$ , defined by  $\bar{x}_0$ , with vertices  $V_{i,t} = \{\bar{v}_i : 1 \leq i \leq \kappa(\bar{x}_0)\}$ , we have that, for any two distinct vertices,  $\{\bar{v}_1, \bar{v}_2\} \subset V_{i,t}$ , the edges  $\bar{e}_{\bar{x}_0, \bar{v}_1}$  and  $\bar{e}_{\bar{x}_0, \bar{v}_2}$  intersect at  $\bar{x}_0$ . Otherwise, without loss of generality,  $\bar{v}_2 \in \bar{e}_{\bar{x}_0, \bar{v}_2}$ , contradicting the definition of a vertex. For three distinct vertices  $\{\bar{v}_1, \bar{v}_2, \bar{v}_3\}$ , we define  $T_{\bar{x}_0, \bar{v}_1, \bar{v}_2, \bar{v}_3}$  to be the  $*$ -convex hull of  $\{\bar{x}_0, \bar{v}_1, \bar{v}_2, \bar{v}_3\}$ , so that  $T_{\bar{x}_0, \bar{v}_1, \bar{v}_2, \bar{v}_3} \subset Z_{i,t}$ . Let  $p \in Z_{i,t}$ , with  $p \neq \bar{x}_0$ , then  $l_{\bar{x}_0, p}$  intersects  $\delta Z_{i,t}$  at  $q$ , with  $q \in V_{ik,t}$ , a  $*$ -connected component of  $\delta Z_{i,t}$ ,  $\dim(V_{ik,t}) = 2$ . If  $\bar{v} \in V_{ik,t}$  is an extreme point of  $V_{ik,t}$ , then it is a vertex, otherwise it would lie on a line  $l_{\bar{y}_1, \bar{y}_2}$ , with  $\dim(l_{\bar{y}_1, \bar{y}_2} \cap V_{ik,t}) = 1$ ,  $\{\bar{y}_1, \bar{y}_2\} \subset Z_{i,t}$ , and  $\{\bar{y}_1, \bar{y}_2\}$  distinct from  $\bar{v}$ , contradicting the fact that  $\bar{v} \in V_{ik,t}$ . As  $\dim(V_{ik,t}) = 2$ , it has at least 3 vertices, and we can pick  $\{\bar{v}_1, \bar{v}_2, \bar{v}_3\}$ , such that  $q \in \mathcal{H}_{\bar{v}_1, \bar{v}_2, \bar{v}_3}$ . In particular, it follows that;

$$Z_{i,t} = {}^* \bigcup_{(\bar{v}_1, \bar{v}_2, \bar{v}_3) \in V_{i,t}^3, \bar{v}_i \neq \bar{v}_j, 1 \leq i < j \leq 3} T_{\bar{x}_0, \bar{v}_1, \bar{v}_2, \bar{v}_3}$$

Let  $R_{\bar{x}_0}$  be the  $*$ -planar graph, with vertices  $\{a_i : 1 \leq i \leq \kappa(\bar{x}_0)\}$ . and an edge  $e_{ij}$  assigned if  $\bar{e}_{\bar{v}_i, \bar{v}_j} \subset \delta Z_{i,t}$ . Clearly,  $R_{\bar{x}_0}$  is 4-connected, so, by [9], there exists a  $*$ -Hamiltonian cycle  $C$ , with  $C(0) = C(\kappa(\bar{x}_0))$ .

Then, clearly;

$$Z_{i,t} = * \bigcup_{0 \leq j \leq \kappa(\bar{x}_0) - 2} T_{\bar{x}_0, \bar{v}_{C(j)}, \bar{v}_{C(j+1)}, \bar{v}_{C(j+2)}}$$

and, moreover, for  $0 \leq j < k \leq \kappa(\bar{x}_0) - 2$ ;

$$\dim(T_{\bar{x}_0, \bar{v}_{C(j)}, \bar{v}_{C(j+1)}, \bar{v}_{C(j+2)}} \cap T_{\bar{x}_0, \bar{v}_{C(k)}, \bar{v}_{C(k+1)}, \bar{v}_{C(k+2)}}) \leq 2$$

**Lemma 1.41.** *We have, for two connected vertices  $\{\bar{v}_1, \bar{v}_2\} \subset V_{it}$ , that, for  $t \in (t_0, t_0 + \epsilon)$ ;*

$$(\bar{v}_1)'_{\nu}(t) = (\bar{v}_2)'_{\nu}(t)$$

and the velocities are all finite. Moreover, if;

$$\bar{e}_{\bar{v}_1, \bar{v}_2}(t) = \bar{v}_1(t) - \bar{v}_2(t)$$

and;

$$l_{\bar{v}_1, \bar{v}_2}(t) = |\bar{e}_{\bar{v}_1, \bar{v}_2}(t)|$$

is the length of the edge connecting  $\bar{v}_1$  and  $\bar{v}_2$ , then  $(l_{\bar{v}_1, \bar{v}_2})'_{\nu}(t) = 0$ , for  $t \in (t_0, t_0 + \epsilon)$ .

We have, for a vertex  $\bar{v}$ , that there exists  $\{L_{\bar{v}, \bar{x}_0, t}, K_{\bar{v}, \bar{x}_0, t}\} \subset \mathcal{R}_{>0}$ , with;

$$|(\bar{v})_{\nu}(t) - (s_{\bar{x}_0})_{\nu}(t)| \leq \frac{L_{\bar{v}, \bar{x}_0, t}}{\sqrt{\eta}}$$

$$|(\bar{v})'_{\nu}(t) - (s_{\bar{x}_0})'_{\nu}(t)| \leq \frac{K_{\bar{v}, \bar{x}_0, t}}{\sqrt{\eta}}$$

*Proof.* By Lemma 1.39, we have that, for the alcove  $Z_{i,t}$  defined by  $\bar{x}_0$ , that, for  $1 \leq i \leq 3$ ,  $1 \leq j \leq 3$ ;

$$\left(\frac{\partial^2 \bar{\Gamma}_{\eta, \nu}}{\partial x_i \partial t}\right)_{\eta, \nu} |_{(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t)} = \bar{0}$$

so that  $(\frac{\partial \bar{\Gamma}_{\eta, \nu}}{\partial t})_{\nu}$  is constant on the  $*$ -connected components of  $U$ . and;

$$\bar{v}_1(t) = \bar{\Gamma}_{\eta, \nu} |_{(\bar{x}_0, \bar{x}_{11}, \bar{x}_{21}, \bar{x}_{31}, t)}, \bar{v}_2(t) = \bar{\Gamma}_{\eta, \nu} |_{(\bar{x}_0, \bar{x}_{12}, \bar{x}_{22}, \bar{x}_{32}, t)}$$

where for  $1 \leq i \leq 3$ ;

$$s_{\bar{x}_{i1}}(t) = \bar{w}_{i1,t}, \quad s_{\bar{x}_{i2}}(t) = \bar{w}_{i2,t}$$

$$\bar{v}_1(t) = \mathcal{H}_{k_{11}t} \cap \mathcal{H}_{k_{21}t} \cap \mathcal{H}_{k_{31}t}, \quad \bar{v}_2(t) = \mathcal{H}_{k_{12}t} \cap \mathcal{H}_{k_{22}t} \cap \mathcal{H}_{k_{32}t}$$

It follows that, for  $t \in (t_0, t_0 + \epsilon)$ ;

$$\begin{aligned} |(\bar{v}_1)'_{\nu}(t) - (\bar{v}_2)'_{\nu}(t)| &= \left| \frac{\partial \bar{\Gamma}_{\eta,\nu}}{\partial t} |(\bar{x}_0, \bar{x}_{11}, \bar{x}_{21}, \bar{x}_{31}, t) - \frac{\partial \bar{\Gamma}_{\eta,\nu}}{\partial t} |(\bar{x}_0, \bar{x}_{12}, \bar{x}_{22}, \bar{x}_{32}, t) \right| \\ &= 0 \end{aligned}$$

so that  $(\bar{v}_1)'_{\nu}(t) = (\bar{v}_2)'_{\nu}(t)$ , (\*). By Lemma 1.36, all the velocities are finite. We have, for  $t \in (t_0, t_0 + \epsilon)$ , that;

$$\begin{aligned} (\bar{e}_{\bar{v}_1, \bar{v}_2})'_{\nu}(t) &= (\bar{v}_1)'_{\nu}(t) - (\bar{v}_2)'_{\nu}(t) \\ &= \bar{0} \end{aligned}$$

so that  $\bar{e}_{\bar{v}_1, \bar{v}_2}$  is constant on  $(t_0, t_0 + \epsilon)$ . In particular,  $l_{\bar{v}_1, \bar{v}_2}$  is constant on  $(t_0, t_0 + \epsilon)$ , and  $(l_{\bar{v}_1, \bar{v}_2})'_{\nu}(t) = 0$ , for  $t \in (t_0, t_0 + \epsilon)$ .

For the penultimate claim, we have that, for  $(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) \in U$ , using the calculations (###), (####), from Lemma 1.36, that;

$$\begin{aligned} |(\bar{v})_{\nu}(t) - (s_{\bar{x}_0})_{\nu}(t)| &= |\Gamma_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) - s_{\bar{x}_0,t}| \\ &= |\bar{A}_{\eta,\nu,t}^{-1} \bar{d}_{\eta,\nu,t} - s_{\bar{x}_0,t}| \\ &= |\bar{A}_{\eta,\nu,t}^{-1} (\bar{d}_{\eta,\nu,t} - \bar{A}_{\eta,\nu,t} s_{\bar{x}_0,t})| \\ &\leq 3\sqrt{3} |\bar{A}_{\eta,\nu,t}^{-1}| |\bar{d}_{\eta,\nu,t} - \bar{A}_{\eta,\nu,t} s_{\bar{x}_0,t}|, \leq 3\sqrt{3} (48(C+1)^2) \eta^{\frac{1}{2}} |\bar{d}_{\eta,\nu,t} - \bar{A}_{\eta,\nu,t} s_{\bar{x}_0,t}| \\ &= 144\sqrt{3} (C+1)^2 \eta^{\frac{1}{2}} |\bar{d}_{\eta,\nu,t} - \bar{A}_{\eta,\nu,t} s_{\bar{x}_0,t}| \quad (\dagger) \end{aligned}$$

We have that, for  $1 \leq j \leq 3$ ;

$$\begin{aligned} &|\bar{d}_{j,\eta,\nu,t} - (\bar{A}_{\eta,\nu,t} s_{\bar{x}_0,t})_j| \\ &= |s_{\bar{x}_j,t}|^2 - |s_{\bar{x}_0,t}|^2 - (s_{\bar{x}_j,t} - s_{\bar{x}_0,t}) \cdot s_{\bar{x}_0,t} \end{aligned}$$

$$= |s_{\bar{x}_j,t} \cdot (s_{\bar{x}_j,t} - s_{\bar{x}_0,t})|$$

so that, from (†);

$$\begin{aligned} & |\Gamma_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) - s_{\bar{x}_0,t}| \\ & \leq 144\sqrt{3}(C+1)^2\eta^{\frac{1}{2}} \sum_{j=1}^3 |s_{\bar{x}_j,t} \cdot (s_{\bar{x}_j,t} - s_{\bar{x}_0,t})| \quad (\dagger\dagger) \end{aligned}$$

For a given  $t_0$ , using a translation by  $s_{\bar{x}_0,t_0}$ , it is easy to see that;

$$|\Gamma_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) - s_{\bar{x}_0,t}| = |\Gamma'_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t) - s'_{\bar{x}_0,t}|$$

where the hashed trajectory  $s_{\bar{x}_0,t}^\# = s_{\bar{x}_0,t} - s_{\bar{x}_0,t_0}$ , and, in the definition of  $\Gamma'$ , we replace the trajectories  $\{s_{\bar{x}_j,t} : 0 \leq j \leq 4\}$  by  $\{s_{\bar{x}_j,t}^\# : 0 \leq j \leq 4\}$ , with  $s_{\bar{x}_j,t}^\# = s_{\bar{x}_j,t} - s_{\bar{x}_0,t_0}$ , for  $1 \leq j \leq 3$ . We can, therefore, assume that  $s_{\bar{x}_0,t_0} = \bar{0}$ , and, using (†) of Lemma 1.19 and Lemma 1.32,  $|s_{\bar{x}_j,t_0}| \leq \frac{R_{\bar{v},\bar{x}_0,t_0}}{\sqrt{\eta}}$ ,  $1 \leq j \leq 3$ , where  $R_{\bar{v},\bar{x}_0,t_0} \in \mathcal{R}_{>0}$ , to obtain, from (††), that;

$$\begin{aligned} & |\Gamma_{\eta,\nu}(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t_0) - s_{\bar{x}_0,t_0}| \\ & \leq 144\sqrt{3}(C+1)^2\eta^{\frac{1}{2}} \sum_{j=1}^3 |s_{\bar{x}_j,t_0}|^2 \\ & \leq 3.144\sqrt{3}(C+1)^2\eta^{\frac{1}{2}} \sum_{j=1}^3 \frac{R_{\bar{v},\bar{x}_0,t_0}^2}{\eta} \\ & = \frac{L_{\bar{v},\bar{x}_0,t_0}}{\eta^{\frac{1}{2}}} \end{aligned}$$

As  $t_0$  was arbitrary, we obtain the result. For the final claim, suppose that  $s_{\bar{x}_0,t_0} = s'_{\bar{x}_0,t_0} = \bar{0}$ , so that  $s_{\bar{x}_0,t_0+\frac{1}{\nu}} = \bar{0}$  and  $(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t_0) \in U$ . Using the calculations (###), (####), (!!), (X) from Lemma 1.36 we have that;

$$\begin{aligned} & |(\bar{v})'_\nu(t_0) - (s_{\bar{x}_0})'_\nu(t_0)| = |(\frac{\partial \Gamma_{\eta,\nu}}{\partial t})_\nu(\bar{x}_0, \bar{x}_1, \bar{x}_2, \bar{x}_3, t_0) - s'_{\bar{x}_0,t_0}| \\ & = |\frac{\partial}{\partial t}(\bar{A}_{\eta,\nu,t_0}^{-1} \bar{d}_{\eta,\nu,t_0} - \bar{s}_{\eta,\nu,\bar{x}_0,t_0})_\nu| \\ & = |\frac{\partial}{\partial t}[\bar{A}_{\eta,\nu,t_0}^{-1} (\bar{d}_{\eta,\nu,t_0} - \bar{A}_{\eta,\nu,t_0} \bar{s}_{\eta,\nu,\bar{x}_0,t_0})]_\nu| \\ & = |(\frac{\partial \bar{A}_{\eta,\nu,t_0}^{-1}}{\partial t})_\nu (\bar{d}_{\eta,\nu,t_0+\frac{1}{\nu}} - \bar{A}_{\eta,\nu,t_0+\frac{1}{\nu}} \bar{s}_{\eta,\nu,\bar{x}_0,t_0+\frac{1}{\nu}}) + \bar{A}_{\eta,\nu,t_0}^{-1} (\bar{d}_{\eta,\nu,t_0} - \bar{A}_{\eta,\nu,t_0} \bar{s}_{\eta,\nu,\bar{x}_0,t_0})'_\nu| \end{aligned}$$

$$\begin{aligned}
&\leq 3\sqrt{3} \left\| \left( \frac{\partial \bar{A}_{\eta, \nu, t_0}^{-1}}{\partial t} \right)_\nu \right\| \left| \bar{d}_{\eta, \nu, t_0 + \frac{1}{\nu}} - \bar{A}_{\eta, \nu, t_0 + \frac{1}{\nu}} \bar{s}_{\eta, \nu, \bar{x}_0, t_0 + \frac{1}{\nu}} \right| \\
&+ 3\sqrt{3} \left\| \bar{A}_{\eta, \nu, t_0}^{-1} \right\| \left| (\bar{d}_{\eta, \nu, t_0} - \bar{A}_{\eta, \nu, t_0} \bar{s}_{\eta, \nu, \bar{x}_0, t_0})'_\nu \right| \\
&\leq 3\sqrt{3} D \sqrt{\eta} \sum_{j=1}^3 |s_{\bar{x}_j, t_0} \cdot (s_{\bar{x}_j, t_0 + \frac{1}{\nu}} - s_{\bar{x}_0, t_0 + \frac{1}{\nu}})| \\
&+ 3\sqrt{3} E \sqrt{\eta} \left( \sum_{j=1}^3 |s'_{\bar{x}_j, t_0} \cdot (s_{\bar{x}_j, t_0 + \frac{1}{\nu}} - s_{\bar{x}_0, t_0 + \frac{1}{\nu}})| + \sum_{j=1}^3 |s_{\bar{x}_j, t_0} \cdot (s'_{\bar{x}_j, t_0} - s'_{\bar{x}_0, t_0})| \right) \\
&= 3\sqrt{3} D \sqrt{\eta} \sum_{j=1}^3 |s_{\bar{x}_j, t_0} \cdot s_{\bar{x}_j, t_0 + \frac{1}{\nu}}| \\
&+ 3\sqrt{3} E \sqrt{\eta} \left( \sum_{j=1}^3 |s'_{\bar{x}_j, t_0} \cdot s_{\bar{x}_j, t_0 + \frac{1}{\nu}}| + \sum_{j=1}^3 |s_{\bar{x}_j, t_0} \cdot s'_{\bar{x}_j, t_0}| \right) \\
&\leq 3\sqrt{3} D \sqrt{\eta} \sum_{j=1}^3 |s_{\bar{x}_j, t_0}| |s_{\bar{x}_j, t_0 + \frac{1}{\nu}}| \\
&+ 3\sqrt{3} E \sqrt{\eta} \left( \sum_{j=1}^3 |s'_{\bar{x}_j, t_0}| |s_{\bar{x}_j, t_0 + \frac{1}{\nu}}| + \sum_{j=1}^3 |s_{\bar{x}_j, t_0}| |s'_{\bar{x}_j, t_0}| \right) (Y)
\end{aligned}$$

By Lemma 1.30, the fact that  $s_{\bar{x}_0, t_0 + \frac{1}{\nu}} = s'_{\bar{x}_0, t_0} = \bar{0}$ , we have, for some  $F \in \mathcal{R}_{>0}$ , that;

$$\max(\{|s_{\bar{x}_j, t_0 + \frac{1}{\nu}}| : 1 \leq j \leq 3\} \cup \{|s'_{\bar{x}_j, t_0}| : 1 \leq j \leq 3\}) \leq \frac{F}{\sqrt{\eta}}$$

so that, using (Y);

$$\begin{aligned}
|(\bar{v})'_\nu(t_0) - (s_{\bar{x}_0})'_\nu(t_0)| &\leq 3\sqrt{3} D \sqrt{\eta} \left( \frac{3F^2}{\eta} \right) + 3\sqrt{3} E \sqrt{\eta} \left( \frac{6F^2}{\eta} \right) \\
&\leq \frac{\sqrt{3}(9D+18E)F^2}{\sqrt{\eta}} = \frac{G}{\sqrt{\eta}}
\end{aligned}$$

Suppose that  $(s_{\bar{x}_0}(t_0), t_0) \in {}^*B(\bar{r}_0, \epsilon) \times (t_1, t_2)$ , where  $\epsilon \in \mathcal{R}_{>0}$ ,  $t_1 < t_2$ ,  $\{t_1, t_2\} \subset \mathcal{R}_{>0}$ ,  $\bar{r}_0 \in \mathcal{R}^3$ . Choose a smooth function  $h(\bar{x}, t)$ , with  $h|_{B(\bar{r}_0, \epsilon) \times (t_1, t_2)} = \bar{J}|_{B(\bar{r}_0, \epsilon) \times (t_1, t_2)}$  and  $h$  supported on  $B(\bar{r}_0, \epsilon') \times (t'_1, t'_2)$ , with  $\epsilon < \epsilon'$ ,  $t'_1 < t_1 < t_2 < t'_2$ ,  $\epsilon' \in \mathcal{R}_{>0}$ ,  $\{t'_1, t'_2\} \subset \mathcal{R}_{>0}$ . Then, clearly,  ${}^*h(s_{\bar{x}_0}(t_0), t_0) = {}^*\bar{J}(s_{\bar{x}_0}(t_0), t_0)$ . Let  $\bar{J}_1(\bar{x}, t) = \bar{J}(\bar{x}, t) - h(\bar{x}, t)$ , then, we have that  ${}^*\bar{J}_1(s_{\bar{x}_0}, t_0) = \bar{0}$ . By the result of Lemma 2.11, replacing  $(\rho, \bar{J})$  with  $(\rho, \bar{J}_1)$ , we can find,  $\bar{y}_0 \in R$ , with;

$$|s_{\bar{y}_0}^\sharp(t_0) - s_{\bar{x}_0}(t_0)| \leq .$$

□

## REFERENCES

- [1] A Non Standard Representation for Ito Integration and Brownian Motion, Robert Anderson, Israel Journal of Mathematics, Vol 25, (1976)
  - [2] Nonstandard Analysis in Practice, Francine and Marc Diener ed., Springer, (1995)
  - [3] Introduction to Electrodynamics, David Griffiths, Pearson, (2008).
  - [4] Conversion from Nonstandard to Standard Measure Spaces and Applications in Probability Theory, Peter Loeb, Transactions of the AMS, Volume 211, (1975)
  - [5] Advances in Nonstandard Analysis, Tristram de Piro, available at <http://www.curvalinea.net>
  - [6] Some Arguments for the Wave Equation in Quantum Theory 2, Open Journal of Mathematical Sciences, Tristram de Piro, (2022).
  - [7] Solving Schrodinger's Equation using Nonstandard Analysis, Tristram de Piro, available at <http://www.curvalinea.net>, (2017).
  - [8] Non-Standard Analysis, Abraham Robinson, Princeton University Press, (1996).
  - [9] A theorem on planar graphs, W. T. Tutte, Transactions of the American Mathematical Society 82, 99116, (1956)
- .